

Booster RF Program for Heavy Ions

A. G. Ruggiero

August 1985

Collider Accelerator Department
Brookhaven National Laboratory

U.S. Department of Energy

USDOE Office of Science (SC)

Notice: This technical note has been authored by employees of Brookhaven Science Associates, LLC under Contract No. DE-AC02-76CH00016 with the U.S. Department of Energy. The publisher by accepting the technical note for publication acknowledges that the United States Government retains a non-exclusive, paid-up, irrevocable, world-wide license to publish or reproduce the published form of this technical note, or allow others to do so, for United States Government purposes.

DISCLAIMER

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, nor any of their contractors, subcontractors, or their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or any third party's use or the results of such use of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof or its contractors or subcontractors. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

Booster rf Program for Heavy Ions

A. G. Ruggiero and G. Young

BNL

8/22/85

BOOSTER RF PROGRAM FOR HEAVY IONS

A. G. Ruggiero and G. Young

BNL

Aug. 22, 1985

In this note we present a scenario for an $h = 1$ rf system to accelerate heavy ions in the Booster. Relevant Booster parameters are

$$2\pi R = 201.75 \text{ m}$$

$$\gamma_T = 4.5$$

and for the beam of gold, the heaviest of the species considered,

$$q = 33 \text{ (partially stripped ions)}$$

$$A = 197$$

$$N = 2 \times 10^9 \text{ (total number of particles)}$$

and also

$$\beta = 0.0478 \text{ at injection}$$

$$= 0.6684 \text{ at top energy}$$

Our considerations are based on two equations:

- i. The rate of acceleration
- ii. The ratio of the rf bucket area to the beam bunch area.

As we said, we shall assume $h = 1$ with one single bunch.

1. The rate of acceleration

The energy gain per turn per nucleon is

$$\Delta E = e \frac{q}{A} V \sin \phi_s$$

where V is the peak voltage and ϕ_s the synchronous phase angle of the rf waveform.

The energy gain per unit of time

$$\dot{E} = \frac{\Delta E \beta c}{2\pi R} = \frac{eq\beta c}{2\pi R A} V \sin \phi_s \quad (1)$$

but

$$E = E_0 \gamma \text{ with } E_0 = 0.9315 \text{ GeV/A}$$

and also

$$\dot{E} = E_0 \gamma^3 \dot{\beta} \quad (2)$$

Combining (1) and (2) gives

$$\dot{\beta} = \frac{eqc}{2\pi R A E_0 \gamma^3} V \sin \phi_s$$

This is our first important equation which relates $\dot{\beta}$ to the product $\gamma^3 V \sin \phi_s$ as all the other quantities are assigned. For the booster accelerating $^{197}\text{Au}^{33+}$,

$$\dot{\beta} = 2.6722 \cdot 10^{-4} \gamma^{-3} V \sin \phi_s \quad (3)$$

with V in volts and $\dot{\beta}$ in sec^{-1} .

Also, a useful equation which relates $\dot{\beta}$ to \dot{B} , the variation with time of the bending field is

$$\dot{B} \rho = 3.10715 \frac{1}{q} \gamma^3 \dot{\beta} \quad (4)$$

where ρ is the bending radius in meters and B is measured in Tesla. For $\beta \ll 1$, that is non-relativistic case, $\dot{\beta}$ constant implies also \dot{B} constant.

2. The rf bucket area

In general,

$$B = \frac{16R}{h^2 c} \alpha(\phi_s) \left[\frac{he V q E_0 \gamma}{2\pi |\eta| A} \right]^{1/2} \quad (5)$$

where $\alpha(\phi_s)$ is a tabulated function of ϕ_s and $\eta = \frac{1}{\gamma_T^2} - \frac{1}{\gamma^2}$.

i. Capture

$$\begin{aligned}\phi_s &= 0^\circ \\ \alpha &= 1 \\ V &= 450 \text{ Volts}\end{aligned}$$

which gives at the beginning

$$B_o = 0.1837 \text{ (eV}\cdot\text{sec)/A} \quad (6)$$

The ion beam is injected with RF voltage on. It is allowed to fill the bucket completely. Therefore B_o is also the bunch area.

ii. The voltage is raised adiabatically with no acceleration ($\phi_s = 0$) to double its original value, that is up to 900 volts.

