

## Intra Beam Scattering in RHIC

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# Intra Beam Scattering in RHIC

## I. Introduction

## II. IBS beam growth and beam loss

emittance, intensity, luminosity

## III. IBS Scaling laws

emittance growth, below & above transition

beam loss, Fokker-Planck equation

## IV. Future

Jie Wei

July 19, 1995

Talk at

"Workshop on Machine Backgrounds at RHIC"

# I. Introduction

IBS : Intra-beam multiple Coulomb scattering

⇒ main cause of beam growth & loss in RH

\* most severe for high charge state ions

scattering cross section  $\sim Z^4/A^2$

\* severe for high intensity beam

$\sim N$  (number of particle per bunch)

\* severe for low emittance beam

$\sim \epsilon_x^{-1} \epsilon_y^{-1} S^{-1}$

\* typically slow process

growth time  $\gg$  synch. osc. period

\* Theoretically approached by many people

A. Piwinsky (1974)

J. Bjorken and S. Mtingwa (1983)

⋮

\* Comparison with experimental study

Within a factor of 2 on growth rate

(theoretical over-estimate, typical)

## II. IBS beam growth and beam loss

$\boxed{Au^{79+}}$   $10^9$  per bunch,  $\sim 57$  bunches

\* transverse emittance growth

$$\epsilon_N: 10 \pi \text{ mm}\cdot\text{mr} \rightarrow 40 \pi \text{ mm}\cdot\text{mr}$$

\* longitudinal growth, beam loss

fill rf bucket in  $< 1$  hour

then significant beam loss ( $\sim 40\%$  in  $10^4$ )

\* luminosity reduction

$$\beta^* = 2 \text{ m}: 1.1 \times 10^{27} \rightarrow 1 \times 10^{26} \text{ (cm}^{-2}\text{s}^{-1}\text{)}$$

$$\beta^* = 1 \text{ m}: 2.2 \times 10^{27} \rightarrow 2 \times 10^{26} \text{ (cm}^{-2}\text{s}^{-1}\text{)}$$

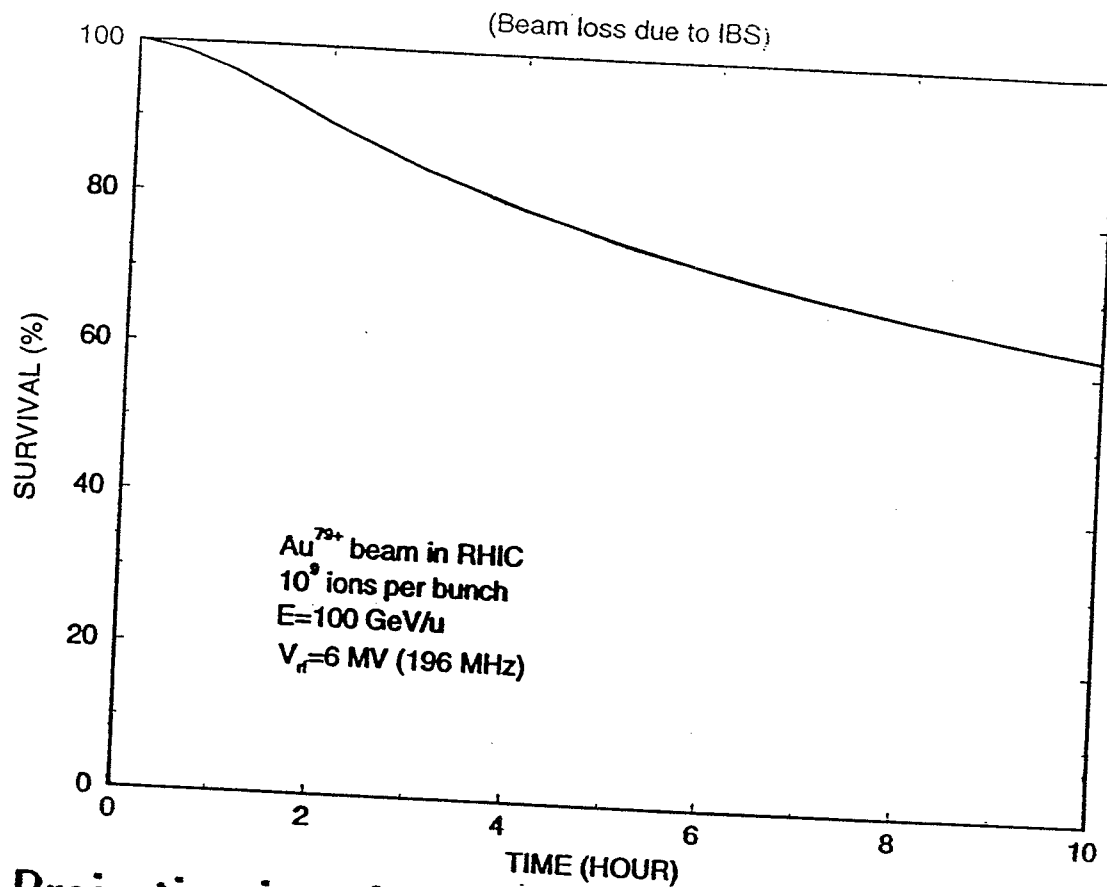
$\boxed{p}$   $10^{11}$  per bunch,  $\sim 57$  bunches

