# DATA SHEETS FOR THE AGS HIGH FIELD QUAD 

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January 1988

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 Brookhaven National Laboratory
## U.S. Department of Energy <br> USDOE Office of Science (SC)

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Accelerator Division<br>Alternating Gradient Synchrotron Department BROOKHAVEN NATIONAL LABORATORY<br>Associated Universities, Inc. Upton, New York 11973<br>Accelerator Division Technical Note<br>AGS/AD/Tech. Note No. 293<br>DATA SHEETS FOR THE AGS HIGH FIELD QUAD<br>E. Bleser and M. Tanaka<br>January 14, 1988

This note presents material we have collected on the 24 high field quads installed in the AGS. The information is summarized in a one page data sheet with various appendices giving back-up information. The nominal name for this type of magnet is 6Q24. In the ring there are twelve magnets connected in series located at horizontal maxima and twelve connected in series at vertical maxima. In AGAST they are designated QHORZ and QVERT. The 4 skew quads, 8 Q16, are not discussed in this note.

New precision power supplies (BNL AGS Spec-879) will be available for these magnets soon. In order to derive full benefits from these supplies it will ultimately be necessary to take account of the hysteresis loops shown in Appendix $B$.

1. Physical Description
A. Core

Core length
24
3.218
60.960

Gap radius
8.174
B. Coil

Turns per pole - 16
Conductor dimensions
Cooling hole diameter
Conductor cross section
Cooling hole diameter
Conductor cross section
Approximate conductor length per pole
C. Vacuum Tube

| Inches |  |
| :--- | :--- |
|  |  |
|  |  |
| 24 | 60.960 |
| 3.218 | 8.174 |

$0.812 \times 0.343 \quad 2.062 \times 0.871$
$0.156 \quad 0.396$
0.2593 sq. in. 1.673 sq. cm.

1056

| O.D. | 5.875 | 14.923 |
| :--- | :--- | :--- |
| I.D. | 5.745 | 14.592 |

Material - SS 304
D. Locations

Horizontal quads, 12 places A17...L17 c
Vertical quads, 12 places A3...L3 c
Distance from mag to mag 64
Distance from mag to quad 25 .
25.58 .4
2. Electrical Properties

2682
d
a
14.923
14.592
b

Per Magnet
Resistance, measured - 11.7 milliohms
calculated at $60^{\circ} \mathrm{F}$ - 11.2 milliohms
calculated at $100^{\circ} \mathrm{F}$ - 12.3 milliohms
Inductance, measured - 2.32 millihenrys
Per 12 Magnets
Resistance, measured - 208 milliohms
Inductance, measured - 30 milliohms

## Properites of the High Field Quadrupole (continued)

3. Magnetic Properties

$$
\frac{R e f}{f}
$$

G x L/I, measured - 41.08 Gauss/Amp
Chosen effective length - 68 cm
Then gradient, $G / I$ - 0.6041 Gauss/(cm x Amp)

Calculated gradient, G/I - 0.6019 Gauss/(cm x Amp)
$\mathrm{B} 5 \mathrm{x} \mathrm{L/I} \quad-\quad-8 \times 10^{-4} \mathrm{Gauss} /\left(\operatorname{Amp} \times \mathrm{cm}^{4}\right)$
4. Tune Control
$\Delta Q_{\mathrm{x}}=\left(0.0262 \times I_{\mathrm{x}}+0.0121 \times I_{\mathrm{y}}\right) / \mathrm{P}$
$\Delta Q_{y}^{x}=-\left(0.0121 \times \mathrm{I}_{\mathrm{x}}+0.0261 \times \mathrm{I}_{\mathrm{y}}\right) / \mathrm{P}$
$I_{x}=$ current in Amps in 12 HF horizontal quads
$I_{y}=$ current in Amps in 12 HF vertical quads
$P=$ beam momentum in $\mathrm{GeV} / \mathrm{c}$

## References

a. Mechanical Drawing D03-M-1222-5.
b. Mechanical Drawing C-D05-M-959-4.
c. Devices in the AGS Ring, M. Zguris, K. Brown. June 1, 1987.
d. WHERE ARE THE AGS MAGENTS? E. Bleser, TN 215, May 20, 1985.
e. Appendix A.
f. Appendix B.
g. Appendix C.


