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Estimated Beam Loss Due to Energy Straggling Through the Charge Exchange Injection Foil

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Introduction

As H⁻ or circulating protons traverses the charge exchange injection foil, it deposits some of its energy in the foil. The energy distribution $f(\Delta, s)$ depends on a parameter κ , which is proportional to the ratio of the mean energy loss over the path length to the largest energy transfer possible in a collision with an atomic electron. If $\kappa \rightarrow 0$, distribution is Landau [1], if κ is between 1 and 10, then the distribution is Vavilov [2], and for $\kappa \ge 10$ then the distribution is Gaussian [2]. For the SNS accumulator ring injection foil, the parameter κ is about 6.61 x 10⁻⁶, and the largest energy transfer possible in a collision with an atomic electron is about 3.3 MeV. Therefore the distribution is Landau, with its tail extending up to 3.3 MeV. The concern was some of the protons, which lie in the tail of Landau distribution, will go through a large synchrotron oscillation and might end up the beam gap. These protons will be lost at time of extraction or at dispersive locations in the ring, and losses might be greater than or equal to 0.01%.

Significant Parameters

The largest energy transfer possible in a collision between a particle of mass 'M' and a free electron of mass 'm' can be calculated by using the conservation of energy and momentum, and is given by

$$\boldsymbol{e}_{\max} = \frac{2mc^2 \, \boldsymbol{b}^2 \boldsymbol{g}^2}{1 + 2\boldsymbol{g}m/M + (m/M)^2} \tag{1}$$

where β , and γ are the relativistic parameters. For 1 GeV proton ε_{max} is 3.3 MeV.

The mean energy loss in path-length s $(g \text{ cm}^{-2})$ is given by the Bethe Bloch formula [3]

$$\overline{\Delta} = 0.30058 \frac{mc^2}{\boldsymbol{b}^2} \frac{Z}{A} s \left[\log \left(\frac{2mc^2 \boldsymbol{b}^2 \boldsymbol{e}_{\max}}{I^2 (1 - \boldsymbol{b}^2)} \right) - 2\boldsymbol{b}^2 - 2\frac{C}{Z} - \boldsymbol{d} \right]$$
(2)

Where Z, A, and I are the atomic mass number, atomic weight, and mean excitation energy of the medium, 2C/Z is the shell correction, and δ is the density effect correction. For 200μ g/cm² thick carbon foil average energy lost by 1Gev proton is 0.44 keV.

The parameter κ is defined as

$$\boldsymbol{k} = \frac{\boldsymbol{x}}{\boldsymbol{e}_{\text{max}}}$$

where $\boldsymbol{x} = 0.30058 \frac{mc^2}{\boldsymbol{b}^2} \frac{Z}{A} s$ (3)

As $\kappa \rightarrow 0$ the distribution f (Δ , s) is Landau, and if $\kappa \gg 1$ then the distribution f (Δ , s) is Gaussian. For other values of κ (=1-10) the distribution f (Δ , s) is Vavilov. For the SNS accumulator ring injection foil, the parameter κ is about 6.61 x 10⁻⁶.

The Landau parameter λ is defined as

$$\boldsymbol{l} = \frac{\Delta - \overline{\Delta}}{\boldsymbol{x}} - 1 - \boldsymbol{b}^2 + C - \log \boldsymbol{k}$$
(4)

Where Δ is the energy loss along the path length s and C =0.577216 is Euler's constant.

The Landau distribution is

$$f_{l}(\Delta, s) = \frac{1}{\mathbf{x}} \mathbf{f}(\mathbf{l})$$
(5)

where $\phi(\lambda)$ is universal function of the variable λ only [2] and given by

$$\boldsymbol{f}(\boldsymbol{l}) = \frac{1}{\boldsymbol{p}} \int_{0}^{\infty} e^{-\boldsymbol{p}y/2} \cos(y \log y + \boldsymbol{l}y) dy$$
(6)

The Landau distribution f_1 is asymmetric, with a tail extending to ε_{max} ; it peaks at λ =-0.233 and has a full width at half maximum of 4.02 λ .

Method and Results

To estimate proton beam losses the computer program ACCSIM [4] was used to simulate the injection and accumulation process. ACCSIM uses the CERN library subroutine DISLAN and DINLAN [5]. These subroutines cut off this distributions at λ =-3.5 and λ = 85.96782. Figure 1 shows $\phi(\lambda)$ vs. λ , as given the subroutine DISLAN, and compared with equation 6. It has been shown by Borsch-Supan [6] that for large values of λ the Landau distribution f (λ) =O (1/ λ^2). Using this fact ACCSIM was modified to accommodate the full Landau tail. A test run was made for one billion (10⁹) foil traversals (foil thickness of 2000 µg/cm^{2, which} is ten times thicker than the proposed SNS charge exchange foil). The resulting energy losses were binned and compared with equation 6, ($\lambda_{max} \sim 8300$, ε_{max} =3.3 MeV for a 1 GeV proton) and shown in figure 2.

To estimate particle loss ACCSIM runs (with modified subroutine) were made for 10^5 particles. The first run was made with normal foil size (8 x 4.56 mm²) and thickness (200

 μ g/cm²). Figure 3 shows the ACCSIM output for x-x', y-y', x-y, and Δ E- ϕ phase space after 1225 turn. There was no particle loss and no particles in the beam gap. The average number of the foil traversals was four.

The second run was made with a full aperture foil (100 x 100 mm²) which was ten times thicker (2000 μ g/cm²). Figure 4 shows the ACCSIM output for x-x', y-y', x-y, and Δ E- ϕ phase space after 1225 turn. There were no particle losses but there were 302 particles in the beam gap. These 302 particles are considered as lost particles, as they will lost during the extraction. The average numbers of the foil traversals were 612.

The ratio of average number of foil traversals of run 2 to run 1 is 153, and foil factor due thickness is 10. We therefore calculate the fractional beam loss due to energy straggling for a normal foil as:

Fractional Loss = $\frac{302}{10^5 \times 10 \times 153} = 2 \times 10^{-6}$



Figure 1: Comparison between the subroutine DISLAN (solid line) and Equation 6 (broken line) for the Landau distribution.



Figure 2: Comparison between the test program and Equation 6 for 10^9 foil traversals (thickness 2000 μ g/cm²).



Figure 3: ACCSIM output for 10^5 particles with foil size 8 x 4.58 mm² and thickness 200 μ g/cm².



Dec 21 13:06:27 1998 ACCSIM v3.5j

Figure 4: ACCSIM output for 10^5 particles with size foil 100 x 100 mm² and thickness 2000 µg/cm².

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