\* transverse:  $20 \pi \text{ mm}\cdot\text{mr} \rightarrow 30 \pi \text{ mm}\cdot\text{mr}$

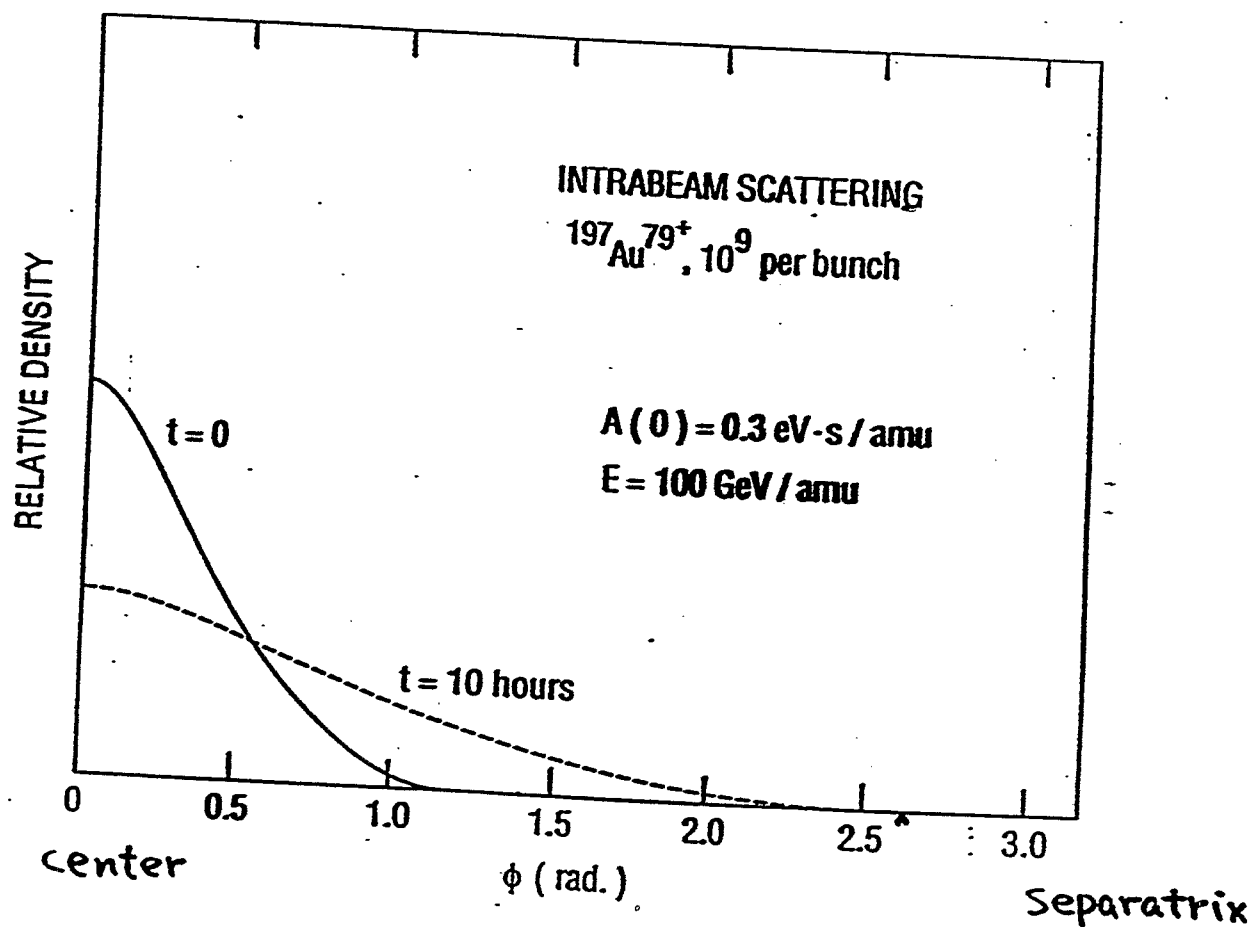
\* longitudinal:  $0.3 \text{ eV}\cdot\text{s} \rightarrow 1 \text{ eV}\cdot\text{s}$

(rf bucket area  $3.1 \text{ eV}\cdot\text{s}$ )

