



BNL-101717-2014-TECH

RHIC/AP/62;BNL-101717-2013-IR

Intrabeam Scattering Results for a High Frequency RF System

G. Parzen

March 1988

Collider Accelerator Department
Brookhaven National Laboratory

U.S. Department of Energy

USDOE Office of Science (SC)

Notice: This technical note has been authored by employees of Brookhaven Science Associates, LLC under Contract No. DE-AC02-76CH00016 with the U.S. Department of Energy. The publisher by accepting the technical note for publication acknowledges that the United States Government retains a non-exclusive, paid-up, irrevocable, world-wide license to publish or reproduce the published form of this technical note, or allow others to do so, for United States Government purposes.

DISCLAIMER

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, nor any of their contractors, subcontractors, or their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or any third party's use or the results of such use of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof or its contractors or subcontractors. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

AD/RHIC-AP-62

Intrabeam Scattering Results for a
High Frequency RF System

G. Parzen

3/25/88

The RF system that was proposed
 has the harmonic number $h = 8 \times 342 = 2736$,
 corresponding to a frequency of about
 $f = 214 \text{ MHz}$. It was found that
 this RF system requires a voltage larger
 than 32 MV in order for
 the RF bucket to be large enough
 to contain the energy spread for
 a beam of Gold ions, after 10 hours,
 at $\gamma = 100$ with $N_b = 1.1 \times 10^9 \text{ ions/bunch}$,
 and ~~b~~ an initial bunch area
 $A = 1 \text{ ev-sec}/\text{amu}$

Intra beam scattering causes the
 beam to grow as shown in the
 following table for Au ions after $10 \frac{\text{hrs}}{\cancel{\text{hrs}}}$,
 and an initial bunch area of $A = 1 \text{ ev-sec}/\text{amu}$
 at $\gamma = 100$.

Table 1 - $\gamma = 100$		final, after 10 hrs						
h	V MV	initial						
		$\sigma_{p_0}/10^3$ cm	σ_{l_0}	$\sigma_p/10^3$ cm	σ_x mm	E_x mm/mrad	$2.5\sigma_p/10^3$ $\Delta p/p/10^{-3}$	bucket
2736	32	1.33	12.8	2.00	19.2	2.2	58	5.0 4.7
342	1.2	0.36	48	1.07	142	1.61	28	2.67 2.6

For comparison, the results for the old RF system with $h = 342$, $V = 1.2 \text{ MV}$ are also shown.

There is a considerable improvement in the bunch length which grows to only $\bar{\tau}_e = 19 \text{ cms}$ after 10 hrs. However, the other dimensions have grown ~~more~~ larger, $\epsilon_x = 58$, $2.5 \sigma_p = 5 \times 10^{-3}$ instead of the old results $\epsilon_x = 28$, $2.5 \sigma_p = 2.6 \times 10^{-3}$.

The 6 σ rule requires a stability limit of $A_{SL} = 9.3 \text{ mm}$ for $\sigma_x = 2.2 \text{ mm}$. ~~the tracking + syn~~ and for a ~~syn~~ particle with $\epsilon_x = \epsilon_y$. The tracking result for a $B^* = 3 \text{ m}$ lattice gives $A_{SL} = 9 \text{ mm}$ at $\sigma_p/p = 5 \times 10^{-3}$.

The larger energy spread will also cause a larger V-spread in the beam due to the ^{iron} saturation generated by and due to random field errors.