## Appendix A

## THE ELECTRICAL PROPERTIES OF THE HIGH FIELD QUADRUPOLES

The material in this appendix was collected or calculated largely by Bi.ll Leonhardt. Figure Al shows the pressure drop across a magnet as a function of the flow rate, calculated at $T=70^{\circ} \mathrm{F}$ and $100^{\circ} \mathrm{F}$ and measured at $70^{\circ} \mathrm{F}$. Figure A2 shows the calculated temperature rise versus pressure drop for DC currents of 600 and 800 Amps. Table A1 gives the measured pressure drop across each magnet as it operates in the AGS. The magnets usually run at less than 500 Amps with a duty cycle of less than $50 \%$, suggesting that they are normally quite cool. There are at present 8 water circuits on a magnet, connected in two series of four. The cooling could be increased by connecting as many as all eight circuits in parallel.

The most recent measurements of the resistance and inductence of a magnet (Joe Funaro and Lou Mazurakis, November 1987) confirm earlier measurements and give:

$$
\begin{array}{ll}
\text { Per Magnet: } & R=11.7 \text { milliohm } \\
& L=2.32 \text { millihenry }
\end{array}
$$

Per 12 Magnet System: $R=208$ milliohm
$L=30$ millihenry

Table AI
WATER SUPPLY READINGS ON HIGH FIELD QUADS

| Location | Supply <br> psi | Return <br> psi | Pressure <br> Drop <br> psid |
| :--- | :---: | :---: | :---: |
|  |  |  |  |
| A-3 | 100 | 32 | 68 |
| B-3 | 103 | 34 | 69 |
| C-3 | 105 | 35 | 70 |
| D-3 | 106 | 44 | 62 |
| E-3 | 106 | 40 | 66 |
| F-3 |  |  |  |
| G-3 | 85 | 18 | 67 |
| H-3 | 87 | 20 | 67 |
| I-3 | 90 | 21 | 69 |
| J-3 | 93 | 23 | 70 |
| K-3 | 96 | 24 | 72 |
| L-3 | 98 | 32 | 66 |
|  |  |  | 67.8 Mean |

PRESSURE DROP vs FLOW RATE

Figure A1
!sd ‘dOyd ヨynSS $\exists y d$

## Appendix B

THE MAGNETIC PROPERTIES OF THE HIGH FIELD QUADRUPOLE MAGNETS

The High Field Quadrupole Magnets were designed in 1958. Other than the mechanical drawings, which seem to be complete and readily available, we have found little information on them. There are 24 magnets installed in the AGS and one spare. This spare was carefully measured by the Magnetic Measurements Group of the ADD Department in March of 1987. We would like to thank Peter Wanderer for supporting this work, John Herrera for stalwart efforts to understand the results, and Bob Gottschalk and his group for carrying out the measurements.

The measurements were made with a rotating coil, 74 inches long, giving the gradient integrated through the length of the magnet. The results are given in Table B1 and plotted in Figure B1. There are some hysteresis effects, not particularly visible in this plot, which give us some options in defining the gradient as a function of the current. For our present purposes, it is sufficient to use a straight line connecting the origin and the point of the highest positive current. These two points then give us:

$$
\mathrm{G} \times \mathrm{L} / \mathrm{I}=41.08 \text { Gauss } / \mathrm{Amp}
$$

This simple result is sufficient for most present uses of the magnets.
It is useful to consider also the gradient rather than the integrated gradient. However, we do not have a point measurement of the gradient in the center-of-the-magnet. Therefore, we must pick an effective length somewhat arbitrarily We have defined the effective length of this magnet to be:

$$
L_{\text {eff }}=68 \mathrm{~cm}
$$

Then

$$
G / I=0.6041 \text { Gauss } /(\mathrm{cm} \times \text { Amp })
$$

A simple calculation gives:

$$
\begin{aligned}
G / I & =0.8 \times \pi \times \mathrm{N} / \mathrm{a}^{2} \\
& =0.6019 \mathrm{Gauss} /(\mathrm{cm} \times \text { Amp })
\end{aligned}
$$

where $N=$ the number of turns and $a=$ the radius of the bore. We have chosen $L_{\text {eff }}$ to be the core length plus $0.86 * a$, which is a very reasonable choice, so we feel the overall picture is consistent.