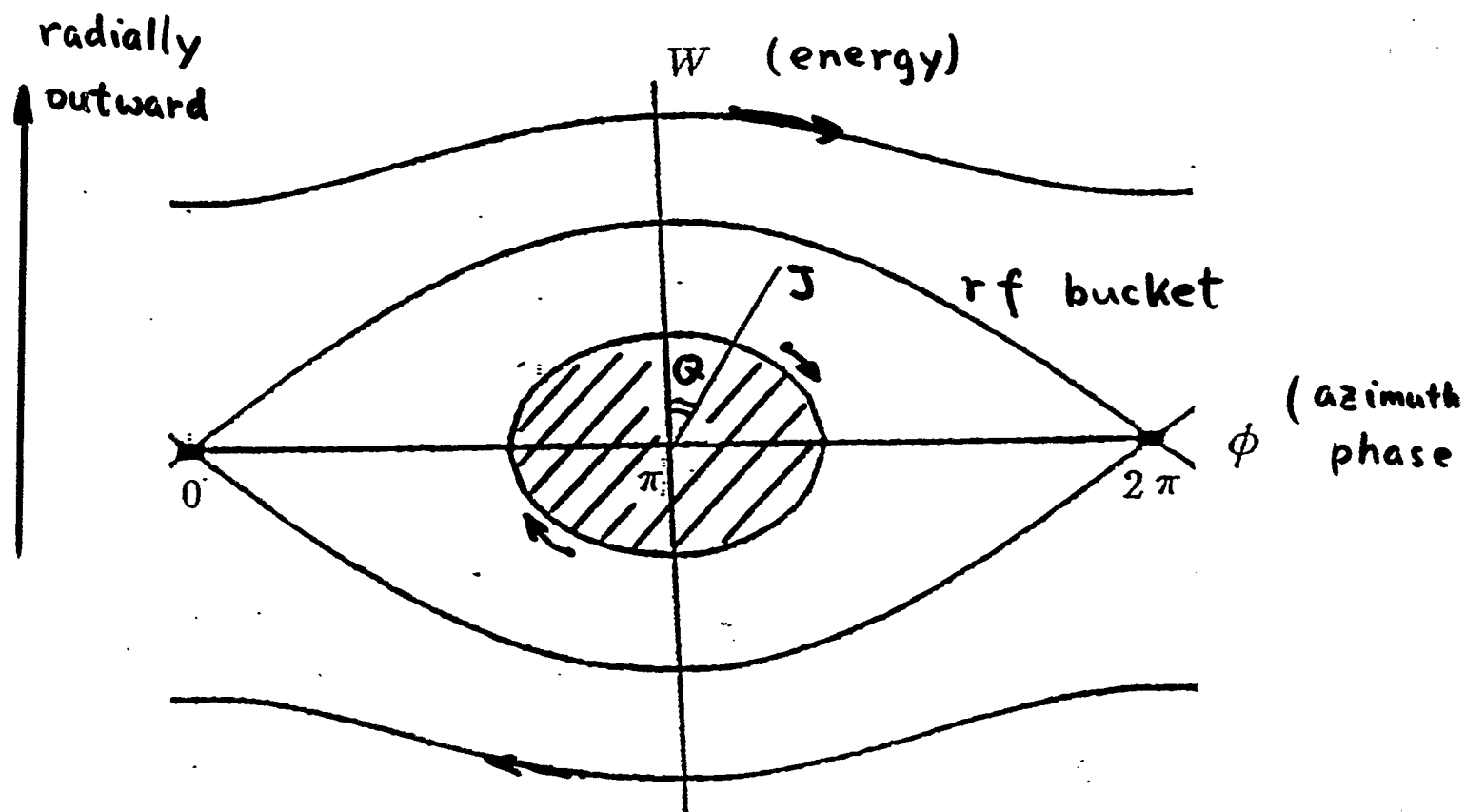
\* no beam loss



Projection in azimuthal direction ( $\phi$ ):



# IBS beam loss



\* longitudinal: (mainly)

IBS diffusion

⇒ particles escape out of the bucket

⇒ becomes dc background, beam halo,

or trapped in empty buckets

transverse: ( $\beta^* < 1$  m operation)

particles of large emittance (action)  $\epsilon_N \gtrsim 40 \pi$

hit physical aperture

Momentum aperture:

depends on  $\beta^*$  in operation

$$\beta^* = 1 \text{ m,}$$

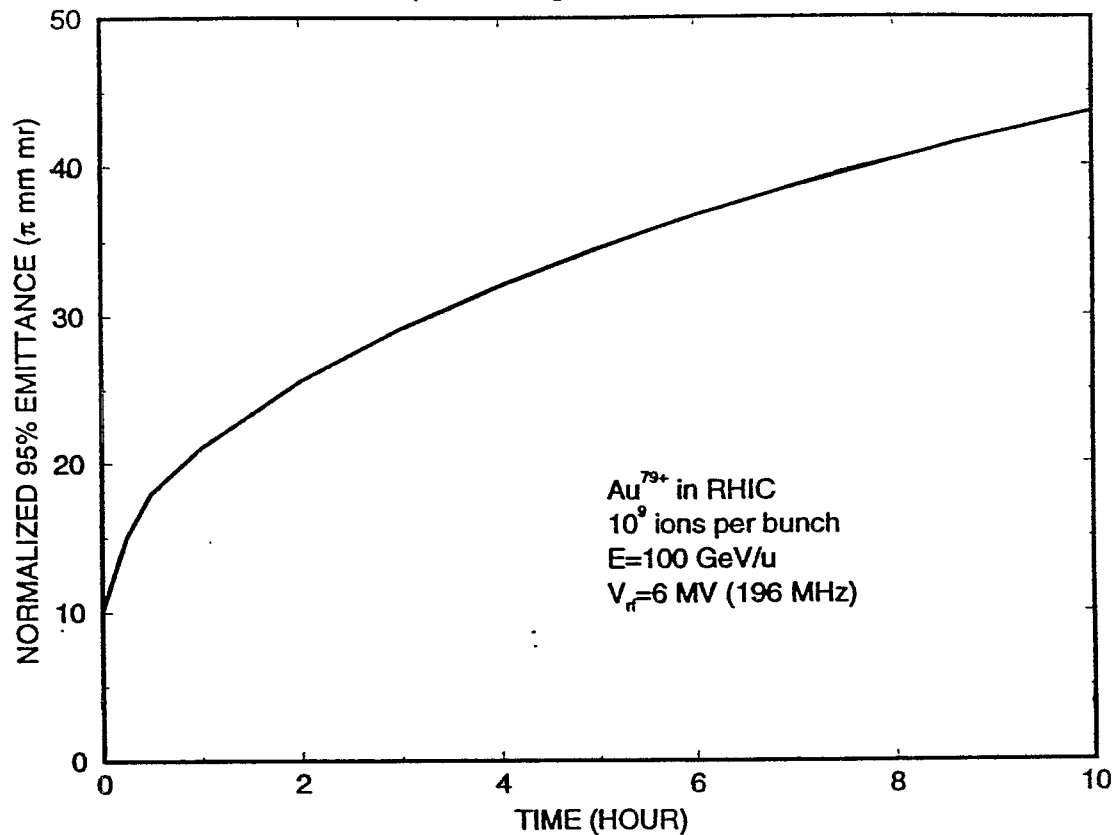
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$\frac{\Delta p}{p}$	Dynamic Aperture	
0	5.5 $\sigma$	O.K.
0.002 (95%)	4.5 $\sigma$	marginal

