

## Simple Estimates of Collisional Excitation Cross Sections

M. Blaskiewicz

March 2006

Collider Accelerator Department  
**Brookhaven National Laboratory**

**U.S. Department of Energy**

USDOE Office of Science (SC)

Notice: This technical note has been authored by employees of Brookhaven Science Associates, LLC under Contract No. DE-AC02-98CH10886 with the U.S. Department of Energy. The publisher by accepting the technical note for publication acknowledges that the United States Government retains a non-exclusive, paid-up, irrevocable, world-wide license to publish or reproduce the published form of this technical note, or allow others to do so, for United States Government purposes.

## **DISCLAIMER**

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, nor any of their contractors, subcontractors, or their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or any third party's use or the results of such use of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof or its contractors or subcontractors. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

# **Simple Estimates of Collisional Excitation Cross Sections**

M. Blaskiewicz



**Collider-Accelerator Department  
Brookhaven National Laboratory  
Upton, NY 11973**

# Simple estimates of collisional excitation cross sections

M. Blaskiewicz\*  
BNL, Upton NY 11973, USA

Transverse emittance reduction has been observed in the AGS with helium like gold [1]. Additionally, the extracted emittance is always close to  $10\pi$  mm – mrad regardless of the care with which the beam is accelerated [2]. It has been suggested [3, 4] that the cooling is due to collisionally exciting the electronic levels with the subsequently emitted photon removing energy from the system. Simple estimates for the collisional excitation of the electronic levels and the cooling rates based upon them are presented. I find no evidence supporting the hypothesis.

## I. INTRODUCTION

Typical parameters for an  $\text{Au}^{+77}$  bunch on the AGS injection porch are given in Table I. Since the bunch is nonrelativistic the rms transverse velocity is  $\sigma(v_x) = \omega_0 Q_x \sigma(x) = 4 \times 10^4 \text{m/s}$ , and the rms longitudinal velocity is  $\sigma(v_s) = C f_s \omega_0 \sigma(t) = 2.7 \times 10^5 \text{m/s}$ . With these rms velocities the typical kinetic energy in the comoving frame is  $\delta E = M[\sigma(v_s)^2/2 + \sigma(v_x)^2] = M c^2 \sigma(v/c)^2/2 \approx 80 \text{keV}$ . Cooling the beam requires a mechanism to reduce  $\delta E$ . Inelastic collisions between the ions in the bunch are a possible candidate for this energy loss and we go on to consider this possibility.

## II. ASSUMPTIONS AND KINEMATICS

We will treat the motion of the nuclei classically and the electrons and photons quantum mechanically. Toward this end consider Figure 1. In Fig 1 A, which is the center of mass frame, each ion has a velocity of magnitude  $v_0$  and the total energy is  $E_i = M v_0^2$ . During the collision one of the ions is excited and radiates a photon of energy  $\Delta E$ , and momentum  $\Delta E/c$ . In the final state let the ion velocities be  $\mathbf{v}_1$  and  $\mathbf{v}_2$  and suppose that the photon travels in direction  $\mathbf{n}$ . Conservation of momentum and energy read

$$M(\mathbf{v}_1 + \mathbf{v}_2) + \mathbf{n}\Delta E/c = 0, \quad (1)$$

and

$$M(v_1^2 + v_2^2)/2 + \Delta E = M v_0^2. \quad (2)$$

TABLE I: AGS injection parameters for  $\text{Au}^{+77}$ .

parameter	value
circumference	$C = 807$
Lorentz factor	$\gamma = 1.10$
revolution period	$T_0 = 6.4 \mu\text{s}$
angular revolution frequency	$\omega = 2\pi/T_0 = 9.8 \times 10^5 \text{s}^{-1}$
nominal betatron tunes	$Q_x = Q_y = 8.2$
rms beam radius	$\sigma(x) = \sigma(y) = 5 \text{ mm}$
$h = 24$ RF voltage	$V_{RF} = 280 \text{ kV}$
synchrotron frequency	$f_s = 6.8 \text{ kHz}$
rms bunch length	$\sigma(t) = 50 \text{ ns}$
ions per bunch	$N_b = 5 \times 10^8$
nuclear mass and charge	$M = 185 \text{GeV}/c^2, Z=79$

---

\*Electronic address: [blaskiewicz@bnl.gov](mailto:blaskiewicz@bnl.gov)

