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# Conceptual Design for the Acceleration of Polarized Protons in RHIC

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# Conceptual Design for the Acceleration of Polarized Protons in RHIC <sup>1</sup>

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# 1. INTRODUCTION

The Relativistic Heavy Ion Collider (RHIC) at Brookhaven allows for the unique possibility of colliding  $250\text{ GeV}$  polarized proton beams at luminosities of up to  $2 \times 10^{32}\text{ cm}^{-2}\text{ s}^{-1}$ . This report describes the technical considerations of accelerating polarized proton beams in RHIC to a top energy of  $250\text{ GeV}$  and of performing collisions of longitudinally and transversely polarized protons in the two interaction regions occupied by the PHENIX and STAR detectors.

Fig. 1 shows all the major components that are required for the acceleration of polarized beams to RHIC top energy. The feasibility of accelerating polarized protons in RHIC was the basis of the proposal the RHIC Spin Collaboration (RSC) submitted to the Brookhaven PAC in October 1992[1].

## 2. POLARIZED ION SOURCE, AGS AND TRANSFER OF THE POLARIZED BEAM TO RHIC

To achieve high luminosity, high energy polarized proton collisions in RHIC the intensity of the present AGS polarized proton source[2] is sufficient. The source produces about  $35 \mu A$  of  $H^-$  with 75 – 80% polarization in  $350 \mu s$  pulses at a repetition rate of 5  $Hz$ . The polarized  $H^-$  ions are accelerated to 200  $MeV$  with an RFQ and the 200  $MHz$  LINAC with an efficiency of about 50%. Twenty pulses of  $H^-$  ions are strip-injected, accumulated, and captured into a single bunch in the AGS Booster with an estimated efficiency of about 50%. The bunch in the Booster will then contain  $N_B = 4 \times 10^{11}$  polarized protons with a normalized emittance of about  $\epsilon_N = 10\pi \text{ mm mrad}$ . The single bunch of polarized protons is accelerated in the Booster to 1.5  $GeV$  kinetic energy and then transferred to the AGS, where it is accelerated to 25  $GeV$ .

During acceleration, the polarization may be lost when the spin precession frequency passes through a depolarizing resonance. These resonances occur when the number of spin precession rotations per revolution  $G\gamma$  ( $G = 1.793$  is the anomalous magnetic moment of the proton,  $\gamma = \frac{E}{m}$ ) is equal to an integer (imperfection resonances) or equal to  $kP \pm \nu_y$  (intrinsic resonances). Here  $P = 12$  is the superperiodicity of the AGS,  $\nu_y = 8.8$  is the vertical betatron tune and  $k$  is an integer. The depolarization is caused by the small horizontal magnetic fields present in all ring accelerators which, at the resonance condition, act coherently to move the spin away from the stable vertical direction. Imperfection resonances are due to the horizontal fields caused by the vertical closed orbit errors and intrinsic resonances are caused by the horizontal focusing fields which are sampled due to the vertical betatron motion. The two resonances in the Booster ( $G\gamma = 3$  and  $G\gamma = 4$ ) are too weak to cause any significant depolarization. Traditionally, the depolarizing resonances in the AGS were corrected by the tedious harmonic correction method for the imperfection resonances and the tune jump method for

the intrinsic resonances[3].

In the approved experiment E-880 at the AGS to be started in 1993, the feasibility of polarized proton acceleration by using a 5% partial Siberian Snake[4] will be tested. Previous experience with polarized proton acceleration in the AGS indicates that a 5% Snake is sufficient to avoid depolarization due to the imperfection resonances without using the harmonic correction method. The remaining six important intrinsic resonances can be corrected by the proven tune jump method. At 25 *GeV*, the polarized protons are transferred to RHIC. At this energy the transfer line between the AGS and RHIC is spin transparent[5].

We estimate that the overall efficiency of the acceleration and beam transfer is better than 50%, giving  $2 \times 10^{11}$  protons per bunch. With proper care in the tune jump procedure the normalized emittance of the bunch is expected to be less than  $20\pi$  *mm mrad*. We repeat the process until all 57 bunches of each ring are filled. Since each bunch is accelerated independently, we have the option of preparing the polarization direction of each bunch independently. Filling both RHIC rings with 57 bunches each and acceleration to full energy will only take about 10 minutes which is short compared to the expected lifetime of the stored polarized proton beams in RHIC of many hours.



## 3. SIBERIAN SNAKES AND SPIN ROTATORS IN RHIC

### 3.1. General Lay-out

By inserting two full Siberian Snakes on opposite sides of each of the two RHIC rings, depolarization from imperfection and intrinsic depolarizing resonances can be avoided up to the top energy of 250 GeV. In addition to the Siberian Snakes, spin rotators are required at the intersection points used by PHENIX and STAR to allow for measurements of spin effects with longitudinally polarized protons. The spin rotators rotate the polarization from the vertical direction into the horizontal plane on one side of the interaction region and restore it to the vertical direction on the other side.

The Siberian Snakes introduce a  $180^\circ$  spin rotation without generating a net orbit distortion. The spin rotators placed around the experiments rotate the spin by  $90^\circ$  to provide longitudinal polarization at the interaction region again without generating net orbit distortions. In both cases the spin rotation is accomplished with a sequence of large bore, constant field, superconducting dipole magnets. The large bore is necessary to accommodate the orbit excursions inside the Snake and the spin rotators.

Each Snake rotates the spin by  $180^\circ$  around a horizontal axis and the two axes of the two Snakes of each ring have to be perpendicular on each other. We are planning to use pairs of Siberian Snakes with one Snake rotating the spin around an axis that points  $45^\circ$  to the outside and the other Snake rotating around an axis that points  $45^\circ$  to the inside of the ring. In this case all Snakes can be constructed in same way. Also, the two Snakes of each ring have to be installed on opposite sides of the ring. In fact, the beam direction in one Snake has to be exactly opposite to the beam direction in the other Snake to within  $0.3\text{ mrad}$ .

The following is a summary of the proposal for the locations and construction of the Siberian Snakes and the spin rotators (see Fig. 1):

- Two pairs of full Siberian Snakes, one pair in each ring, are installed at the 4 o'clock and 10 o'clock regions as shown in Fig. 2. These Snakes rotate the spin around axes that point  $45^\circ$  to the inside or outside of the ring. We are planning to install the Siberian Snakes in the  $13\text{ m}$  long, cold straight sections between Q7 and Q8.
- The two pairs of spin rotators, for PHENIX at the 8 o'clock region and STAR at the 6 o'clock region, are installed in the  $40\text{ m}$  long straight sections between Q3 and Q4 on either side of the interaction region. The beam direction in the straight sections is different from the direction in the collision area by  $3.67\text{ mrad}$  which introduces a spin rotation that is larger by a factor of  $G\gamma$ . This means that the spin rotators have to prepare a horizontal polarization direction such that after this spin rotation the desired longitudinal polarization direction is obtained at the interaction point. Fig. 3 shows the size of the maximum orbit excursion inside the spin rotators as a function of beam energy. To limit the orbit excursions in the spin rotators to the same value as in the Snakes, which is about  $70\text{ mm}$ , we are planning to inject with the spin rotators turned off and only turn them on at an energy of  $75\text{ GeV}$  or higher if needed by the experiments. The turn-on can be done adiabatically without loss of polarization.

### 3.2. Siberian Snake and Spin Rotator Design

For both the Siberian Snake and the spin rotators we used the  $200\text{ mm}$  bore, superconducting RHIC DX dipole magnets as a basis for the design. The Snake consists of seven magnets with dimensions listed in Table 1. The parameters are a result of an optimization using a three dimensional orbit and spin tracking program that includes the effects of fringe fields. The maximum beam excursion at the injection energy is  $70\text{ mm}$ . The result of the orbit and spin tracking is shown in Fig. 4. The horizontal and vertical dipole magnets are powered by separate power supplies. This allows for an adjustment of the spin tune to  $1/2$  and also for small changes of the direction of the rotation axis to compensate for the effect of the detector solenoids as will be discussed below. This is shown in Fig. 5

The spin rotators are constructed of 3 vertical and 3 horizontal dipole magnets. With independent power supplies for the vertical and horizontal dipoles, respectively, it is possible to rotate the polarization into any direction of the horizontal plane according to the requirements outlined in the previous section. To ensure

Deflection direction	H	V	-H	-V	H	V	-H
Dipole Magnetic Field [Tesla]	3.742	4.060	3.742	4.060	3.742	4.060	3.742
Magnet Start [m]	0.000	0.599	1.600	2.619	4.340	5.359	6.962
Magnet End [m]	0.218	1.219	2.238	3.959	4.978	5.979	7.180

Table 3.1: Parameters of the Siberian Snake magnets

Deflection direction	V	H	-2V	-2H	V	H
Magnet Start [m]	0.000	1.789	3.295	4.388	6.826	8.004
Magnet End [m]	0.278	2.704	3.960	6.321	7.104	8.920
Exitation for radial polarization:						
Dipole Magnetic Field [Tesla]	3.588	1.348	3.588	1.348	3.588	1.348
Exitation for long. polarization:						
Dipole Magnetic Field [Tesla]	3.588	4.040	3.588	4.040	3.588	4.040

Table 3.2: Parameters of the Spin Rotator magnets

that the spin lies in the horizontal plane the spin rotation angles  $\psi_H$  and  $\psi_V$  in the horizontal and vertical dipoles, respectively, have to obey the equation[8]

$$\sin \psi_H \sin \psi_V = \pm \frac{1}{2}.$$

The relative sign of  $\psi_H$  and  $\psi_V$  for different directions in the horizontal plane are shown in Fig. 6. Fig. 7 shows the result of the tracking calculations for the spin rotator that produces longitudinal polarization (top) and radial polarization (bottom).

### 3.3. Effect on Lattice

Calculations of the changes to the RHIC lattice parameters were performed using MAD. It is worth noting that the survey algorithm in MAD does not work properly since the Snakes introduce torsion into coordinate system. The following changes in the lattice parameters are caused by the introduction of two Snakes, each located between Q7 and Q8 as described above:

Comment	$v_x$	$v_y$	$D_x^{\max}$	$D_y^{\max}$	$\beta_x^*$	$\beta_y^*$
No Snakes	28.827	28.823	1.674 m	0	2.0 m	2.0 m
With Snakes, Injection energy	28.837	28.836	1.738 m	0.110 m	2.0 m	2.1 m
With Snakes, Top energy	28.827	28.823	1.678 m	0.009 m	2.0 m	2.0 m

Clearly the changes of the lattice parameters are very small and can easily be corrected for.

### 3.4. Compensation for Detector Solenoids

The STAR and PHENIX detectors use solenoid magnets as spectrometers. With transverse polarization at the collision point the solenoid contributes to the imperfection resonance strength,

$$\epsilon_{imp.sol} = \frac{1 + G}{2\pi} \frac{\int B_{\parallel} d\ell}{B\rho}$$

For a 5 Tesla-meters integrated solenoid field strength, the resulting spin resonance strength is about 0.02 at the injection energy and 0.003 at 250 GeV/c.

With longitudinal polarization at the collision point the longitudinal field rotates the polarization around its axis and thus changes the spin tune. The spin tune is changed by 0.03 to 0.003 at 25 GeV and 250 GeV, respectively, by a 5 Tm solenoid. This can be compensated by adjusting the direction of the axes around which the Snakes rotate the spin. By adjusting up to  $\pm 2.5^\circ$  the spin tune can be adjusted for energies down to 50 GeV.