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## EMITTANCE

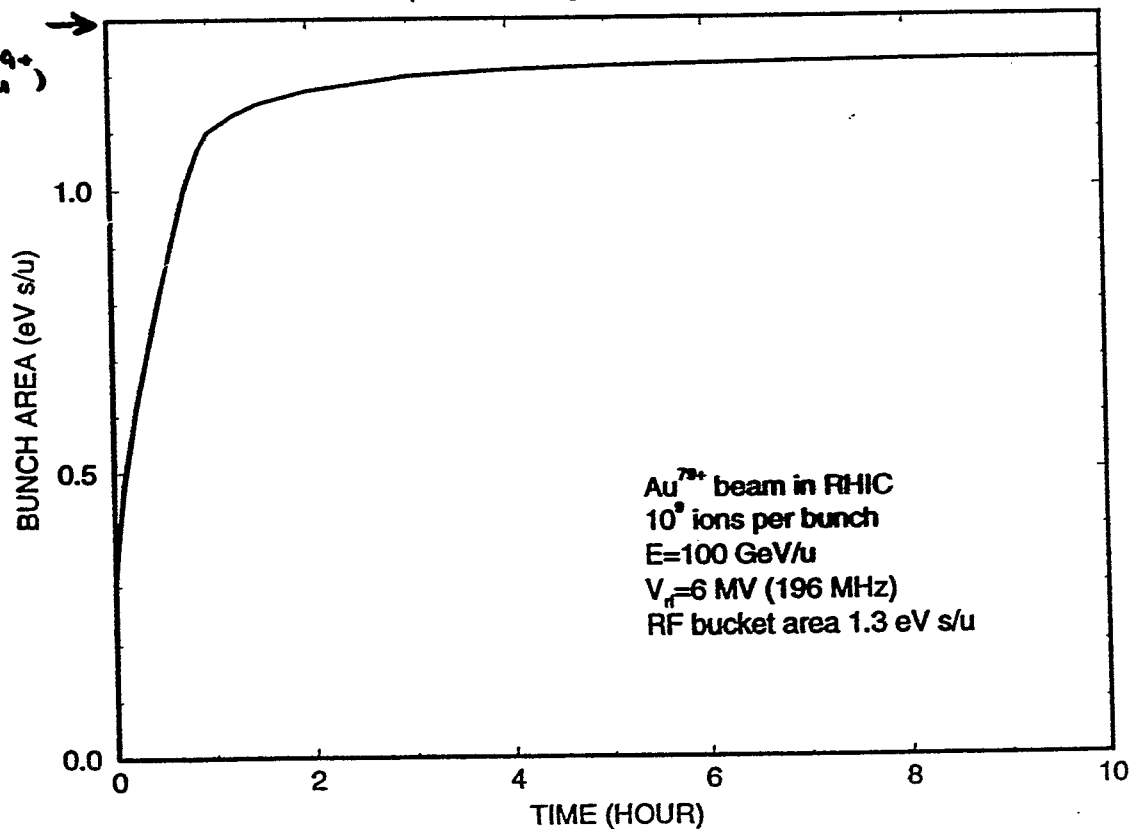
(Emittance growth due to IBS)