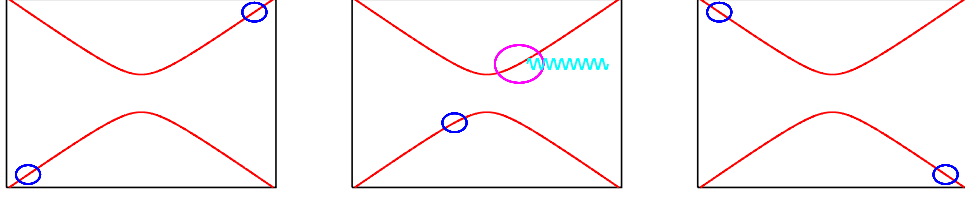


FIG. 1: A schematic of collisional excitation. Initially (A, left) the two ions are in the electronic ground state. As they approach (B, center) one of the ions undergoes an electronic excitation and radiates a photon of energy  $\Delta E$ . Finally (C, right) the ions separate with the total kinetic energy reduced by  $\Delta E$ .

We now show that the photon momentum in Eq (1) can be neglected for  $v_0 \ll c$ . From Eq (1) the maximum difference in the final speeds of the particles occurs when  $\hat{n}$  is parallel to either  $\mathbf{v}_1$  or  $\mathbf{v}_2$ . For definiteness take  $\hat{n}$  parallel to  $\mathbf{v}_2$ . Then the final speeds are related by  $v_2 = v_1 - \Delta E/Mc$ . Substituting this in equation (2) one finds that the maximum speed for a particle in the final state is

$$v_{max} = \sqrt{v_0^2 - \frac{\Delta E}{M} \left[ 1 - \frac{\Delta E}{Mc^2} \right] + \left[ \frac{\Delta E}{2Mc} \right]^2} + \frac{\Delta E}{2Mc}. \quad (3)$$

Additionally, from Eq (1) we have  $|v_1 - v_2| \leq |\mathbf{v}_1 + \mathbf{v}_2| = \Delta E/Mc$  when all photon directions are considered. Therefore

$$\frac{|v_1 - v_2|}{\min(|v_0 - v_1|, |v_0 - v_2|)} \leq \frac{\Delta E/Mc}{v_0 - \sqrt{v_0^2 - \frac{\Delta E}{M} \left[ 1 - \frac{\Delta E}{Mc^2} \right] + \left[ \frac{\Delta E}{2Mc} \right]^2} - \frac{\Delta E}{2Mc}} \leq \frac{v_0/c}{1 - v_0/c}, \quad (4)$$

where the final inequality follows from multiplying top and bottom by  $v_0 + v_{max} - \Delta E/Mc$  and simplifying.

For the problem at hand  $v_0/c \ll 1$  so that we may take  $v_1 = v_2$  and the change in velocity of either ion is  $v_0 - \sqrt{v_0^2 - \Delta E/M}$ . From a physical standpoint one may suppose that the ion velocities change when the ion is excited and that the subsequent photon emission doesn't change the velocity significantly.

While the final state velocities can be calculated, the details of the excitation process need to be treated quantum mechanically so that a classical description of the ion trajectories cannot be obtained. A full blown quantum mechanical description will not be attempted here. Instead, the excitation probabilities will be obtained by assuming that the ions undergo classical Rutherford scattering. This yields a time dependent classical potential which can lead to excitation of the electrons.

To obtain the ion trajectories let  $v_r = 2v_0$  be the relative velocity of the ions before the collision and let  $b$  be the impact parameter. The distance of closest approach between the nuclei satisfies (cgs units,[5] p108)

$$\frac{1}{r_{min}} = \frac{2Z^2e^2}{Mv_r^2b^2} \left[ \sqrt{1 + \left( \frac{Mv_r^2b}{2Z^2e^2} \right)^2} - 1 \right]. \quad (5)$$

The radial acceleration is a maximum at  $r_{min}$  and is given by

$$\ddot{r}_{max} = \frac{2Z^2e^2}{Mr_{min}^2} \sqrt{1 + \left( \frac{Mv_r^2b}{2Z^2e^2} \right)^2}. \quad (6)$$

When  $r = r_{min}$  the direction of the internuclear axis is rotating at its maximum rate and is given by

$$\dot{\theta}_{max} = \frac{bv_r}{r_{min}^2}. \quad (7)$$

For purposes of easy calculation note that

$$\frac{Mv_r^2b}{Z^2e^2} = \left( \frac{v_r}{c} \right)^2 \frac{b}{r_c}$$

where  $r_c$  is the classical particle radius,  $r_c = (Z^2/A)1.54 \times 10^{-18}$  meters which yields  $r_c = 4.63 \times 10^{-17}$  m for helium-like gold.