The beam growth at $\gamma = 30$ is shown in the following table

Table 2		initial		final after 10 hrs					
h	V	$\sigma_{p_0}/10^{-3}$	σ_{e_0}	$\sigma_p/10^{-3}$	σ_e	σ_x	ϵ_x	$2.5\sigma_p/10^{-3} \Delta p/p/10^{-3}$	bucket
MV			cm		cm	mm	mm.mrad		
2736	32	4.45	12.9	4.42	12.8	4.0	57	11	15
342	1.2	1.26	45	1.99	71	3.0	32	5.0	10.0

The bunch length is again considerably improved; the bunch length grows to only $\sigma_e = 12.8$ cm after 10 hours. The other dimensions have grown larger, $\epsilon_x = 57$, $2.5\sigma_p = 11 \times 10^{-3}$ instead of the old results $\epsilon_x = 32$, $2.5\sigma_p = 5 \times 10^{-3}$

The 6 σ rule requires a stability limit of $A_{SL} = 17$ mm for $\sigma_x = 4$ mm and for a particle with $\epsilon_x = \epsilon_y$. The tracking results for a $\beta^* = 6$ m lattice give $A_{SL} = 7$ mm at $\Delta p/p = 11 \times 10^{-3}$ and $A_{SL} = 17$ mm at $\Delta p/p = 0$. This situation can be somewhat improved by reducing the voltage at $\gamma = 30$.

If the RF Voltage V is reduced available
 the energy spread in the $\frac{\Delta E}{E}$ beam will grow
 to the boundary of the bucket in
 less than 10 hrs.
 The following table shows the
 time it takes for the beam
 to reach the edge of the bucket
 as a function of the rf Voltage
 for a beam of Au ions at $\gamma = 100$

V MV	Time to reach bucket (hrs)	$\delta p/p$ $/10^{-3}$	Final σ_x mm	Initial A ev-sec/nm	Final $\sigma_p/10^{-3}$ σ_e (cm)	Final $\frac{1}{\sigma_p} \frac{d\sigma_p}{dt}/10^{-2}$ (hr^{-1})
32	5.76	4.7	1.99	1	1.38	18 1.87
20	1.4	3.7	1.5	1	1.48	18 6.50
15	.56	3.2	1.28	1	1.28	18 13.0
10	.10	2.62	1.02	1	1.05	18 36.0

Possible Conclusions

The high frequency rf system $h = 2736$, $f = 214 \text{ MHz}$ might be considered acceptable. It has less margin for error. It ~~will probably~~ ^{may} lead to shorter beam lifetimes (less than 10 hrs), and to ~~smaller~~ smaller luminosities.

The high frequency 214 MHz RF system has the advantage of a shorter bunch length. In other respects, it appears to give less favorable results than the present 27 MHz RF system.