## BUNCH AREA

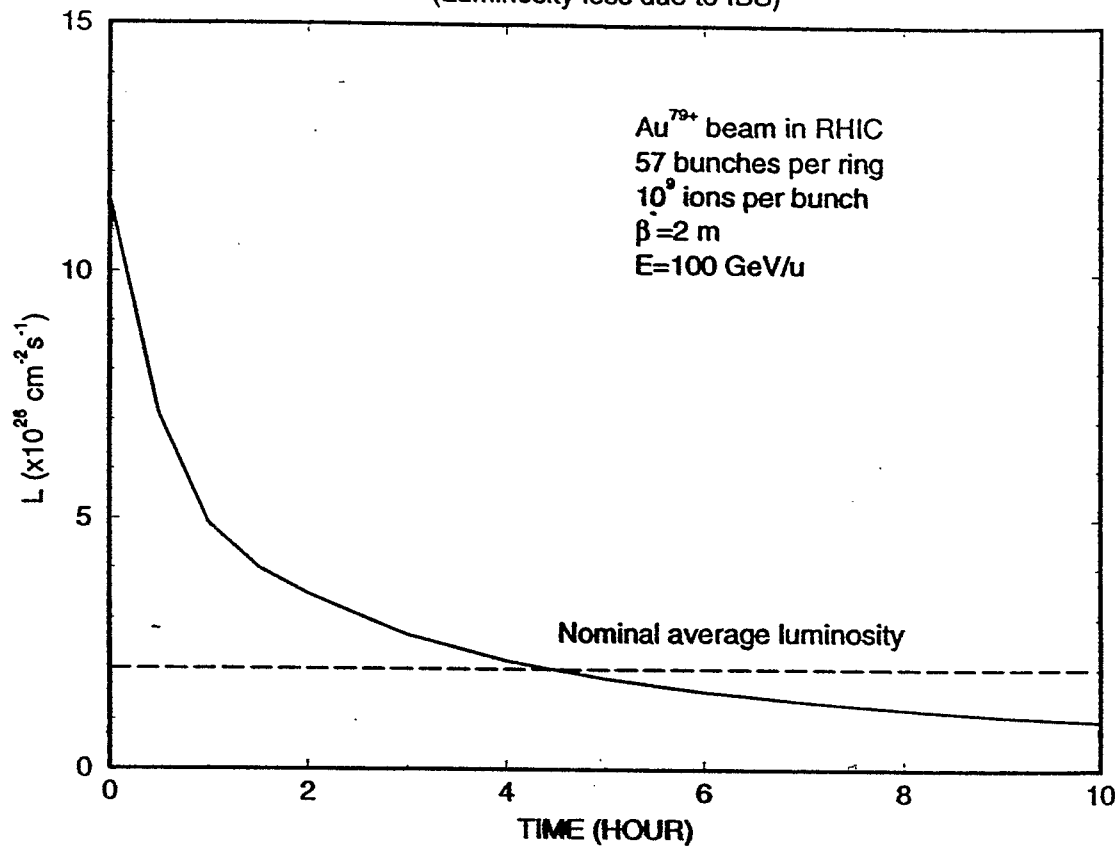
(Bunch area growth due to IBS)

bucket  
area (Au<sup>79+</sup>) →



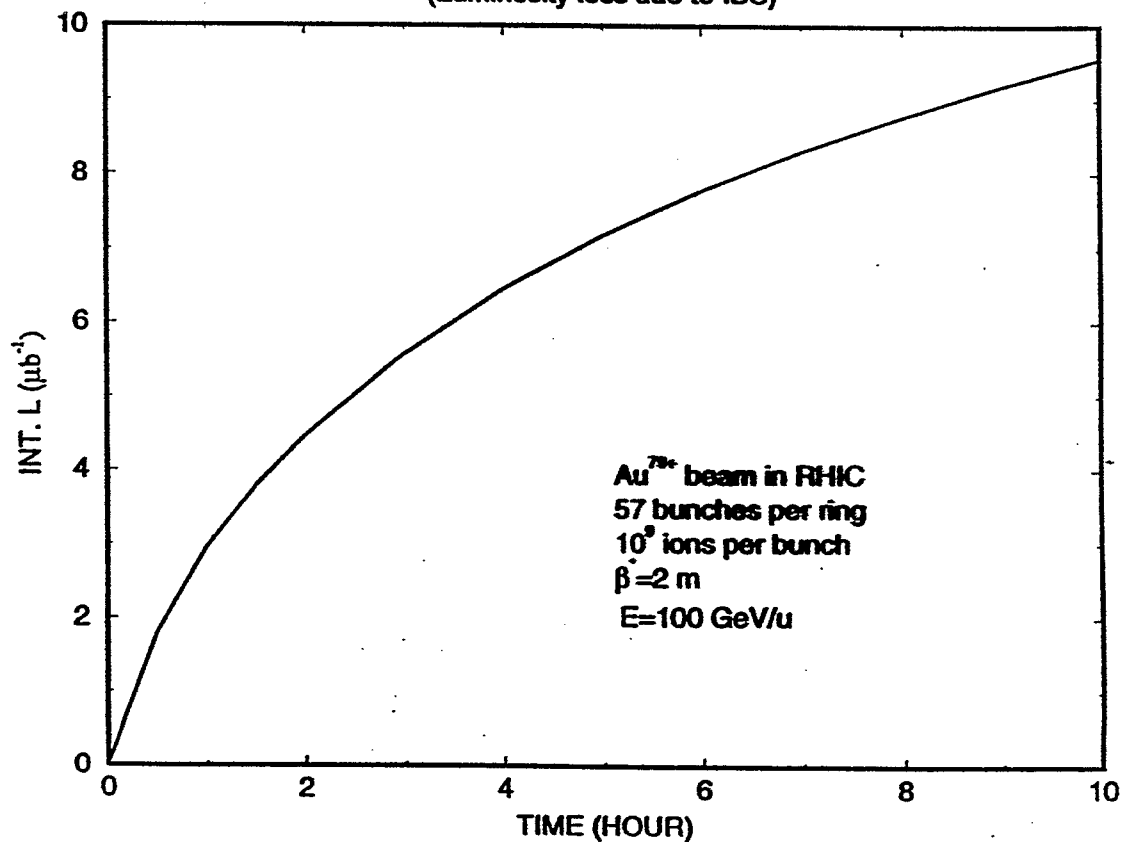
# INSTANTANEOUS LUMINOSITY

(Luminosity loss due to IBS)



# INTEGRATED LUMINOSITY

(Luminosity loss due to IBS)