### III. ELECTRONIC DIPOLE TRANSITIONS

Consider the orders of magnitude in the problem. During the collision the electric field at one nucleus due to the other is  $\mathbf{E}(t) = \hat{n}Ze/r^2(t)$ , where  $\hat{n}$  is a unit vector. For  $v_r = 10^6$  m/s and  $b = 0$ ,  $r_{min} = 0.17 \times 10^{-10}$  m. The electron cloud of helium-like gold has an extent  $\sim a_0/Z$  where  $a_0 = 0.53 \times 10^{-10}$  m is the Bohr radius. The ratio of the extent of the electron cloud to the distance of closest approach is  $\lesssim 0.04$ . Therefore, the electron clouds of the two ions remain well separated and we consider the evolution of each ion independently. We take our coordinate origin as one of the nuclei and neglect the effects of rotation and acceleration of the nucleus. Therefore the interaction hamiltonian is  $H_i(t) = ze^2Z/R^2(t)$  where  $z$  is the electron coordinate along a fixed axis and we have gone over to  $R$  to describe the internuclear distance. The screening effects are small so consider a single electron with hamiltonian  $H = H_0 + H_i$ . Work in the two state approximation with  $H_0|0\rangle = E_0|0\rangle$ ,  $H_0|1\rangle = E_1|1\rangle$ , and interaction energy  $T = \langle 0|H_i|1\rangle$ . Choose a basis so that  $T$  is real and  $E_0 < E_1$ . Next, for a fixed value of  $R$ , define the adiabatic states  $|-\rangle$  and  $|+\rangle$ .

$$|-\rangle = \cos \chi |0\rangle - \sin \chi |1\rangle \quad (8)$$

$$|+\rangle = \sin \chi |0\rangle + \cos \chi |1\rangle \quad (9)$$

where

$$\langle 0|H_0 + H_i|-\rangle = E_- \langle 0|-\rangle = E_0 \cos \chi - T \sin \chi = E_- \cos \chi \quad (10)$$

$$\langle 1|H_0 + H_i|-\rangle = E_- \langle 1|-\rangle = T \cos \chi - E_1 \sin \chi = -E_- \sin \chi \quad (11)$$

and a similar equations hold for  $|+\rangle$ . Solving the algebra yields

$$E_{\pm} = (E_0 + E_1)/2 \pm \sqrt{(E_0 - E_1)^2/4 + T^2} \quad (12)$$

with

$$\tan(2\chi) = 2T/(E_1 - E_0).$$

Define the phases

$$\Phi_{\pm} = \int_0^t E_{\pm}(t_1) dt_1 / \hbar$$

and let the wavefunction describing the system be

$$|\Psi\rangle = C_-(t)|-\rangle \exp[i\Phi_-(t)] + C_+(t)|+\rangle \exp[i\Phi_+(t)]$$

with  $C_-(-\infty) = 1$ . At time  $t$ ,  $|C_+(t)|^2$  is the probability that the electron is in the excited state. Remember that both  $|+\rangle$  and  $|-\rangle$  depend on time, since  $\chi$  does. The equations for  $C_+$  and  $C_-$  follow from

$$\langle + | i\hbar \frac{\partial}{\partial t} - H_0 - H_i | \Psi \rangle = 0,$$

and

$$\langle - | i\hbar \frac{\partial}{\partial t} - H_0 - H_i | \Psi \rangle = 0.$$

After taking the inner products one finds

$$\dot{C}_+ = -\dot{\chi} C_- \exp(i[\Phi_+ - \Phi_-]) \quad (13)$$

$$\dot{C}_- = +\dot{\chi} C_+ \exp(i[\Phi_- - \Phi_+]). \quad (14)$$

By construction  $E_+ > E_-$  so for times well before the collision we may take  $C_- = 1$  and  $C_+ = 0$ . Under the supposition that  $C_+$  remains small compared to 1 a first approximation is given by

$$C_+(t) = - \int_{-\infty}^t dt_1 \dot{\chi}(t_1) \exp(i[\Phi_+(t_1) - \Phi_-(t_1)]). \quad (15)$$

$m_e c^2 - E_{n,j} \text{ keV}$	$n$	$j$
93.46	1	1/2
23.93	2	1/2
21.69	2	3/2

TABLE II: Binding energies for  $Z = 79$ . Both the  $2S_{1/2}$  and  $2P_{1/2}$  states have  $n = 2$ ,  $j = 1/2$ . The energy degeneracy is removed by the Lamb shift.