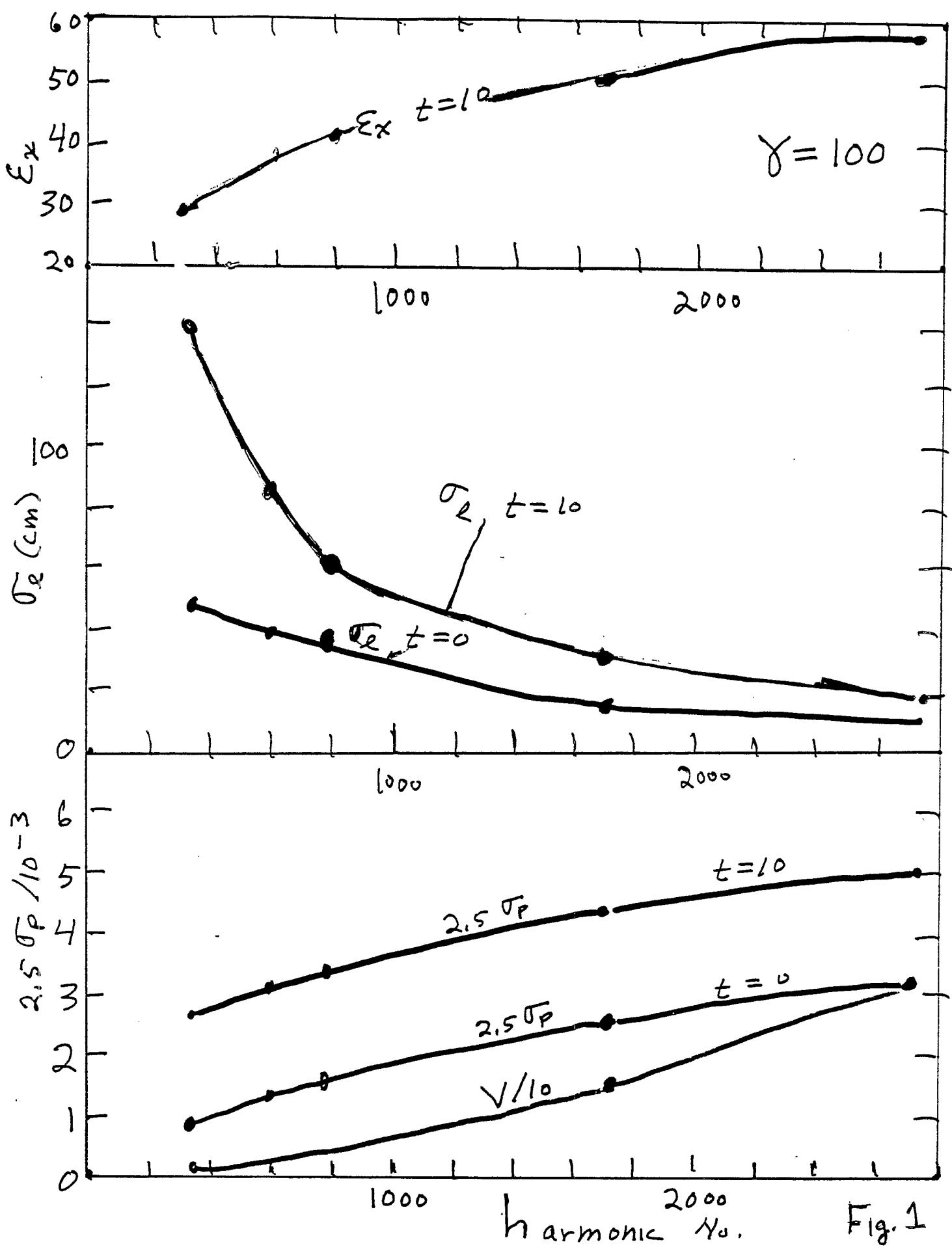
Frequency Dependence and Other Solutions

Figures 1 and 2 show how the intra beam scattering results depend on the harmonic number h or the frequency.

For each harmonic number, the voltage V has been adjusted so that the beam energy spread will just fit inside the rf bucket at $\delta=100$ after 10 hours

Figs. 1 and 2 show σ_p , σ_e and E_x ~~after~~ after 10 hours as a function of the harmonic number h .

Fig 1. shows results at $\delta=100$, Fig 2 shows results at $\delta=30$



harmonic No.