## Luminosity for a round beam

$$\mathcal{L} = \frac{3}{2\pi} f_{\text{rev}} \cdot N_B \cdot N^2 \cdot \frac{\beta\gamma}{\epsilon_N \cdot \beta^*}$$

$N_B$ : number of bunches

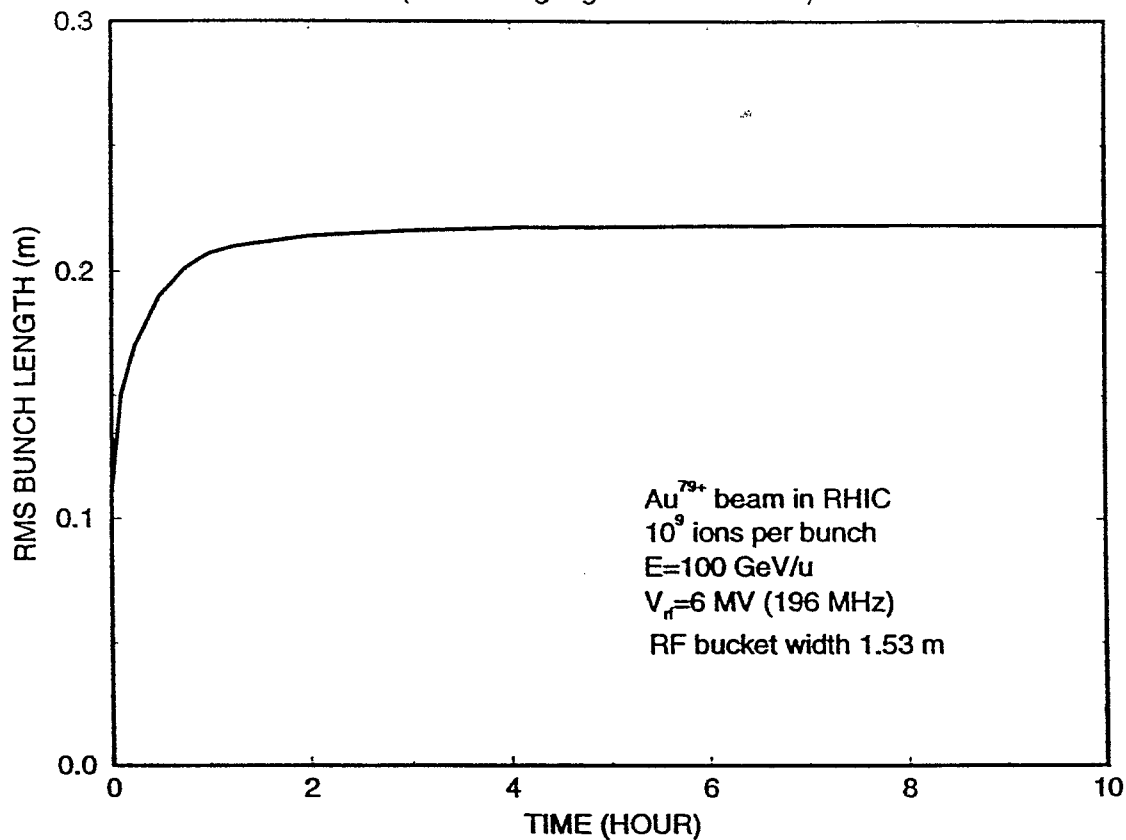
$N$ : number of ions per bunch

$$\epsilon_N : \frac{6 \cdot \beta\gamma \sigma_x \sigma_y}{\beta^*}$$

$f_{\text{rev}}$ : revolution frequency

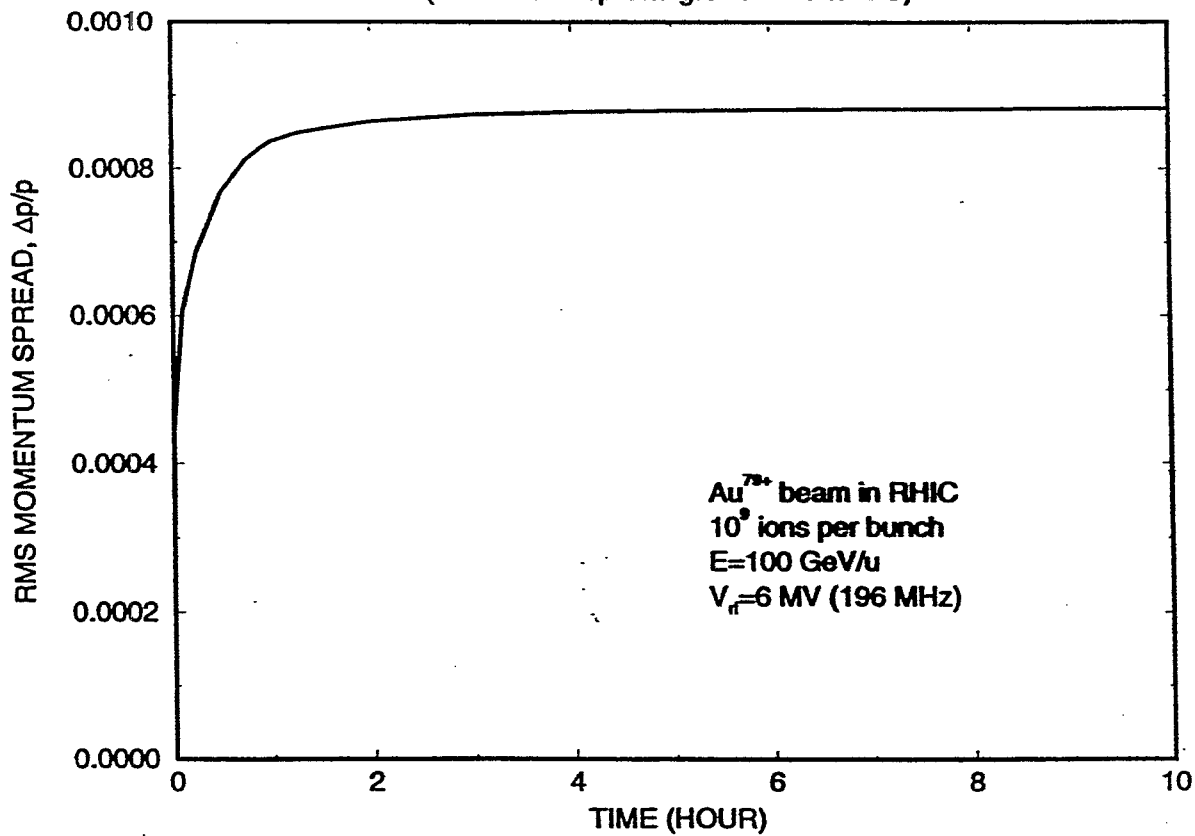
# BUNCH LENGTH

(Bunch length growth due to IBS)



# MOMENTUM SPREAD

(Momentum spread growth due to IBS)



### III. IBS scaling laws

#### \* below transition energy :

similar to gas scattering

⇒ asymptotically approaches isotropic state

$$\langle \frac{\sigma_x}{\beta_x} \rangle \sim \langle \frac{\sigma_y}{\beta_y} \rangle \sim \frac{\sigma_p}{\gamma}$$

growth due to machine lattice variation

heat absorption in a time-dependent system

#### \* above transition energy :

no equilibrium due to negative mass

⇒ asymptotically dispersion related horizontal

and longitudinal dimension

$$n_c \cdot \langle \sigma_x^2 \rangle \sim \langle x_p \rangle^2 \sigma_p^2$$

$$n_c = \begin{cases} 1 & \text{uncoupled} \\ 2 & \text{coupled} \end{cases}$$

continuous growth in both horizontal and longitudinal dimension

Beam growth at high energy  $\gamma \gg \gamma_T$

$$\begin{bmatrix} \frac{1}{\sigma_p} \frac{d\sigma_p}{dt} \\ \frac{1}{\sigma_x} \frac{d\sigma_x}{dt} \end{bmatrix} = \frac{Z^4 N}{A^2} \frac{C_0}{\gamma_T \epsilon_x \epsilon_y S} \cdot \begin{bmatrix} (1-d^2)/d \\ d/n_c \end{bmatrix}$$

$$C_0 \equiv \frac{\pi r_0^2 m_0 c^2 L_c}{16}$$

$$d = \langle x_p \rangle \sigma_p / (\sigma_x^2 + \langle x_p^2 \rangle \sigma_p^2)^{1/2}; \quad 0 < d < 1$$

