

Comparison of intrabeam scattering results for a RHIC lattice and a FODO approximation including coupling effects

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December 2004

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U.S. Department of Energy

USDOE Office of Science (SC)

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C-A/AP/#180
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for a RHIC lattice and a FODO approximation
including coupling effects**

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October 28, 2004

Abstract

This note compares the results for the intrabeam scattering growth rates of a Rhic lattice with the growth rates of a similar fodo lattice. The latter lattice has only identical fodo cells and no matching insertions. This comparison will allow us to use the simpler fodo lattice to find more simple analytical results for the intrabeam scattering growth rates. Approximate analytical results for the fodo lattice have been given in [1].

1 Introduction

This note compares the results for the intrabeam scattering growth rates of a Rhic lattice with the growth rates of a similar fodo lattice. The latter lattice has only identical fodo cells and no matching straight sections. This comparison may allow us to use the simpler fodo lattice to find more simple analytical results for the intrabeam scattering growth rates. Approximate analytical results for the fodo lattice have been given in [1]. The results given below show that the intrabeam scattering growth rates of the fodo lattice can be used to give good approximate results for the intrabeam scattering growth rates of a Rhic lattice for the lattice used in the calculations as presented below..

For the transverse growth rate, G_x , one can assume that the contribution to G_x from the Rhic insertions can be neglected . This assumes that G_x

comes mostly from coupling with G_s through the horizontal dispersion which is small in the insertions. One can then find the Rhic growth rate $G_{x,rhic}$ from the fodo growth rate $G_{x,fodo}$ using

$$G_{x,rhic} = G_{x,fodo} \frac{L_{fodo}}{L_{rhic}} \quad (1)$$

L_{fodo} , and L_{rhic} are the lengths of these two lattices respectively.

For the longitudinal growth, the contribution to G_s from the Rhic insertions is appreciable but not large. One can write

$$G_{s,rhic} = G_{s,fodo} \quad (2)$$

with about a 10% error.

Similar results have been found by A. Fedotov.

2 The fodo lattice.

The fodo lattice used below has 72 identical cells, each cell is the same as the cells of the inner arc of a Rhic lattice. This lattice has a length of 2100 meters compared to the Rhic lattice length of 3830 meters.

3 Computed growth rates for a Rhic lattice and the fodo lattice.

Growth rates were computed using the IBS.P program [2] for a beam of ionized gold atoms with 1 e9 ions/bunch, using a RF cavity with $h=360$ and $V=300$ volts, and with $\gamma=100$. The growth rates G_i are related to the average emittances of the bunch, $\bar{\epsilon}_i$ by

$$G_i = \frac{1}{2} \frac{1}{\bar{\epsilon}_i} \frac{d\bar{\epsilon}_i}{dt} \quad i = x, y, s \quad (3)$$

Because of the factor of 1/2 in Eq.(3), the growth rates given below are the growth rates for σ and not that of the emittance.

The comparison given below for the growth rates for the Rhic and fodo lattices uses the growth rates at $t=0$, before the emittances have grown.

For the Rhic lattice the computed results are

$$\begin{aligned} G_x &= .216/hour \\ G_s &= 4.71/hour \end{aligned} \tag{4}$$

For the fodo lattice the computed results are

$$\begin{aligned} G_x &= .351/hour \\ G_s &= 5.29/hour \end{aligned} \tag{5}$$

$L_{fodo} = 2130m$, and $L_{rhic} = 3834m$. $L_{fodo}/L_{rhic}=.556$. Eq.1 gives $G_{x,rhic}=.195$ which is in error by 10%. Eq. 2 for $G_{s,rhic}$ also is off by 10%.

The result for G_x depends on how one treats the coupling between the x and y motions. A complete treatment of the coupling was given by A. Piwinski [3]. His results are difficult to use as they require knowing all the sources of coupling in the ring such as the skew quadrupoles in each element. What may be a more useable treatment of coupling is the one given in [4]. There it is assumed that the x and y motions are completely coupled, which is interpreted to mean that

$$\langle \epsilon_x \rangle = \langle \epsilon_y \rangle \tag{6}$$

where $\langle \rangle$ indicates an average over all particles in a bunch. ϵ_x and ϵ_y now are not constants of the motion, and the particle distribution is now assumed to be a function of the total emittance, $\epsilon_t = \epsilon_x + \epsilon_y$ which is a constant of the motion to a fairly good approximation. Equations can then be found for $d \langle \epsilon_t \rangle / dt$ and $d \langle \epsilon_s \rangle / dt$ which can be integrated to find $\langle \epsilon_t \rangle$ and $\langle \epsilon_s \rangle$ as a function of time. One result of this approach is that the growth rate for $\langle \epsilon_t \rangle$, G_t , is related to the growth rates of $\langle \epsilon_x \rangle$ and $\langle \epsilon_y \rangle$ as computed in the absence of coupling, and denoted here by G_x and G_y , by

$$G_t = \frac{1}{2}(G_x + G_y) \tag{7}$$

In the operation of Rhic without electron cooling, the coupling is only about a 10% effect on the beam size after about an hour [4]. Although the initial growth rate may be reduced by as much as a factor of two, the beam

size after about an hour is only reduced by about 10%. When electron cooling is present, the possible reduction in the initial growth rate of a factor of 2, may become more important as the electron cooling has to fight the initial growth rate due to intrabeam scattering.

The above results were found using an older Rhic lattice with 6 insertions each with $\beta = 6$ at the crossing points.

Acknowledgements

I thank Jie Wei and Alexei Fedotov for their helpful comments.

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