Using non-relativistic wavefunctions the interaction energy is given by

$$T = \langle 1|H_i|0 \rangle = Ga_0 eE(t)/Z,$$

where  $G$  is a pure number depending on which orbitals are involved. For  $1S \rightarrow 2P$ ,  $G = 8(2/3)^5/\sqrt{2} \approx 0.745$ ; for  $2S \rightarrow 2P$ ,  $G = 3$ .

A photon can be emitted either during the collision or well after the collision. Consider the second case and also assume  $\chi \ll 1$  during the collision so that  $(E_+ - E_-) \approx E_1 - E_0 \equiv \hbar\omega$  and  $\tan(2\chi) \approx 2\chi$ . Then

$$C_+(\infty) \approx - \int_{-\infty}^{\infty} \frac{\langle 1|H_i|0 \rangle}{\hbar} \exp(i\omega t) dt \quad (16)$$

$$= \frac{Gea_0}{\hbar Z} \int_{-\infty}^{\infty} E(t) \exp(i\omega t) dt \quad (17)$$

$$\approx \frac{Gea_0}{\hbar Z} \int_{-\infty}^{\infty} \frac{Ze}{[r_{min} + a_{max}t^2/2]^2} \exp(i\omega t) dt \quad (18)$$

$$\approx \frac{Ge^2a_0}{\hbar r_{min}^2} \int_{-\infty}^{\infty} \frac{1}{1 + a_{max}t^2/r_{min}} \exp(i\omega t) dt \quad (19)$$

$$\approx \pi \frac{Ge^2a_0}{\hbar r_{min}^2} \tau_c \exp(-\omega\tau_c), \quad (20)$$

where  $\tau_c = \sqrt{r_{min}/a_{max}}$  is the characteristic time scale of the collision. For a given  $v_r$  the characteristic time is minimum for  $b = 0$ ,  $\tau_{c,min} = \sqrt{32}(c/v_r)^3 r_c/c$ .

For a hydrogen-like atom with nuclear charge  $Ze$  the energies are [7] p75.

$$E_{n,j} = m_e c^2 \left( 1 + \left[ \frac{Z\alpha}{n - j - 1/2 + \sqrt{(j + 1/2)^2 - Z^2\alpha^2}} \right]^2 \right)^{-1/2} \quad (21)$$

where  $\alpha = e^2/\hbar c = 1/137.036$  is the fine structure constant,  $n = 1, 2, \dots$  is the principle quantum number and  $j = 1/2, 3/2, \dots, n - 1/2$  is to total angular momentum of the state. Table II gives a few values for  $Z = 79$ . For helium-like gold the  $1S \rightarrow 2P$  transition has excitation energy 71.8keV while the  $2S \rightarrow 2P$  transition has excitation energy 2.24keV. Taking  $v_r = 1 \times 10^6 \text{ m/s}$  we find  $\omega\tau_{c,min} = 75$  for the  $2S \rightarrow 2P$  transition. So in this approximation there is essentially no probability of emitting a photon with either helium-like or lithium-like gold ions in the AGS.

Radiative corrections remove the degeneracy between the the  $2S_{1/2}$  and  $2P_{1/2}$  states, which both have have  $n = 2$ ,  $j = 1/2$ . For lithium-like uranium the measured value is  $E_{2S_{1/2}} - E_{2P_{1/2}} = 280.6 \text{ eV}$  [8]. The energy shift is given by [7][p 365]

$$E_{2S_{1/2}} - E_{2P_{1/2}} = \frac{m_e c^2 \alpha^4 Z^5}{6\pi} \left( \ln \left[ \frac{\langle E_{2P} \rangle}{Z^2 \alpha^2 \langle E_{2S} \rangle} \right] + \frac{91}{120} \right), \quad (22)$$

where  $\langle E_{2S} \rangle$  and  $\langle E_{2P} \rangle$  are defined in terms of expectation values for energy shifts due to interactions with the zero-point fluctuations [7][p 364]. We assume that lithium-like gold is sufficiently close to lithium-like uranium so that