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Appendix B (continued)
```

For convenience we replot Fig. Bl with the integrated gradient divided by the effective length to get Fig. B2. In order to see the hysteresis effects, we replot in Fig. B3 the data from Fig. B2 with $0.6041 * I$, the postulated straight line, subtracted from the data. The squares indicate increments of positive current; the circles increments of negative current. Figure B3 clearly shows the hysteresis effect and the residual fields, which correspond to a current error of about 1 Ampere. If we require greater accuracy we shall have to undertake a program of cycling the magnets. From the tune shifts calculated in Appendix C, we see that at heavy ion injection, the residual fields in the high field quadrupoles can give tune shifts as large as $\pm 0.1$.

For completeness we show plots of $G x \operatorname{L} / \mathrm{I}$ vs I and $\mathrm{I} /(\mathrm{G} \times \mathrm{L}$ ) vs ( G x
L) in Figures B4 and B5.

For a quadrupole magnet, the gradient Bl , is the primary allowed term. The allowed higher harmonics are the duodecupole, B5, and the twenty-pole B9. These have been measured and the results are shown in Figs. B6 and B7. These fields are very small and can presumably be ignored. The measuring system is a tangential coil, rotating at a radius of 3.81 centimeters, whose voltage output is analyzed to give harmonics through B19. The data suggests that this remarkably fine system is even able to measure the next allowed term, the twenty-eightpole, B13. The non-allowed terms have been measured, are all small, and since they are presumably due to random errors in magnet construction, no meaningful statements can be made about the results from only one magnet. Our conclusion is that we can treat these magnets as perfect quadrupoles.

TABLE EI
INTEGRATED FIELD MEASUREMENTS

|  | B1＊L | EStL | E 3 ＊L |
| :---: | :---: | :---: | :---: |
| CURRENT | Gauss | Gauss／cm＊＊4 | Gauss／cm＊＊0 |
| Q20 | －36．4E | Q． 8121 | －1．37E－67 |
| 99．9 | 4096．12 | $-10.1278$ | －1． $52 E-65$ |
| 193．6 | 818日． 98 | －12． 150 | －3．04E－65 |
| 299．4 | 12e95． 55 | $-2.303$ | －4． $56 E-105$ |
| 397 E | 16393.60 | －0． 317 | －6． 23 SE － 05 |
| 449． $\mathrm{Q}_{1}$ | 1844日． 23 | $-4.357$ | －E． $78 E-85$ |
| 399．2 | 1641E． 29 | －40．317 | －6． Q3E－65 $^{\text {a }}$ |
| 299．4 | 1ころころ．20 | － 2.30 | －4．53E－b5 |
| 199．6 | 8こ30．19 | －2．162 | －3．84E－25 |
| 97.9 | 4130.87 | － 4.8188 | $-1.53 E-25$ |
| Q． 21 | 4 20.6 | － 2.8021 | －7． $814 \mathrm{E}-8 \mathrm{C}$ |
| －99．9 | －4080． 82 | Q． 079 | 1． $51 E-105$ |
| $-199.6$ | －8177．41 | 0.159 | 2．93E－Q5 |
| －299．4 | －1ここ76．55 | Q． E E 38 | 4．48E－125 |
| －399．1 | $-16371.54$ | Q． 318 | 6． 21 E－ |
| $-449.0$ | －184ED． 34 | 6． 357 | E．75E－15 |
| －399． 1 | $-16389.66$ | Q． 318 | E． $2=E-25$ |