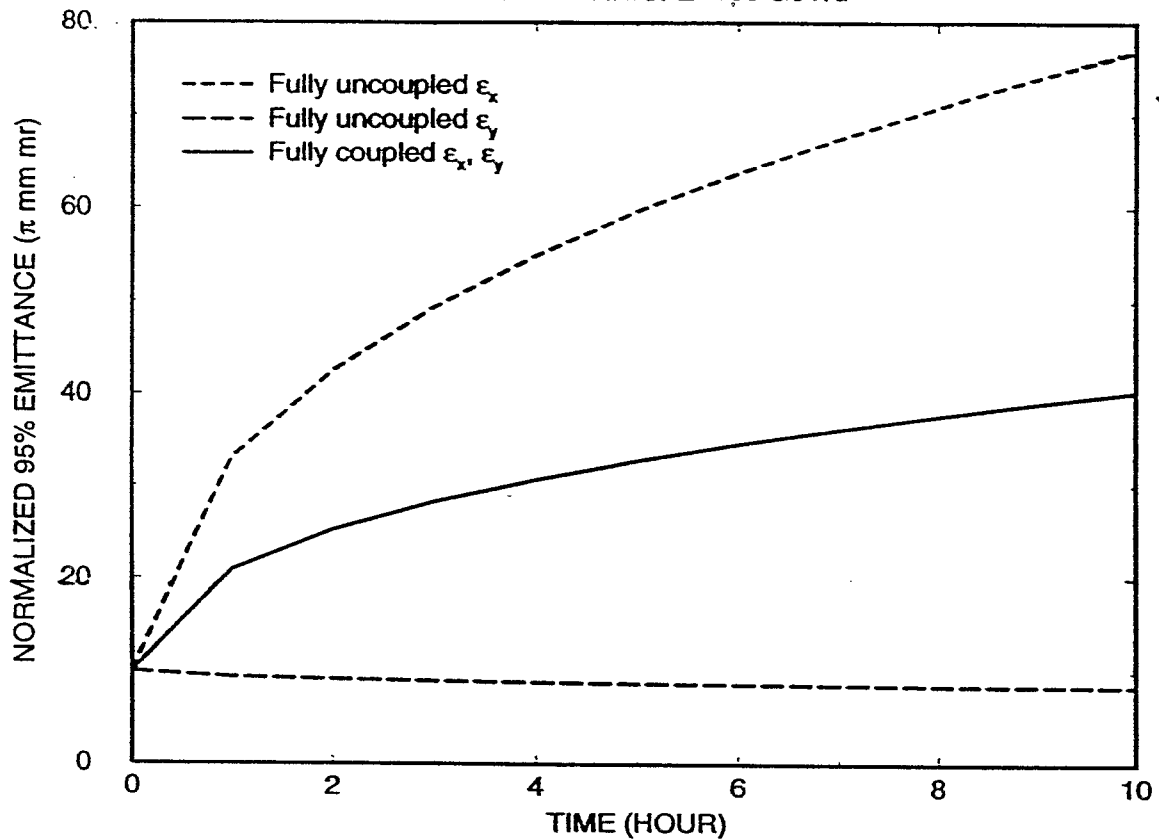
\* not sensitive to energy

$$* \quad \tau^{-1} \sim \frac{Z^4}{A^2} \frac{N}{\epsilon_x \epsilon_y S}$$

\* transverse growth sensitive to coupling

# EFFECTS OF COUPLING ON IBS

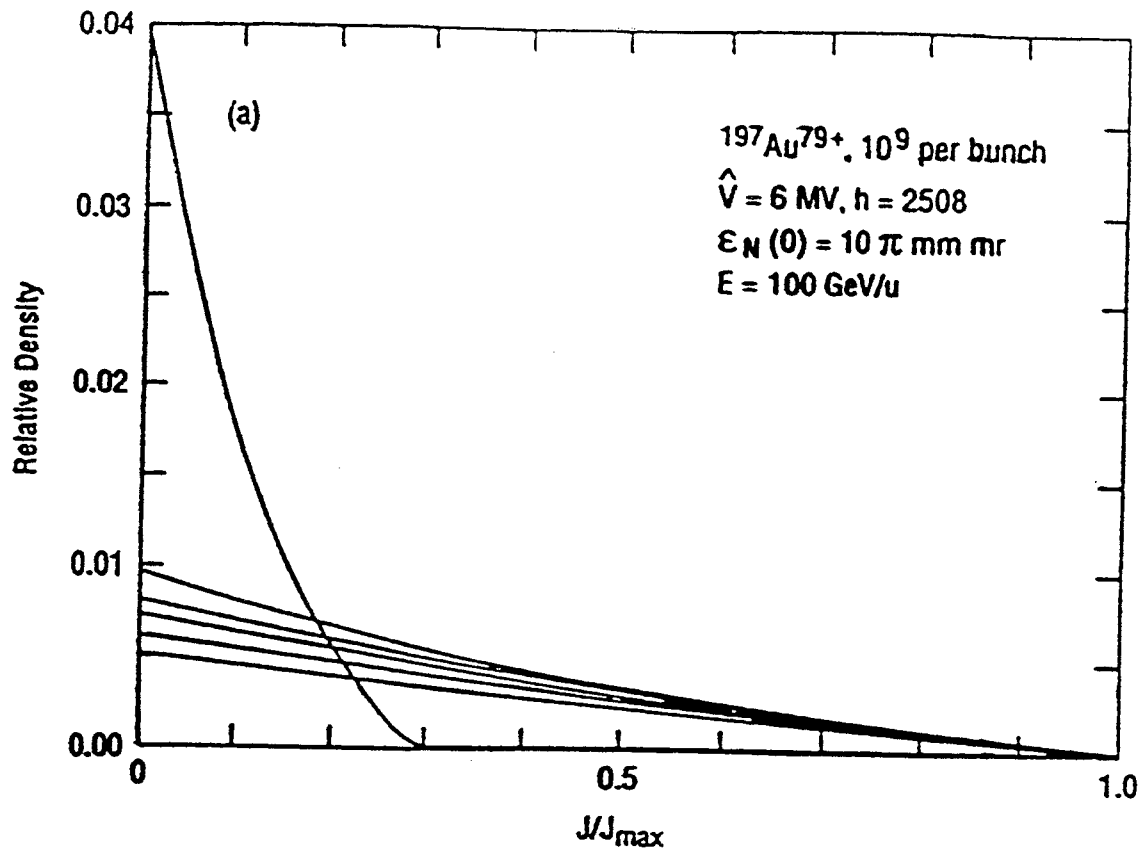
Au<sup>79+</sup> beam in RHIC. E=100 GeV/u



beam loss

for  $\beta^* = 1$  m operation :

dynamic aperture limit  $\Rightarrow \epsilon_{x,y} \sim 40 \pi$  mm.m



Use Fokker - Planck equation to evaluate beam evolution and beam loss

$$\frac{\partial \psi}{\partial t} = - \frac{\partial}{\partial J} (F(J) \psi) + \frac{1}{2} \frac{\partial^2}{\partial J^2} (D(J) \frac{\partial \psi}{\partial J})$$

$J$ : action

$\psi$ : density distribution function

## IV. Future

- \* Increase the number of bunches per ring
- \* An upgrade in rf system (to 16 MV) can reduce beam loss to 2%.

The gain in luminosity is  $\sim 30\%$ .

may have aperture problem with  $\beta^* = 1\text{m}$  operation and  $\epsilon_N \sim 48 \pi \text{ mm}\cdot\text{mr}$

- \* Methods like stochastic cooling are desirable to preserve beam quality, alleviate dynamic aperture requirement, and increase luminosity.