Fig. 1

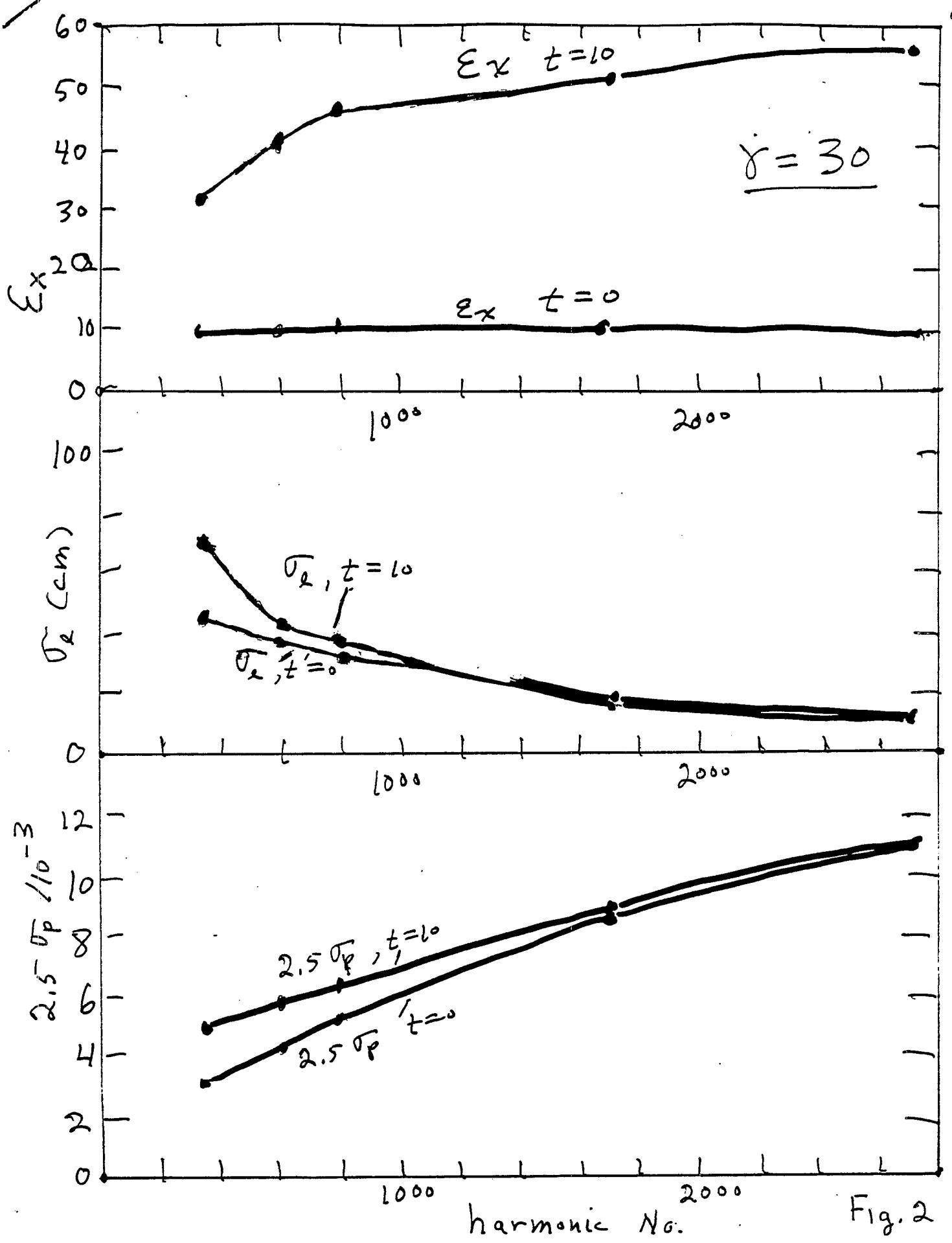


Fig. 2

Figs 1 and 2 show that, except for the growth in σ_e , the performance generally improves at the lower frequencies, which give smaller values for the final σ_p and E_x .

~~Other~~ other solutions, at lower frequencies, that may deserve consideration are the following:

Solution A. $h = 785$, $f = 61 \text{ MHz}$, $V = 4.7 \text{ MV}$

γ	30	100
$\sigma_e(t=0) \text{ cm}$	26	28
$\sigma_e(t=10) \text{ cm}$	31	61
$E_x(t=10)$	45	41
$2.5\sigma_p(t=10)/10^{-3}$	6.4	3.4

This solution has $\sigma_e = 28$ at $t=0$ at $\gamma=100$. The addition of stochastic cooling would be required to keep σ_e from growing.

Solution B. $h = 1700$, $f = 130 \text{ MHz}$, $V = 16.5 \text{ MV}$

γ	30	100
$\sigma_e (t=10) \text{ cm}$	18	28
$\varepsilon_x (t=10)$	52	52
$2.5 \sigma_p (t=10) / 10^3$	9	4.5

This solution has $\sigma_e \leq 28 \text{ cms}$ for $t = 10 \text{ hours}$, and $V = 16.5 \text{ MV}$

Solution C $h = 2052$, $f = 160 \text{ MHz}$, $V = 20 \text{ MV}$

γ	30	100
$\sigma_e (t=10)$	18	26
$\varepsilon_x (t=10)$	55	55
$2.5 \sigma_p (t=10) / 10^3$	10	4.6