## 4. ACCELERATION OF POLARIZED PROTONS IN RHIC

### 4.1. Depolarizing Resonance Strengths

Without Siberian Snakes there are numerous depolarizing resonances in RHIC, both intrinsic and imperfection resonances. The strengths of the intrinsic resonances can be calculated quite accurately from the appropriate integral over the horizontal focusing fields. Fig. 8 shows the result for the RHIC92 lattice with  $\beta^* = 10 m$  at all intersections. A calculation with  $\beta^* = 1 m$  gave only a slightly different result. The calculation was performed for a particle with a normalized Courant-Snyder invariant of  $\epsilon_0 = 10 \pi mm mrad$ . For a different value of the invariant the strength scales according to

$$\epsilon = \epsilon_0 \sqrt{\frac{\epsilon}{\epsilon_0}}$$

where  $\epsilon_0$  is the resonance strength for the invariant  $\epsilon_0$ .

Important intrinsic spin resonances are located at

$$G\gamma = kP \pm \nu_y \approx mPM \pm \nu_B, \quad (4.1)$$

where  $k$  and  $m$  are integers,  $P$  is the superperiodicity of the accelerator,  $M$  is the number of FODO cells per superperiod, and  $2\pi\nu_B = 2\pi(\nu_y - 6)$  is the accumulated phase advance of all FODO cells, which contain bending dipoles. The locations of the 3 strongest intrinsic resonances are

$$\begin{array}{lll} G\gamma = 3 \times 81 + (\nu_y - 6), & 5 \times 81 - (\nu_y - 6), & 5 \times 81 + (\nu_y - 6) \\ E = & 139 GeV, & 200 GeV, & 224 GeV \end{array}$$

where 81 is the product of superperiodicity, 3, and the “effective” FODO cells per superperiod, 27, which includes dispersion suppressors. The strengths of all 3 strong resonances are less than 0.5.

Important imperfection resonances are located at an integer closest to strong intrinsic resonances. This is clearly shown in the top part of Fig. 9 which shows the calculated imperfection resonance strengths for an uncorrected closed orbit obtained from a random sample of magnet misalignments with a rms spread of  $\pm 0.5 \text{ mm}$ , dipole roll angles with a spread of  $\pm 1 \text{ mrad}$ , dipole field errors of  $\pm 5 \times 10^{-4}$ , and position monitor errors of  $\pm 0.5 \text{ mm}$ . After the closed orbit correction scheme MICADO was applied the strengths are greatly reduced as shown in the lower part for Fig. 9. The strengths of the imperfection resonances generally increase linearly with the beam energy and are bounded by

$$\epsilon_{imp} = 0.25 \frac{\gamma}{250} \sigma_y$$

where  $\sigma_y$  is the rms value of the residual closed orbit excursions. The strength is smaller than 0.04 for all energies.

## 4.2. Effect of Siberian Snakes

With the installation of Siberian Snakes, which are local  $180^\circ$  spin rotators, the spin tune becomes  $1/2$ , independent of the beam energy. Clearly the depolarizing resonance conditions cannot be met anymore as long as the fractional betatron tune  $\Delta\nu_y \neq 1/2$  and therefore, in principle, no depolarization would occur. This is in fact true as long as the depolarizing resonances are not too strong. However, in the presence of strong resonances depolarization can occur from resonance conditions extended over more than just one turn. This leads to additional possible depolarizing resonance conditions:

$$\Delta\nu_y = \frac{\nu_{sp} \pm k}{n}$$

They are called Snake resonances[6] and  $n$ , the number of turns, is called the Snake resonance order. For two Snakes, as proposed here for RHIC, significant depolarization from Snake resonances only occurs for an intrinsic resonance strength of about 0.5 and even order Snake resonances require in addition an imperfection resonance strength of about 0.05. Fig. 10 shows the result of tracking through a region with an intrinsic resonance of strength 0.5 and an imperfection resonance of strength 0.05. There are clearly regions of the betatron tune that do not experience any depolarization. Since the betatron tunes of RHIC were chosen to be located between  $4/5$  and  $5/6$ , the betatron tune could be fit between the Snake

resonances  $13/16 = 0.8125$  and  $5/6 = 0.8333$ . With the betatron tune including its spread located between 0.815 and 0.830, a 0.015 range, no depolarization will occur over the whole RHIC energy range up to the top energy.

If the betatron tune spread is too large to fit into this range, some depolarization will be caused by the Snake resonance  $\Delta\nu_y = 13/16$ . Tracking calculations performed with an acceleration rate of  $\dot{\gamma} = 3.9$  /sec showed that a gaussian beam with  $\epsilon_{N,95\%} = 20\pi$  mm mrad and with 10 % of the beam overlapping the 13/16 Snake resonance less than 5 % of the polarization is lost for each passage through one of the strong intrinsic resonances. The final polarization after passing all 3 strong resonances would then be at least 86 % of the injected polarization. It is important to note that, unlike for electron beams, the betatron tune distribution within the proton beam is basically static and does not get mixed. This means that depolarization experienced by a part of the beam is confined to this part only and will not affect the whole beam. In other words, there is no diffusion of polarization.

### 4.3. Sextupole Depolarizing Resonances

Spin resonances arising from sextupoles are located at

$$\nu_{sp} = kP \pm \nu_x \pm \nu_y = mPM \pm (\nu_x - 6) \pm (\nu_y - 6)$$

with resonance strength given by,

$$\epsilon_K \approx \frac{1 + G\gamma}{8\pi} \sqrt{\frac{\epsilon_x \epsilon_y}{\pi^2}} \sqrt{\beta_x \beta_y} PM (|S_F| + |S_D|),$$

where  $S_F, S_D$  are respectively strengths of sextupoles located at the focusing and defocusing quadrupole locations. Because the emittance decreases with energy, the sextupole spin resonance strength is energy independent in hadron storage rings. For RHIC, the resonance strength is about 0.0003 at a normalized emittance of  $20\pi$  mm mrad. Such a sextupole driven resonance has been observed in the IUCF Cooler Ring. However, because of their small spin resonance strength, depolarizing resonances due to sextupoles are not important as long as the betatron tunes are chosen to avoid the resonance condition, which for a spin tune of  $1/2$  is

$$\frac{1}{2} = kP \pm \nu_x \pm \nu_y$$

The current working point for RHIC certainly avoids this condition.

#### 4.4. Spin Tune Spread and Modulation

With snakes, the spin tune is independent of energy. Therefore the synchrotron motion does not give rise to spin tune spread. This has been verified indirectly in the snake experiment at the IUCF Cooler Ring, where one finds that there is no depolarization at the synchrotron side band for a 100% snake. However the spin tune modulation may still arise from imperfect spin rotation in the snake, and imperfect orbital angle between snakes. The errors in orbital angle may arise from survey error, closed orbit error, and/or betatron motion.

For an imperfection resonance strength  $\epsilon = 0.05$  with two snakes, the perturbed spin tune shift is given by,  $\Delta\nu_{sp} = \frac{\pi|\epsilon|^2}{4} = 0.002$ . The error in the spin rotation angle of two snakes contributes also to the imperfection resonance strength. Assuming that the relative error of the integrated field strength of each snake dipole is  $10^{-3}$ , the error in the spin rotation angle of a snake should be about  $\sqrt{8} \times 180 \times 10^{-3} = 0.5^\circ$ , where we have assumed that a snake is constructed from 8 dipoles. The effect of 2 snakes in the accelerator will give a resonance strength of the order  $\epsilon_{imp}^{eq} \approx 0.004$ , which is smaller than the imperfection resonances due to closed orbit errors.

The error  $\Delta\theta$  in the orbital angle between the two snakes can give rise to a spin tune shift of  $\Delta\nu_{sp} = \frac{1}{\pi}G\gamma\Delta\theta$ . Since the error in the orbital angle gives rise to spin tune shift and not a spin tune spread, one can compensate the effect by adjusting the spin rotation axes of the snakes. A survey error of about  $\Delta\theta \approx 0.1 \text{ mrad}$  leads to a spin tune shift of 0.01 at the highest energy. For such a survey error, active compensation by adjusting the snake spin rotation axes is needed but is well within the tuning range of the Snake design.

The closed orbit can also cause orbital angle error between snakes. Let us assume that the maximum closed orbit is about  $\hat{a} \approx 6\sigma \approx 1.0 \text{ mm}$ . The angular deviation is of the order of  $\Delta x'_{co} \approx \frac{\hat{a}}{\sqrt{\beta\hat{\beta}}}$ , where  $\hat{a}$  is the maximum orbit error

and  $\sqrt{\beta\hat{\beta}} \approx \frac{R}{\nu}$  is the average betatron amplitude function. The expected error is about  $\Delta x'_{co} \approx 2 \times 10^{-5}$  for RHIC, which gives rise to a spin tune shift of 0.002.

Similarly the betatron oscillation can cause orbital angle modulations. The spin tune modulation is given by

$$\Delta\nu_{s,\beta} \approx \frac{1}{\pi}G\gamma\sqrt{\frac{\epsilon}{\beta}} = \frac{1}{\pi}G\sqrt{\frac{\gamma\epsilon_N}{\beta}}$$

The resulting spin tune spread is about 0.007 for a beam with  $20\pi$  mm-mrad normalized emittance at 250 GeV.



Combining all the possible sources, we expect that the total spin tune *spread* to be about 0.009 (imperfection resonance and betatron motion) and a correctable spin tune *shift* of 0.012 (Snake survey error and closed orbit error).

#### 4.5. Betatron Tune Spreads and Modulations

In avoiding snake resonances up to the 13th order, the available tune space is about 0.015. With a spin tune spread of 0.009, the betatron tune needs to be controlled to better than 0.006. The tight requirements for the spin and betatron tune spread are only relevant while accelerating through the 3 strong intrinsic resonances when the beam-beam tune shift is negligible.

At the injection energy, the space charge tune spread can be as large as 0.02 for RHIC. However, the corresponding spin resonance strength at low energy is also about a factor of 3 smaller and therefore the available tune space is much larger.

## 5. STORAGE OF POLARIZED PROTONS IN RHIC

In storage mode even a very small depolarizing resonance strength can in principle lead to significant depolarization. This was observed at the ZGS where the effect of high order depolarizing resonances were studied on a 1 second flat top as shown in Fig. 11[10]. In an accelerator without Snakes like the ZGS the spin tune is energy dependent ( $\nu_{sp} = G\gamma$ ) and therefore if the resonance condition is within the energy spread of the beam each proton will cross the resonance condition repeatedly and eventually the whole beam can be depolarized. With Snakes, however, the spin tune is energy independent and therefore all Snake resonance conditions are energy independent. As pointed out earlier and shown in Fig. 10 the beam can overlap higher order Snake resonances due to the betatron tune spread. Results of spin tracking calculations for a proton with  $\Delta\nu_y = \frac{13}{16}$  and a Snake resonance strength of 0.15 showed that the spin vector is precessing about the vertical axis. At a resonance strength of 0.15, the vertical projection is about 90%. This precession is in fact a spin closed orbit in a broader sense of 16 turns. Spin tracking calculations were calculated over  $8 \times 10^9$  turns without significant deviations from this spin closed orbit. Only a small fraction of particles are located within the width of the resonance and therefore effective depolarization in the storage mode is small.