| －Е99．4 | －12368．75 | Q． 239 | 4． $50 .-65$ |
| $-199.6$ | －8玉19．43 | Q． $1 \in Q$ | E．99E－b5 |
| $-99.9$ | $-4125.36$ | Q． 2181 | 1．48E－05 |
| $\square 0^{0}$ | $-36.36$ | Q． 812 C | －E． $77 E-27$ |

(SONVSNOHL) 'SSn $\forall$ O

Figure B2
GRADIENT HYSTERISIS LOOP
EFFECTIVE LENGTH $=68 \mathrm{~cm}$

Figure B3

Figure B5
(s-OL S SWNIL)
$8^{\mathrm{mo} / \text { SSn }} \mathrm{SO}$

## Appendix C

## THE TUNE SHIFT PRODUCED BY THE HIGH FIELD QUADRUPOLES

The usual expression for a tune shift produced by a quadrupole perturbation is:

$$
\Delta Q=-\frac{1}{4 \pi} \quad \int \beta_{K} d L
$$

The quadrupoles are located at the appropriate extremum of the beta functions, and the beta functions are therefore very flat through the quads. Values calculated by MAD are shown in Table C1 and Figures Cl and C2. In addition:

$$
\begin{aligned}
\mathrm{K} & =-\mathrm{G} / \mathrm{Bp} \\
\mathrm{GdL} & =41.08 \mathrm{I} \\
\frac{1}{\mathrm{BP}} & =\frac{3 \times 10^{-7}}{\mathrm{P}}
\end{aligned}
$$

Giving

$$
\begin{aligned}
\Delta Q & =\frac{1}{4 \pi}{ }_{\sum}^{12} 41.08 \times 3 \times 10^{-7} \mathrm{BI} / \mathrm{P} \\
& =117.7 \times 10^{-7} \mathrm{BI} / \mathrm{p}
\end{aligned}
$$

The units for these formula are Gauss, centimeters, Amperes, and $\mathrm{GeV} / \mathrm{c}$. Using the $\beta$ values for $15 \mathrm{GeV} / \mathrm{c}$ we have:

$$
\left[\begin{array}{l}
\Delta Q_{\mathrm{x}} \\
\Delta Q_{\mathrm{y}}
\end{array}\right]=\left[\begin{array}{cc}
0.0262 & 0.0121 \\
-0.0121 & -0.0261
\end{array}\right]\left[\begin{array}{l}
\mathrm{I}_{\mathrm{x}} / \mathrm{p} \\
\mathrm{I}_{\mathrm{y}} / \mathrm{p}
\end{array}\right]
$$

TAELEECA
BETA FUNOTTONS CALCLIMTED EY MAD

| $55-3$ |  |  |  | 55－17 |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| F | 1 | EETA $X$ | EETA $\gamma$ | $\operatorname{ETA} \mathrm{X}$ | BETA $V$ |
| GEV／E | 1 | met゙をrs | meterss | meters | meters |
|  |  |  |  |  |  |
| Qu | 1 | 14．35边 | 17．791 | 17.794 | 14n $=34$ |
| D， | 1 | 14． 170 | 17．94E | 17：951 | wased |
| 12． 23 | 1 | 13．8玉w | 18．1家 | 18． 1.8 | 13，718 |
| 砋： | 1 | 13,445 | 18.401 | 18．410 | 13， 3 B |
| D． 62 | 1 | 11． 3 3 E | EQ 58 | ED． 583 |  |
| $\square_{\square} \square_{6}$ | 1 | 11． 54.3 | En， 656 | E17，7E | 11． 1 ¢ ${ }^{\text {a }}$ |
| 7，78 | 1 | 11.1 .77 | E8．7501 | ED．ATS | 11． 18.87 |
| 15 | 1 |  | EE．16E | EEES3 |  |
| 8 | 1 | 121035 | E\＃， 4 4 | $E \mathrm{Em}$ | 10． E \％ |
| E5 | 1 | 10．420 | $\sum \pm .156$ | E E． 47 | 108， 5 |
| E7 | 1 | 10.4030 | Ex．$E 47$ | Emigl | 10， 251 |
| 9 | 1 | 120487 | $\pm \pm \pm 73$ | EE，59 | 10．438 |
| 3 F | 1 | 12065 | E\＃5 54 | Eizum | 10，ERS |
|  |  |  |  |  |  |

MINIMUM BETA vs MOMENTUM

Figure C1

```
in
- 21 -
```

MAXIMUM BETA vs MOMENTUM

| 23 |
| :--- |

Figure C 2


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