The new bucket area is

$$B_a = 0.2598 \text{ (eV}\cdot\text{sec)/A}$$

and the bunch area is still given by (6). The ratio of B_o and B_a represents at this time a filling factor of $\frac{1}{\sqrt{2}}$.

iii. Acceleration

We propose that the programs of V and ϕ_s for acceleration be chosen so as to preserve the original filling factor of 0.7 during the entire cycle. This criterion would minimize the total amount of rf voltage.

Inserting all the known values in (5) we then derive the second important equation.

$$\frac{\alpha^2(\phi_s) V \gamma}{|\eta|} = 924.8 \text{ volts} \quad (7)$$

Equations (3) and (7) are all that are needed and they are quite general. If, for example, the $\dot{\beta}$ dependence with time is assigned during the entire acceleration cycle, one then derives β and γ , and the two equations (3 and 7) can be solved to yield V and ϕ_s versus time.

Finally the program of the bending field can be obtained from (4).

Examples

Consider the first stage 0-50 msec where at the end $\beta = 3\beta_{inj} = 0.1434$. We can take non-relativistic approximation and set $\gamma \sim 1$.

i. Constant $\dot{\beta}$ (also \dot{B})

$$\dot{\beta} = \frac{3\beta_{inj} - \beta_{inj}}{50 \text{ msec}} = 1.912 \text{ sec}^{-1}$$

with this, from (3) and (7) and consulting the tables for $\alpha(\phi_s)$, we derive

$$\begin{aligned} \phi_s &= 35^\circ, \text{ constant} \\ V &= 12.5 \text{ kV}, \text{ constant} \end{aligned}$$

Observe this program is not realistic because both ϕ_s and V have a discontinuity. The first jumps instantly from 0° to 35° and the latter from 0.9 kV to 12.5 kV.

ii. Parabolic variation of β with time - at 50 msec $\dot{\beta}$ is twice the value at the same instant for constant program. That is

$$\dot{\beta} = 3.824 \text{ at } t = 50 \text{ msec}$$

whereas $\dot{\beta} = 0 \text{ sec}^{-1}$ at the beginning. V and ϕ_s are not constant but vary as one can derive from (3 and 7).

t	0 msec	50 msec
f_{rf}	71.03 kHz	213.09 kHz
ϕ_s	0°	42°
V	0.9 kV	21.4 kV

iii. Exponential variation.

$$\beta = \beta_{inj} \exp(Kt)$$

$$\text{with } K = 21.972 \text{ sec}^{-1},$$

and

$$\dot{\beta} = (1.0503 \text{ sec}^{-1} \exp(Kt))$$

The results are summarized in the table:

t	0 msec	50 msec
$\dot{\beta}$	1.0503 s ⁻¹	3.1508 s ⁻¹
ϕ_s	30°	39°
V	7.89 kV	19.37 kV

Here also there is a discontinuity for both ϕ_s and V which have to jump from 0° to 30° for the first and from 0.9 kV to 7.89 kV for the latter.

Probably the best program is a combination of the parabolic ramp until $\dot{\beta}$ reaches the value for the linear ramp, followed by a linear ramp at $\dot{\beta} = 1.912 \text{ sec}^{-1}$, then ending the cycle with a parabolic ramp back to $\dot{\beta} = 0$.

General Procedure

This is to include both equations (3 and 7) into a computer program with the goal of optimizing the best choice of V and ϕ_s according to some criterion that the rf engineering experts can develop (power, flux, losses...).