When the spin vector is acted on with an adiabatic modulation within the tolerable limit, the spin vector will follow the spin closed orbit adiabatically. Non-adiabatic processes, arising from rf noise at the spin precession frequency, can indeed cause beam depolarization. Let us consider that a single dipole with strength  $\theta_k$  is modulated at  $\nu_{sp}f_0$ , which is about 39 kHz for RHIC. The corresponding induced spin precessing kick is  $G\gamma\theta_k$ . The number of turns that occur before the spin is perturbed to 80% of the original polarization is given by,

$$N_p = \frac{\arccos[0.8]}{G\gamma\theta_k}$$

Let us now consider the same angular kick to the orbital motion. If there is an rf source at  $\nu_{sp}f_0$ , one expects a similar angular kick at the frequency  $qf_0$ , where  $q$  is the fractional part of the betatron tune. The orbital survival turn number is given by

$$N_o = \frac{A}{\langle\beta\rangle\theta_k}$$

where  $\langle\beta\rangle$  is the average betatron amplitude, and  $A$  is the dynamic aperture. Using  $A = 0.01$  m,  $\langle\beta\rangle = 20$  m for RHIC, we find that the orbital lifetime is only half as long as the polarization lifetime.

Indeed, any rf source at high frequencies around the synchrotron and betatron tunes, are dangerous to the orbital stability of particles in accelerators. Similarly, any rf source at the spin tune can cause beam depolarization. These high frequency rf sources should be addressed carefully in hardware design.

## 6. MEASURING THE BEAM POLARIZATION IN RHIC

We propose to construct two polarimeters that are capable of measuring the polarization of the circulating beam in each ring independently at various stages in the acceleration cycles.

These polarimeters utilize the asymmetries ( $A_n$ ) in inclusive pion production at high  $X$  that were measured at the Argonne ZGS [11] and at Fermilab[12]. The asymmetries increase linearly with  $X$  and appear to be independent of the incident polarized beam momentum. Moreover, the ZGS data which were taken with both liquid hydrogen and deuterium targets show no dependence on target nuclei.

The polarimeters are designed to probe the kinematic region of 0.5 in  $X$  and  $0.8\text{ GeV}/c$  in transverse momentum of the scattered pions. This was optimized from the Fermilab data based on the fact that the asymmetry rises linearly with  $X$  while the cross section is falling as  $(1 - X)^2$ , and that the error in the asymmetry measurement is the inverse of the square root of the total number of events.

We propose to use  $\pi^-$  asymmetry, since  $\pi^-$  are a relatively large fraction of negatives, and particle identification is therefore not necessary. At the chosen parameters, the measured  $\pi^-$  asymmetry is 18% and the invariant cross section is about  $100 \frac{\mu b}{\text{GeV}^2}$  [13].

The polarimeters are designed to fit in the 40 meter long straight sections between Q3 and Q4. The apparatus consists of a  $5\ \mu\text{m}$  carbon fiber target, followed by a 2 meter long dipole magnet with a  $1\text{ GeV}/c$  transverse momentum kick, and then a hadron calorimeter. The distance from the fiber target to the calorimeter would be about 40 meters, with the spectrometer magnet at 30 meters. The scattering angle and selected spectrometer momentum depend on the beam energy: at RHIC injection of  $25\text{ GeV}/c$ , the scattering angle would be  $64\text{ mrad}$  and the  $\pi^-$  momentum  $12.5\text{ GeV}/c$ ; at  $250\text{ GeV}/c$  RHIC beam momentum, the scattering angle is  $6.4\text{ mrad}$  and the  $\pi^-$  momentum is  $125\text{ GeV}/c$ . The smallest angle,  $6.4\text{ mrad}$ , sets the center of the magnet's aperture at  $20\text{ cm}$  from beam line

center. This places some constraints on the design of the magnet. Upstream of this magnet we envisage using collimation to cut down on background. The magnet kick and the positioning of the downstream detection elements will minimize the contamination from straight through neutrals. The detection apparatus will consist of several trigger counters followed by a hadron calorimeter with energy resolution  $\frac{\Delta E}{E} = 40\%/\sqrt{E}$ , with  $E$  in GeV. This provides an 11% measurement at 12.5 GeV/c. The whole assembly will rotate in order to accommodate the changes in the scattered angle as the circulating beam energy increases. We also envision that the assembly will translate together with the fiber target when the target is moved into the beam, maintaining a fixed geometry independent of the transverse position of the target.

A 5  $\mu\text{m}$  carbon fiber target and a detector with a solid angle of  $2\ \mu\text{sr}$  that is fixed for all energies will result in the following negative pion rates:

Energy	$\pi^-$ rate [ $\text{s}^{-1}$ ]
250 GeV	14000
25 GeV	140

Thus a 1% measurement of the pion asymmetry ( $10^4\ \pi^-$ , a 5% measurement of the polarization) will require less than 2 minutes at the worst case of 25 GeV/c.

The fraction of the beam lost on such a target per second is estimated to be  $2.5 \times 10^{-4} \sqrt{\frac{P_{\text{beam}}}{25\ \text{GeV}/c}}$ , with 1.8% of the beam interacting for the 25 GeV/c measurement. The solid angle was chosen to be compatible with the small angle high energy measurement - we used  $\Delta\theta_x = 1\ \text{mrad}$  and  $\Delta\theta_y = 2\ \text{mrad}$ ; this is an aperture of  $3\ \text{cm} \times 6\ \text{cm}$  at 30 meters from the target. The target thickness was chosen based on consideration of the time required to measure the polarization (to collect  $10^4\ \pi^-$ ) and rates in the polarimeter. The fraction of beam loss which is required to collect  $10^4\ \pi^-$  is independent of the target thickness, but can be adjusted by changing the acceptance. This exercise is meant to show the viability of such a system with refinements to follow.

For the 5  $\mu\text{m}$  carbon fiber, the equilibrium temperatures were calculated to be  $614^\circ\ \text{C}$  and  $1225^\circ\ \text{C}$ , well below the melting temperature for carbon of  $3550^\circ\ \text{C}$  for beam momenta of 25 and 250 GeV/c. For a 1  $\mu\text{m}$  fiber,  $320^\circ\ \text{C}$  and  $730^\circ\ \text{C}$  were calculated.

## 7. LUMINOSITY OF POLARIZED PROTON COLLISIONS

The luminosity at the RHIC beam energy  $E[GeV]$  is given by[14]

$$\mathcal{L} = 5.7 \times 10^{31} \frac{\left(\frac{N_B}{10^{11}}\right)^2 \left(\frac{B}{57}\right) \left(\frac{E}{250}\right)}{\left(\frac{\epsilon_N}{20\pi}\right)^2} \text{ cm}^{-2}\text{s}^{-1}$$

where  $B$  is the number of bunches in each ring,  $N_B$  is the number of particles per bunch, and  $\epsilon_N$  is the normalized emittance. Thus the polarized proton luminosity will be  $2 \times 10^{32} \text{ cm}^{-2}\text{s}^{-1}$ , which corresponds to two interactions per crossing for the  $B = 57$  bunches per ring. Note the importance of smaller emittance in obtaining higher luminosity. The luminosity will gain a factor of four when the emittance is smaller by a factor of two. This is due to the fact that the betatron amplitude function  $\beta^*$  at the interaction point can be squeezed by a factor of two at the smaller emittance while maintaining the same dynamic aperture at the high betatron amplitude quadrupole triplets. This is also the reason behind the linear dependence of the luminosity on energy, since the unnormalized emittance depends on  $E^{-\frac{1}{2}}$ .

We show in Fig. 12 the luminosity vs. energy as described above ("enhanced luminosity"), emphasizing the linear dependency of the luminosity with energy. The RHIC design luminosity for proton-proton is also shown which, for the emittance used above, would represent  $5 \times 10^{10}$  protons per bunch.

## 8. SPIN REVERSAL OF THE STORED BEAMS

Since the proposed asymmetry measurements are high precision measurements, frequent polarization sign reversal is imperative to avoid systematic errors. Possible sources for systematic errors are luminosity variations, crossing angle variations, and detector efficiency variations. As mentioned earlier different bunches will have different polarization sign and therefore different bunch crossings will measure interactions with different combinations of incoming beam polarization signs. Although this will greatly reduce systematic errors it is still true that one pair of bunches would always cross with the same combination of polarization signs during the whole lifetime of the stored beams which is at least several hours. To eliminate the possibility of systematic errors from this situation we propose to install a spin flipper in each ring which is capable of reversing the polarization sign of all bunches. The spin flipper consists of three horizontal DC dipole magnets interleaved with three laminated vertical dipole magnets[8]. Exciting the vertical magnets with about  $40\text{ kHz}$  AC current would drive an artificial spin resonance which can be used to adiabatically reverse the polarization direction. Fig. 13 shows the result of a test of this concept performed at the Indiana University Cyclotron Facility (IUCF)[9]. We estimate that complete spin reversal would take less than 1 second.

The same device will be used to accurately measure the spin tune. This is instrumental to adjust the spin tune to 0.500.

## 9. OPERATIONAL MODES OF RHIC FOR POLARIZED PROTON RUNNING

For normal operation no special operational modes of RHIC are required for the acceleration and storage of polarized protons. Since depolarizing resonances are driven predominantly by vertical orbit excursions, particular care has to be given to the corrected vertical closed orbit and the vertical beam emittance. This means that beam blow-up from beam-beam interactions and stop-bands should be minimized. To avoid depolarization from snake resonances the fractional vertical betatron tune including tune spread has to be kept well within  $\frac{13}{16} = 0.8125$  and  $\frac{5}{6} = 0.8333$ , especially at the energies of the 3 strongest intrinsic resonances:  $G\gamma = 3 \times 81 + (\nu_y - 6)[E = 139\text{GeV}]$ ,  $5 \times 81 - (\nu_y - 6)[E = 200\text{GeV}]$ ,  $5 \times 81 + (\nu_y - 6)[E = 224\text{GeV}]$ . Also the acceleration rate would have to be at least  $\dot{\gamma} = 4.2\text{ s}^{-1}$ , which corresponds to  $\frac{dB}{dt} = 0.05\text{ T/s}$ .

During commissioning, acceleration cycles that allow for polarization measurements at various energies are required, in particular at the injection energy, at energies below and above the 3 strongest intrinsic resonances. and at some intermediate energies:  $25\text{ GeV}$ ,  $50\text{ GeV}$ ,  $100\text{ GeV}$ .



## 10. CONCLUSIONS

Polarized Proton acceleration and storage in RHIC have been carefully examined. We found that

1. The current polarized ion source at the AGS is capable of providing a sufficient number of polarized protons for RHIC
2. The partial Snake experiment at the AGS, E880, will provide the bunch beam intensity of  $N \geq 2 \times 10^{11}$  and  $\varepsilon \leq 20\pi mm mrad$  at 25 GeV.
3. Two Snakes per ring in RHIC will allow for the acceleration of polarized beams.

Relevant issues such as spin tune and betatron tune spread, the Snake and spin rotator design and their effect on the lattice functions, polarimeters in RHIC, and the polarization lifetime have been carefully addressed. It is indeed technically feasible to achieve polarized proton beam collisions in RHIC at a luminosity of  $2 \times 10^{32} cm^{-2} s^{-1}$ .

## 11. FIGURE CAPTIONS

1. Layout of the AGS-RHIC accelerator complex showing all the major component that are necessary for the acceleration of polarized proton beams to the RHIC top energy.
2. View of RHIC overemphasizing the interaction regions to show the location of the Siberian Snakes and the spin rotators placed around the collider experiments STAR and PHENIX. Also shown are the polarization directions around the rings and around the detectors for collisions with longitudinal polarization.
3. Maximum beam orbit excursions inside the spin rotators for longitudinal polarization at the collision point shown as a function of beam energy.
4. Orbit and spin tracking through the seven magnet Siberian Snake at  $\gamma = 200$ . The spin tracking shows the reversal of the vertical polarization.
5. Change of the direction of the Snake rotation axes as function of the excitation of the horizontal and vertical dipole strings.
6. Top view of the horizontal polarization immediately after the spin rotator which rotates the vertical polarization into the horizontal plane. The sign of the horizontal ( $\psi_H$ ) and vertical ( $\psi_V$ ) bending magnets is shown in order to obtain polarization in the indicated quarter of the plane.
7. Orbit tracking through the spin rotator. The top figure shows the tracking for the excitation that produces longitudinal polarization, the bottom figure is for radial polarization.
8. Strengths of the intrinsic depolarizing resonances in RHIC calculated for the RHIC92 lattice and for both  $\beta^* = 10 m$  and  $\beta^* = 1 m$  at all six intersection points. There is no noticeable difference of the calculated strengths for the two values of  $\beta^*$ .