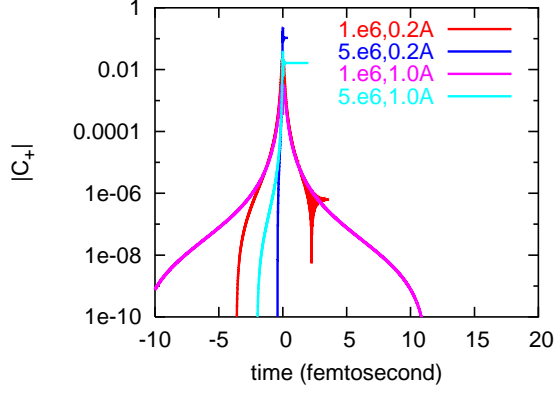


FIG. 2: Numerically obtained plots of  $|C_+|$  for the Lamb transition. Various values of  $v_r$  (m/s) and  $b$  (angstroms) are shown.

only the leading  $Z^4$  scaling is important. This gives  $E_{2S_{1/2}} - E_{2P_{1/2}} \approx -149$  eV. Taking the same relative velocity as before we get  $\omega\tau_c \approx 5$ . However, for these same parameters  $T/(E_1 - E_0) \approx 5.8$  at closest approach so equation (20) is inapplicable. Numerical integration of (15) will be used for this case. Representative results are shown in Figure 2.

Once  $C_+(\infty, b, v_r)$  is obtained the total excitation cross section is given by

$$\sigma(v_r) = 2\pi \int_0^\infty b db |C_+(\infty, b, v_r)|^2, \quad (23)$$

For the three transitions considered the results are very small. For  $v_r = 10^6$  m/s the  $2S_{1/2} \rightarrow 2P_{1/2}$  cross section is  $4 \times 10^{-33} \text{m}^2$  and this rises to  $5 \times 10^{-26} \text{m}^2$  for  $v_r = 2 \times 10^6$  m/s. The cross sections for the other two transitions are even smaller.

It might be possible that the photon is emitted *during* the collision. To estimate the probability of this we use

$$P_\gamma \approx \max_t |C_+(t)|^2 \tau_c / \tau_r,$$

where  $\tau_r$  is the radiative lifetime of the excited state. As a rough approximation to  $\tau_r$  we take the nonrelativistic quantum mechanical expression [9] p179,

$$\frac{1}{\tau_r} = \frac{4\alpha\omega^3 D^2}{3c^2}$$

Where  $D = Ga_0/Z$  is the electric dipole moment of the transition for radiation of frequency  $\omega$ . We estimate  $|C_+(t)| \sim |\dot{\chi}|/\omega$  and obtain cross sections  $\sim 10^{-30} \text{m}^2$  for the relevant parameter ranges.

#### IV. CONCLUSIONS

The possibility of electronic excitation of gold ions in the AGS due to collisions has been considered. We found cross sections  $\sigma \lesssim 10^{-30} \text{m}^2$  for reasonable parameters. For typical particle densities of  $n = 10^{11} \text{m}^{-3}$  and a relative velocity of  $v_r \lesssim 10^6 \text{m/s}$  the mean free time between collisions is  $\tau_{mf} = 1/(n\sigma v_r) \gtrsim 10^{13} \text{sec}$  which is 13 orders of magnitude too long to be of consequence.

#### Acknowledgements

Thanks to Thomas Roser and Dejan Trbojevic for useful discussions and encouragement.

---

[1] Wilem Van Asselt *private communication* circa 1995.



- [2] Leif Ahrens *private communication*.
- [3] Thomas Roser *private communication*.
- [4] K. Kilian, Proceedings of the Workshop on Electron Cooling and New Cooling Techniques, World Scientific, Singapore, New Jersey, London, Hong Kong, p 238, 1990.
- [5] H. Goldstein, *Classical Mechanics* Addison-Wesley, 1981.
- [6] See eg L.I. Schiff *Quantum Mechanics* McGraw-Hill, (1968).
- [7] C. Itzykson and J.B. Zuber *Quantum Field Theory* McGraw-Hill, (1980).
- [8] J. Schweppe, A. Belkacem, L.Blumenfeld, N. Claytor, B. Feinberg, H. Gould, V.E. Kostroun, L.Levy, S. Misawa, J.R. Mowatm H.H. Prior, PRL **66**, 1434-1437, (1991).
- [9] R. Loudon, *The Quantum Theory of Light* Oxford, 1986.