9. Strengths of the imperfection depolarizing resonances in RHIC calculated after simulating the MICADO orbit correction scheme.
10. Vertical component of the polarization after acceleration through a strong intrinsic resonance and a moderate imperfection resonance shown as a function of the fractional vertical betatron tune.
11. Depolarization on the 1 second flat top at the ZGS.
12. Expected luminosity (enhanced luminosity) as a function of energy for a normalized 95 % emittance of  $20 \pi \text{ mm rad}$  and  $\beta^* = 2 \text{ m}$ . The RHIC design luminosity for proton collision is also shown.
13. Adiabatic spin reversal of a stored beam of polarized protons at the IUCF Cooler ring. The RF voltage is proportional to the strength of the artificial spin resonance driving the spin reversal.

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# POLARIZED PROTON COLLISIONS AT BROOKHAVEN

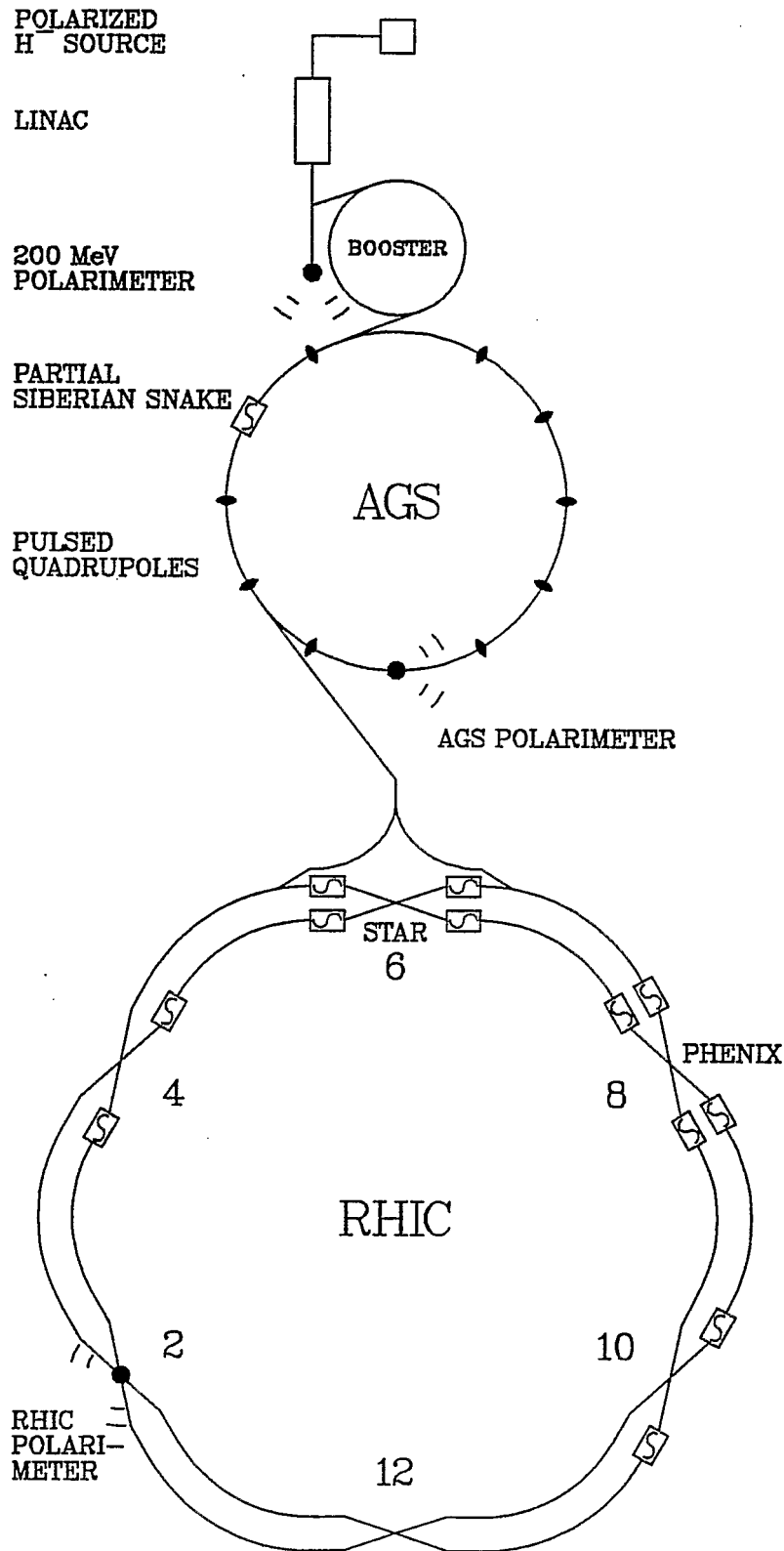


Figure 1

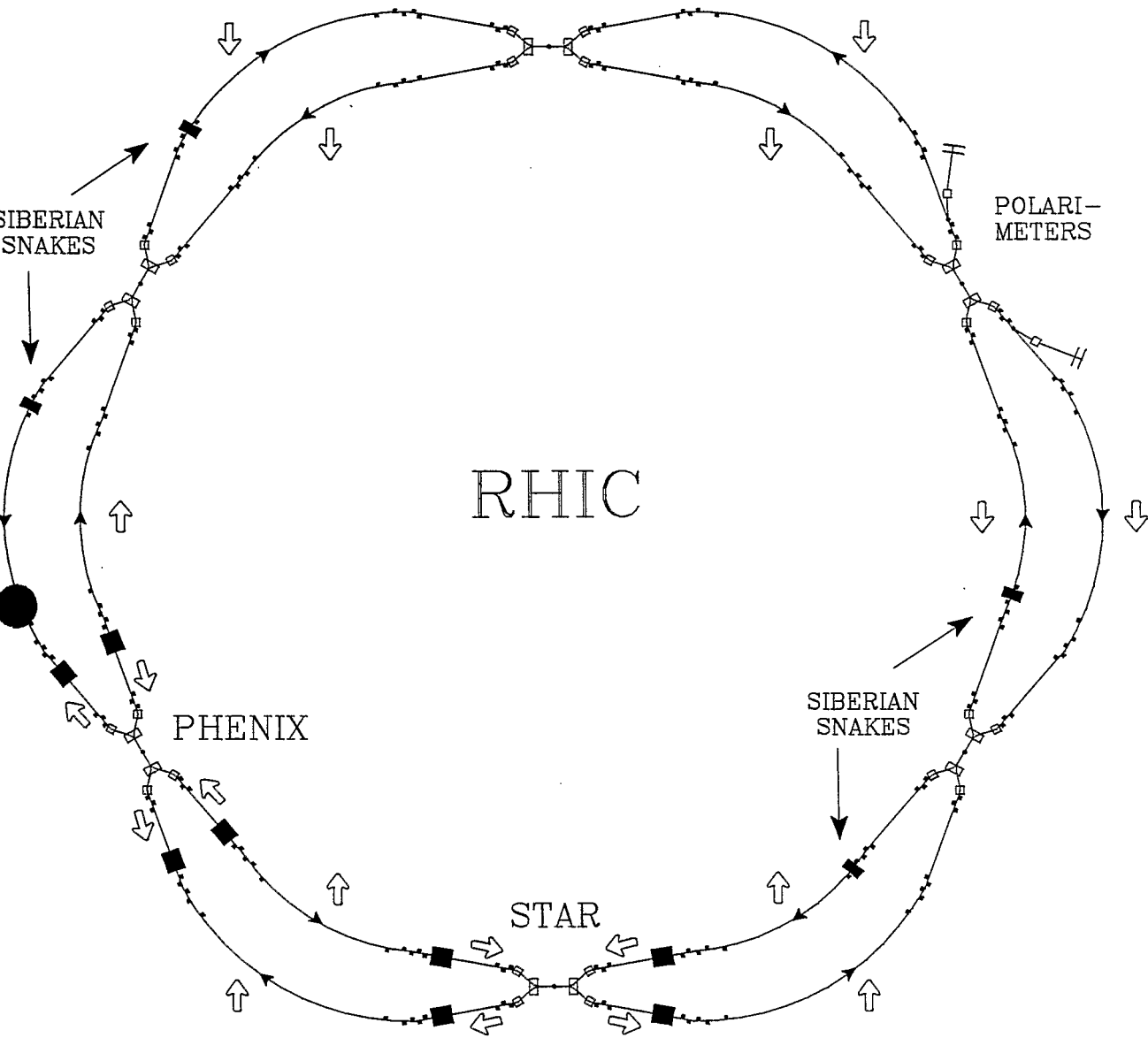
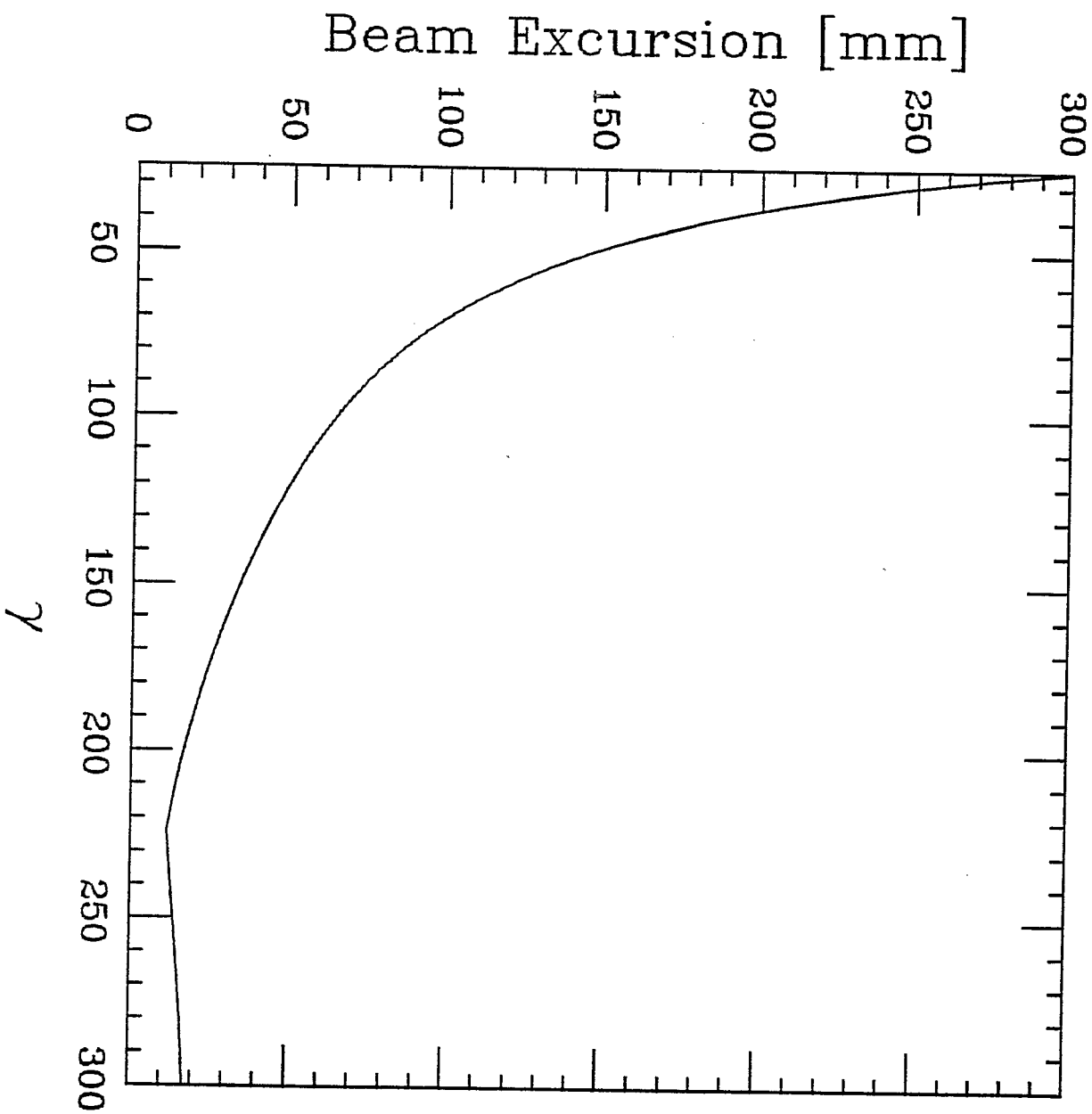


Figure 2

Figure 3



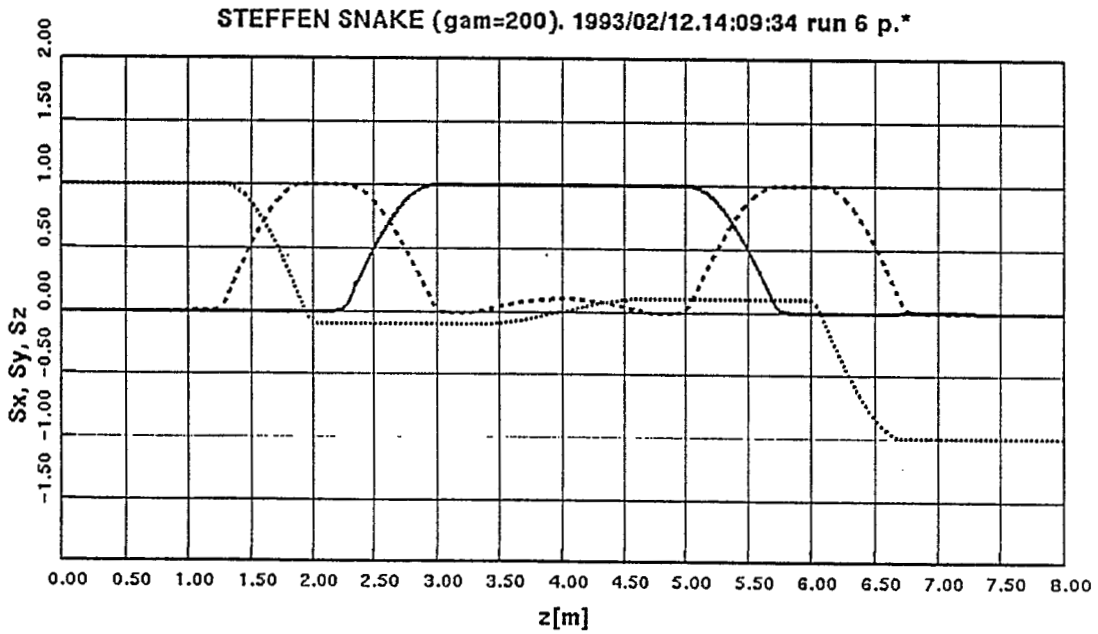
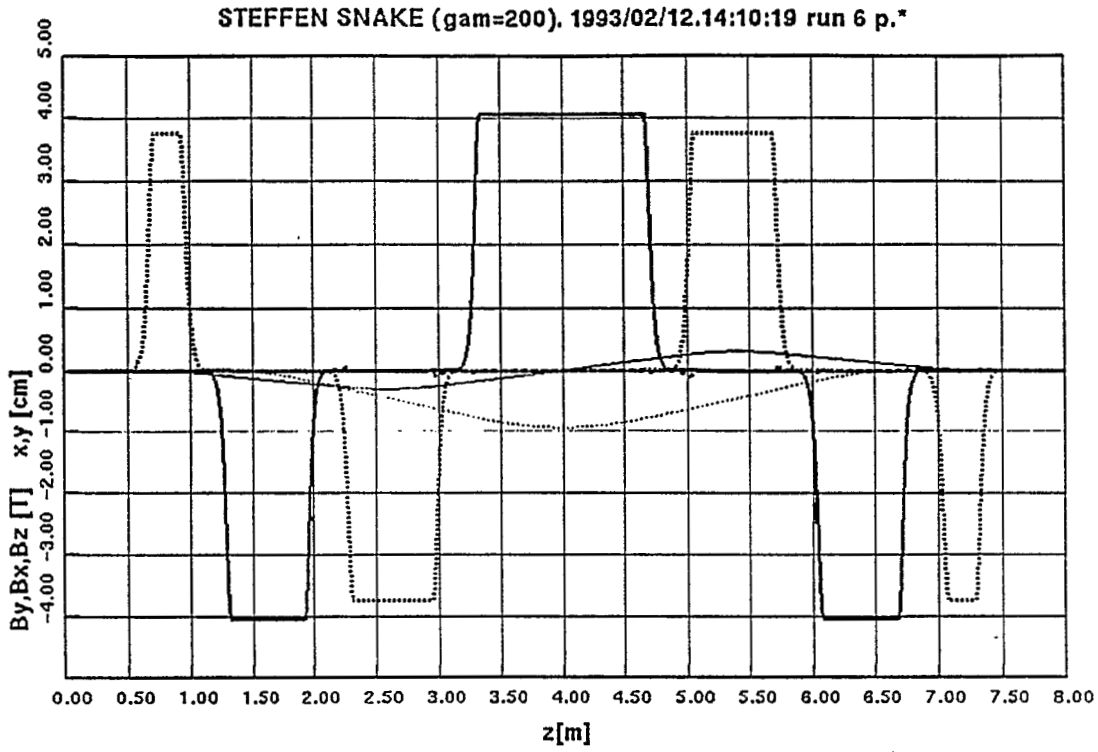


Figure 4



STEFFEN SNAKE (gam=200). 1993/02/13.12:22:51 run 6 p.\*

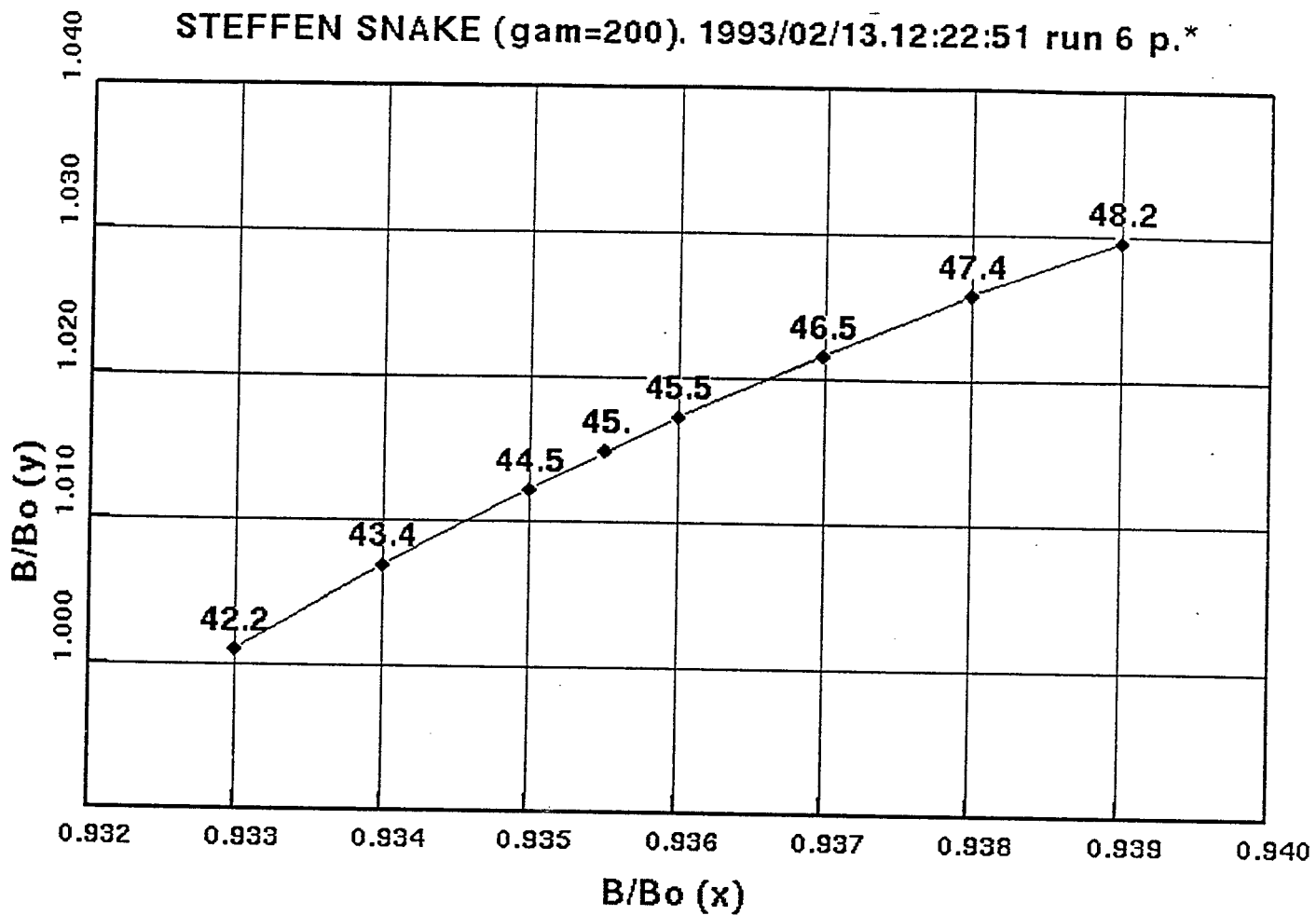


Figure 5

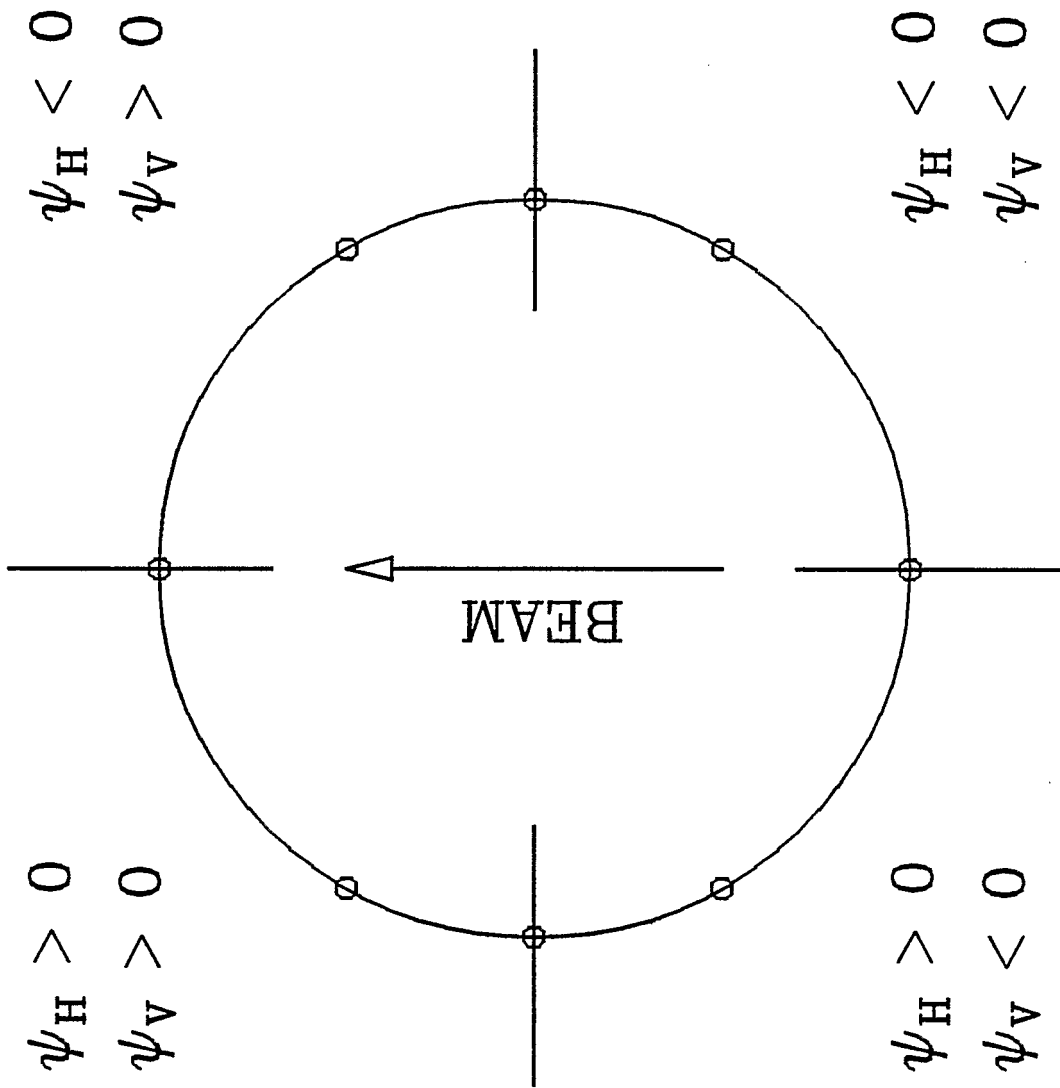
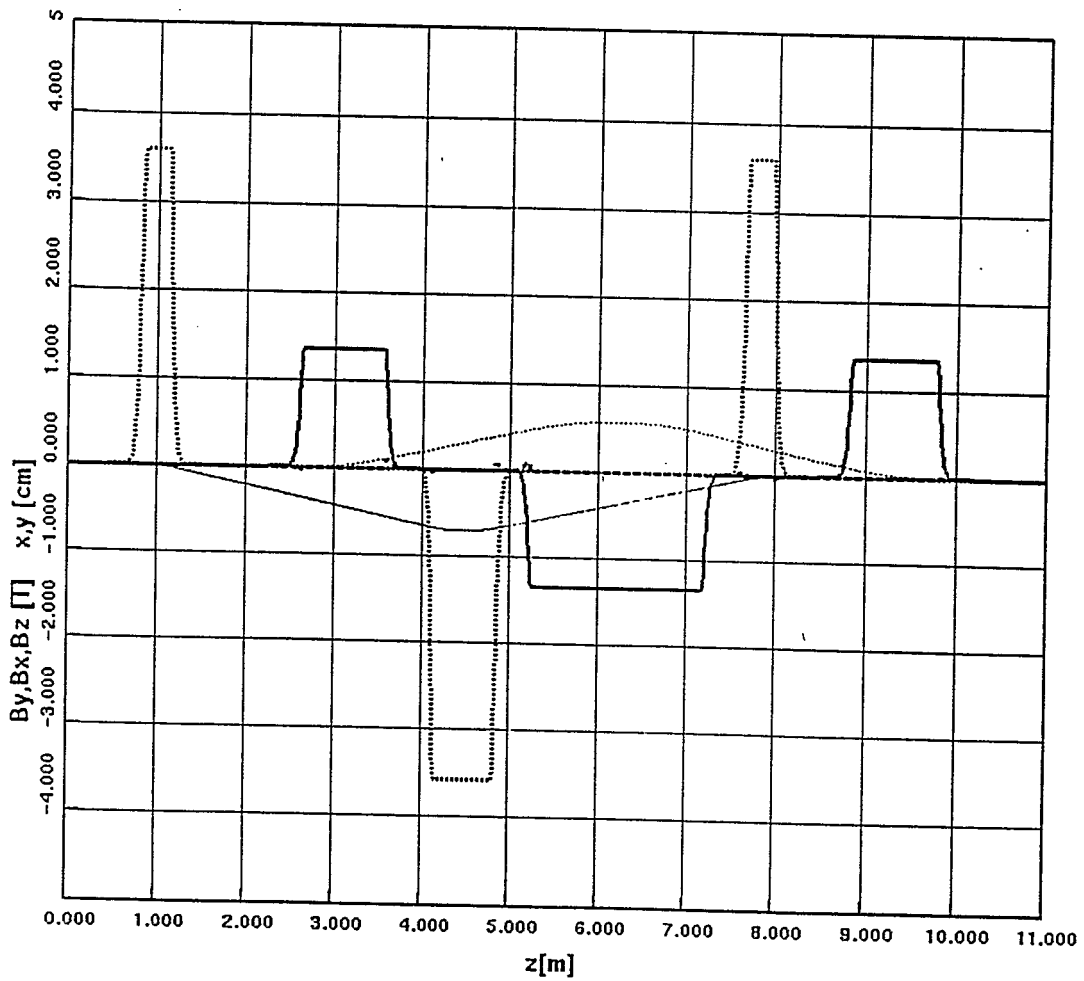
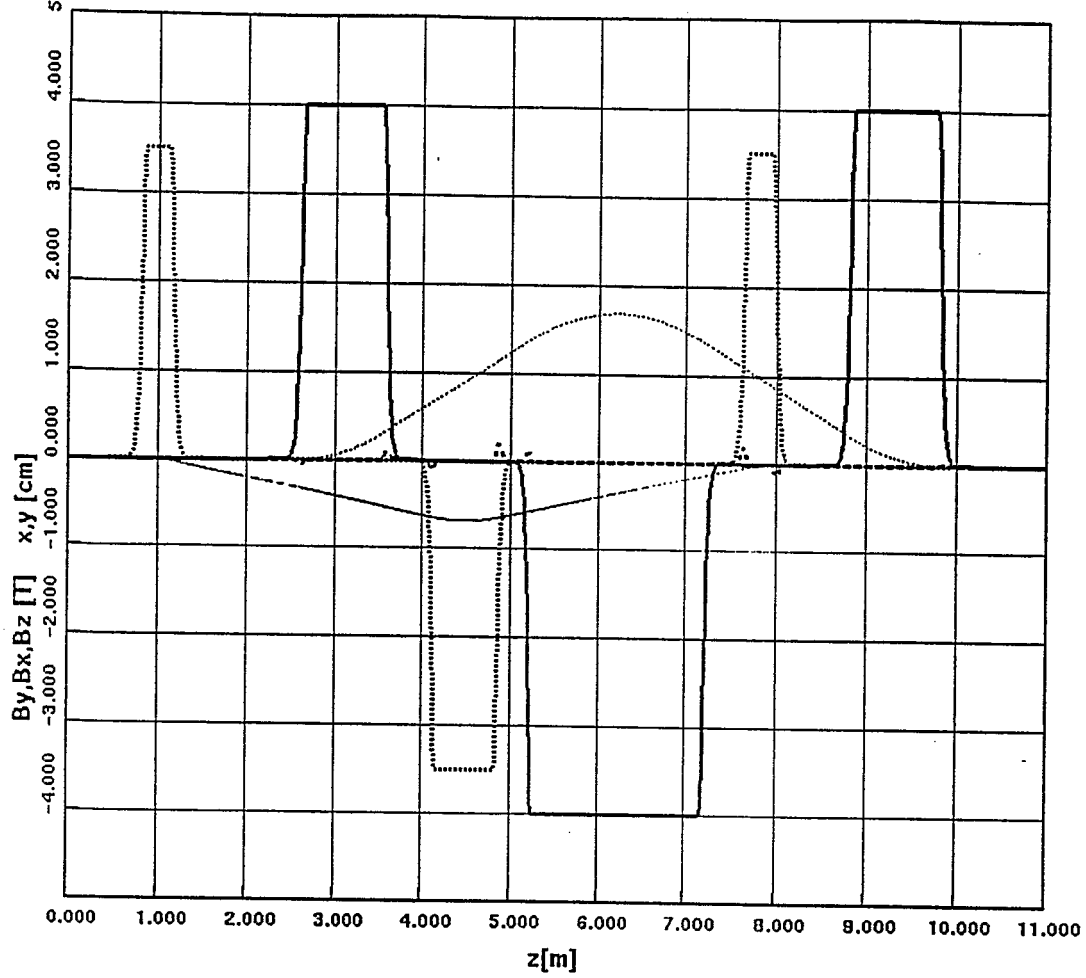
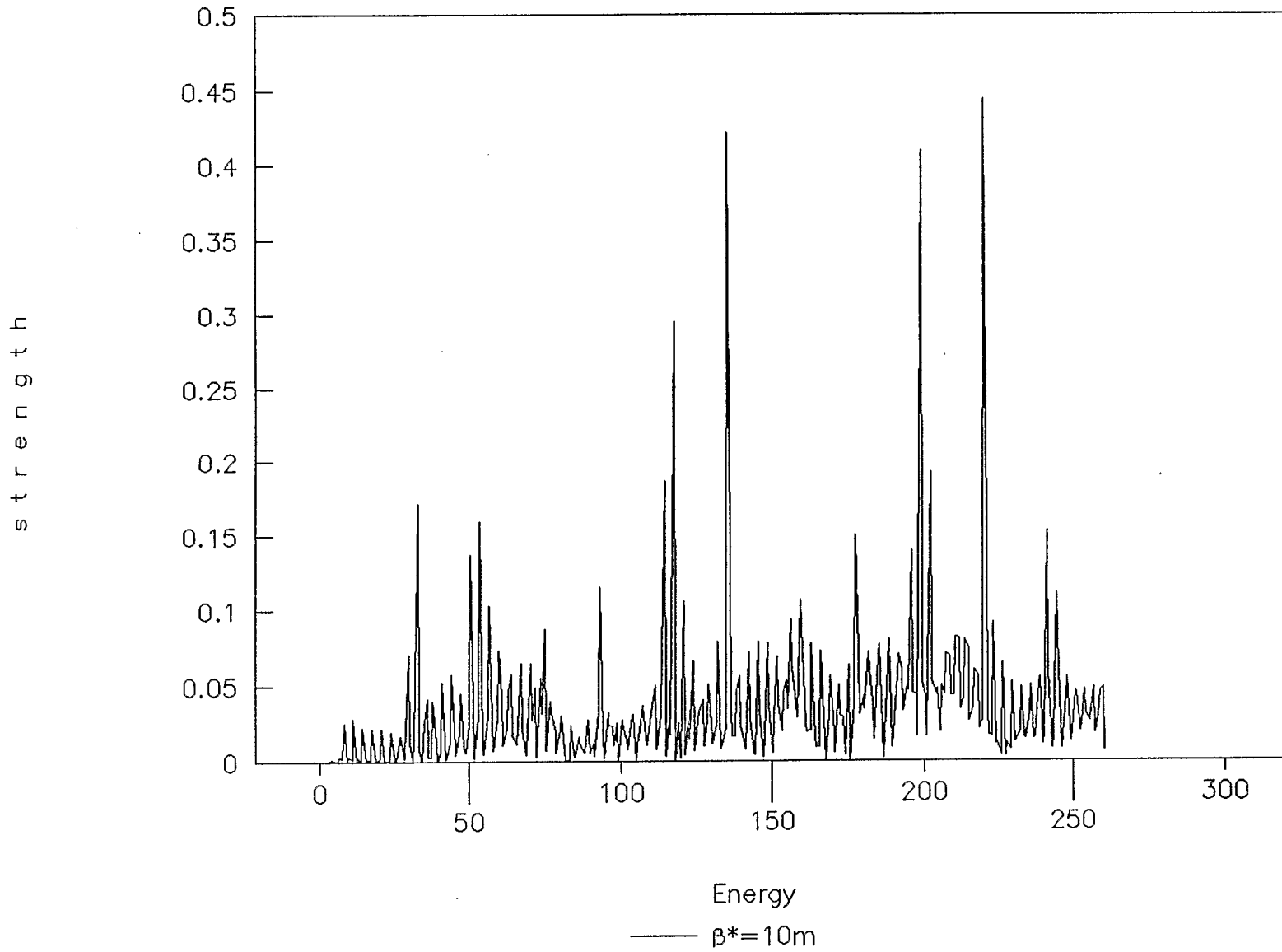


Figure 6



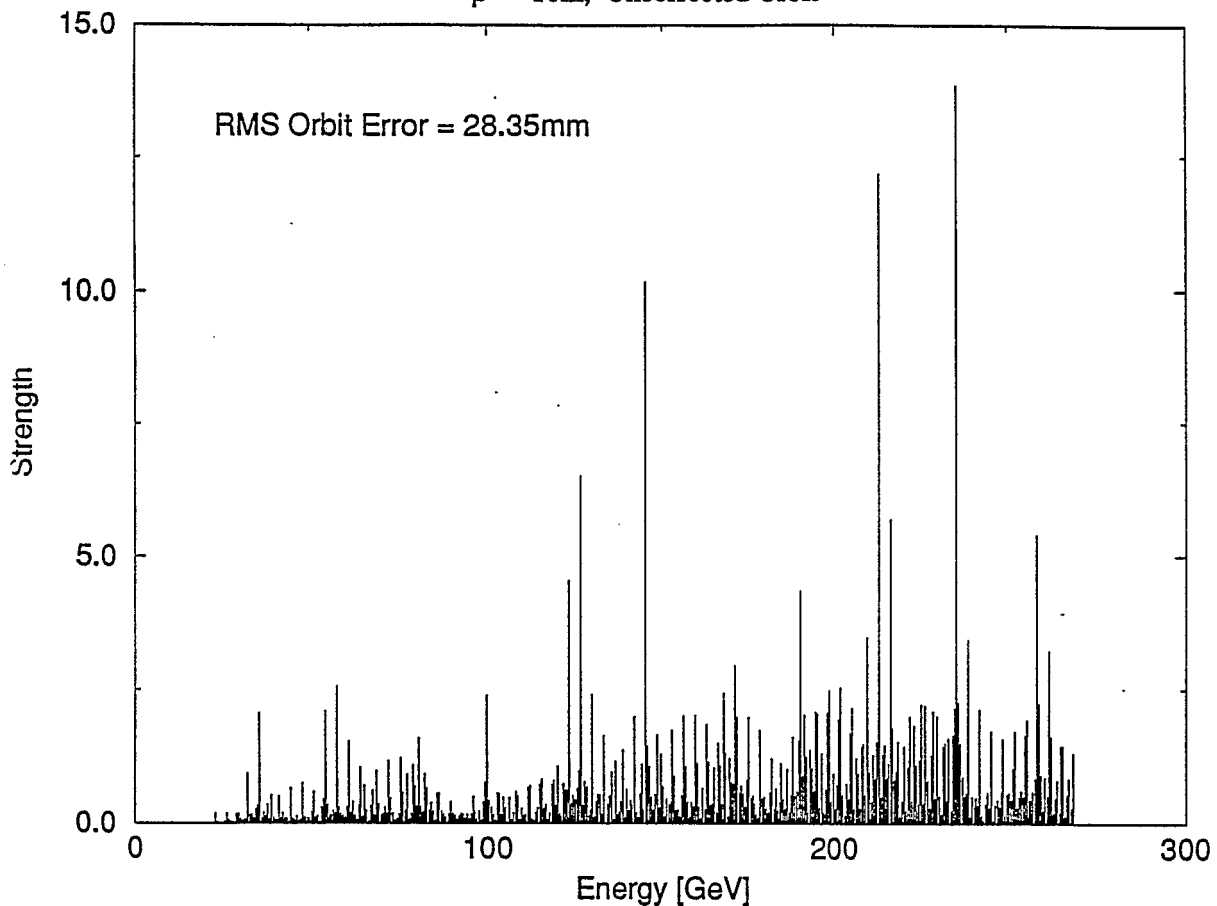
# Intrinsic resonance strength

RHIC; invar. emittance  $10 \mu\text{rad}\text{-m}$

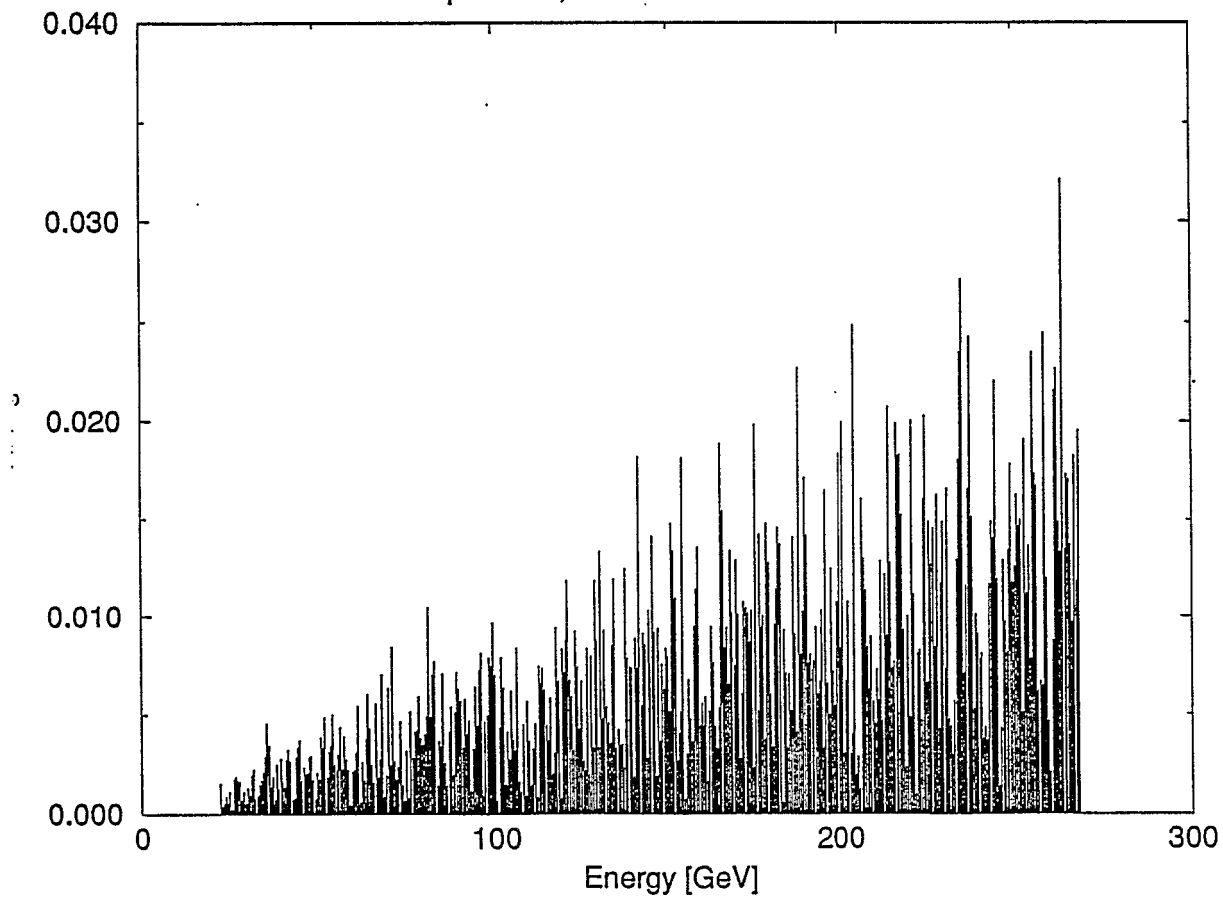


# Imperfection Depolarizing Resonances

$\beta^* = 10\text{m}$ , Uncorrected Orbit



$\beta^* = 10\text{m}$ , Orbit Error = 0.155mm



$\epsilon_{\text{int}}=0.5$     $\epsilon_{\text{imp}}=0.05$     $N_s=2$     $\nu_s=1/2$

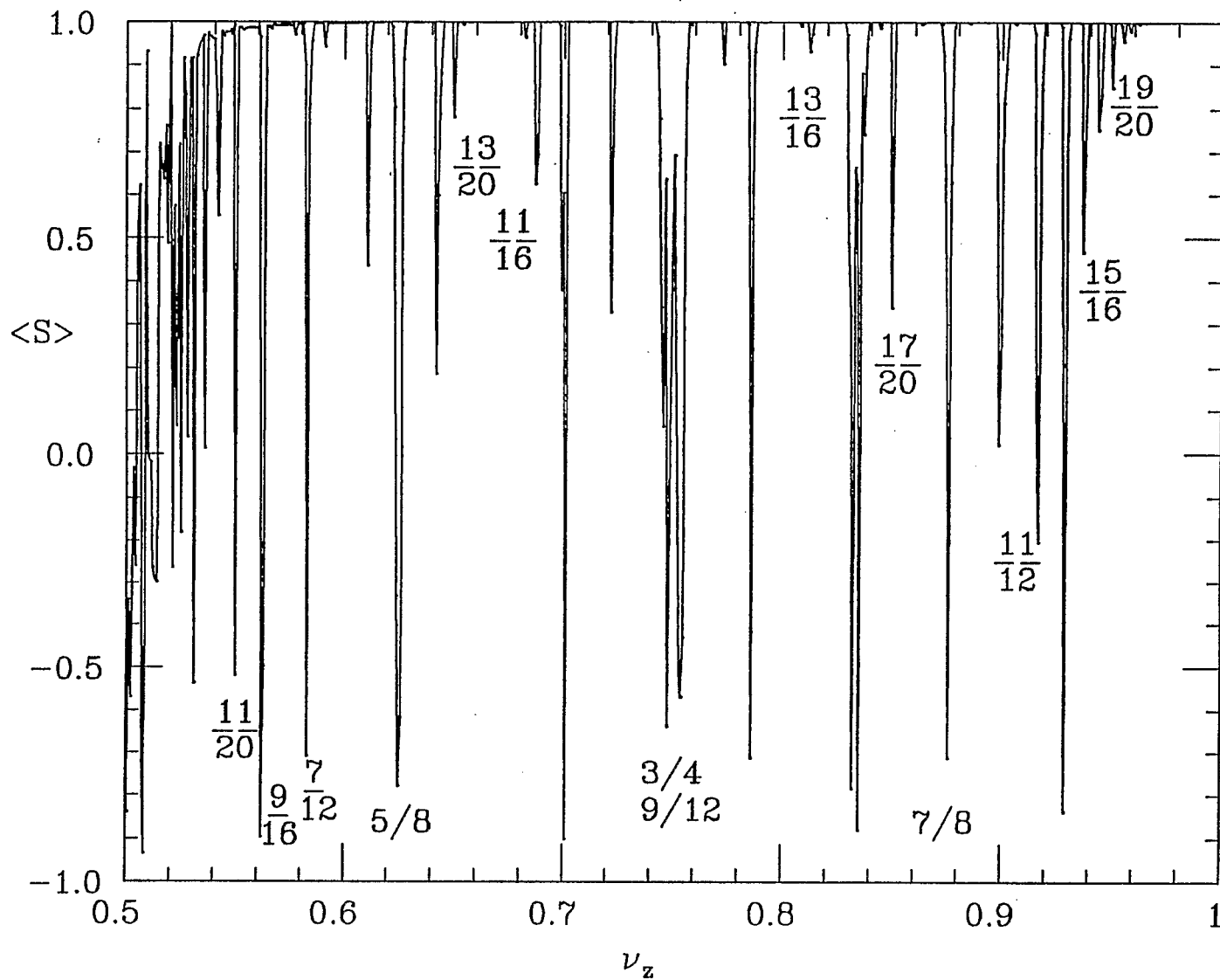
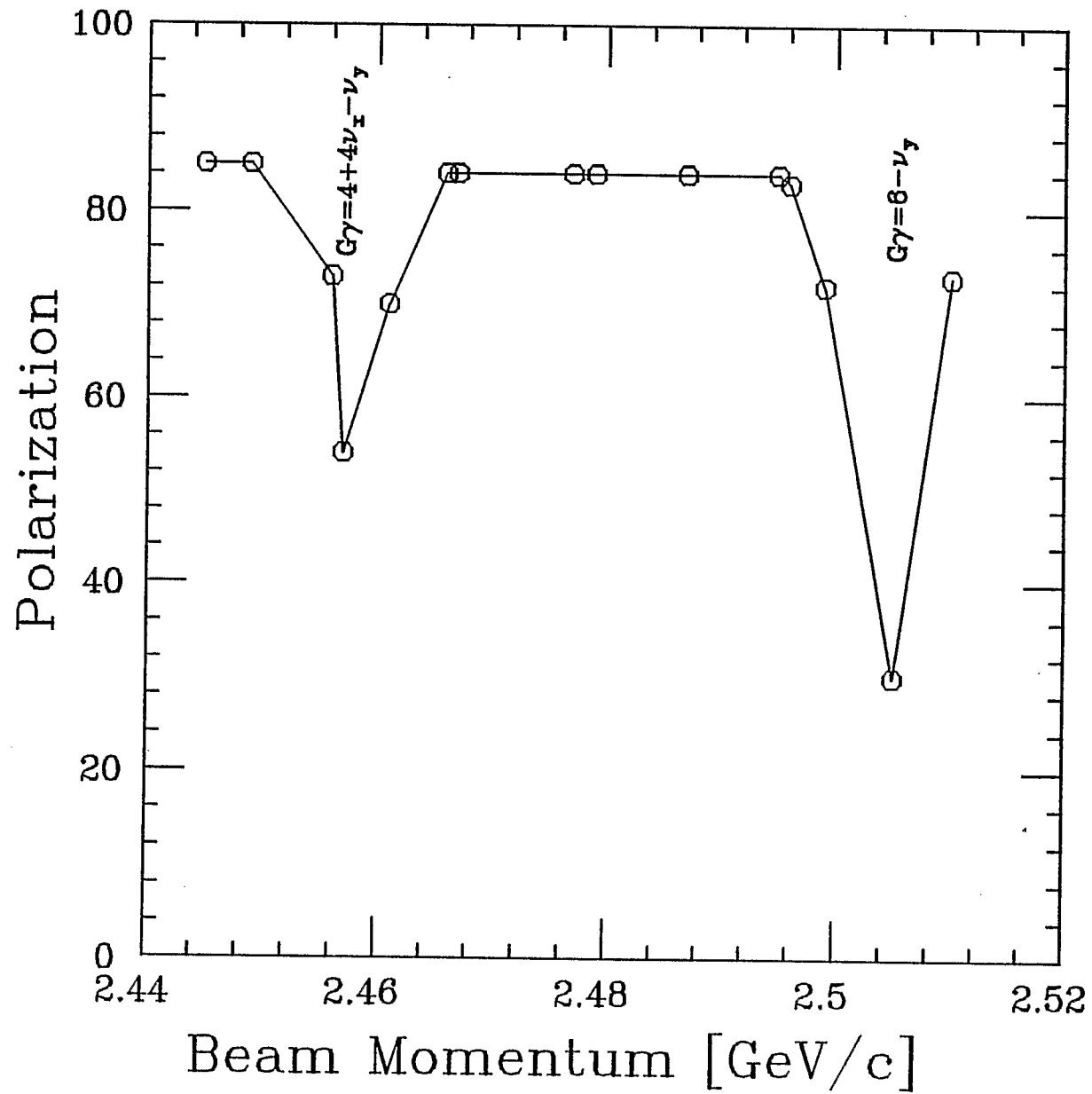


Figure 10

Figure 11



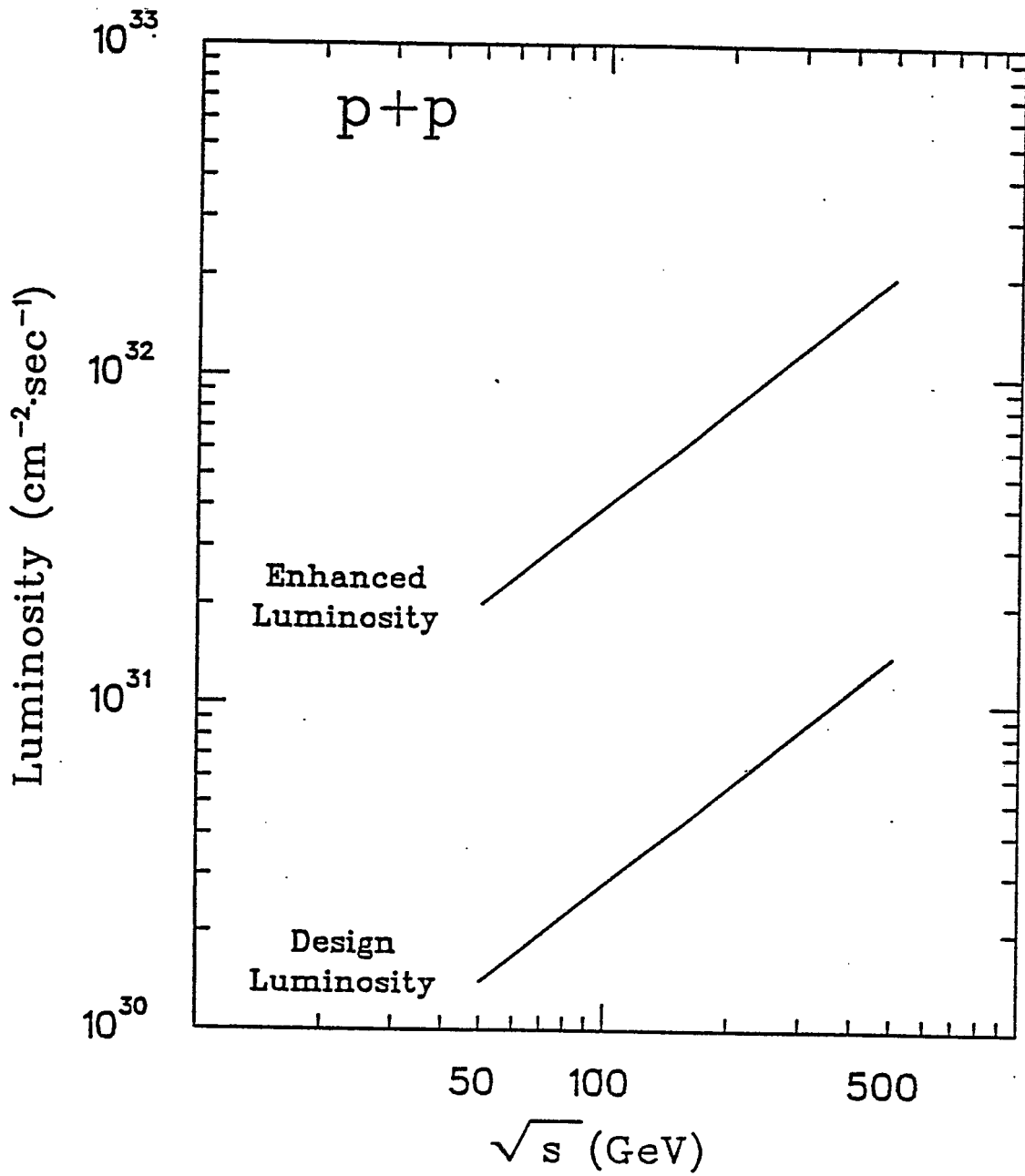


Figure 12



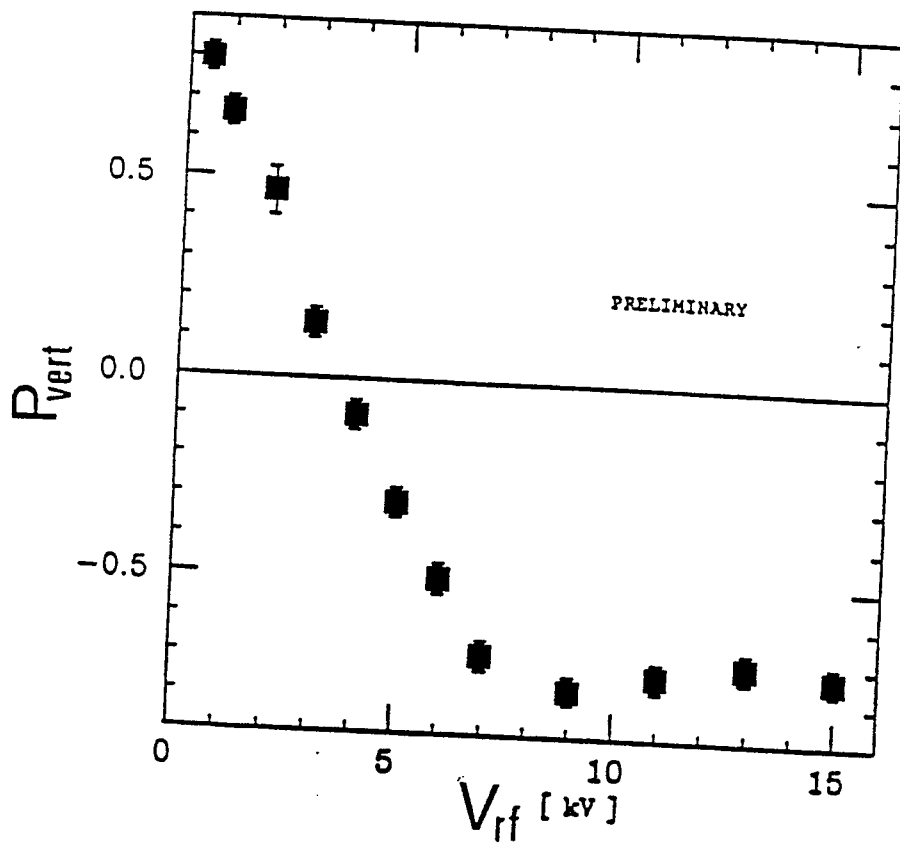


Figure 13