



BNL-104863-2014-TECH

AGS/AD/Tech Note No. 447;BNL-104863-2014-IR

MAD Simulation: HEP/SEB Extraction (FY1996)

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October 1996

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U.S. Department of Energy

USDOE Office of Science (SC)

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Technical Note

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11 October 1996

Version 1.1

(Ver. 1.1: Fixed an error on Dx in Fig.4.5 and Table 3.)
(Ver. 1.0 : 10.Oct.96)

ABSTRACT

The slow extracted beam (SEB) at the AGS has been used successfully since 1968 for the High Energy Physics Program. This note describes the results from a MAD tracking study of SEB resonant extraction up to straight section F13 using the actual FY96 HEP/SEB run parameters and an estimated emittance for the high intensity beam of $(5-6) \cdot 10^{13}$ ppp.

1. Introduction

The slow extracted beam (SEB) from the AGS to the East Area has been operational since March 1968 for the High Energy Physics Program (HEP)[1]. The present SEB system with an electrostatic wire septum at straight section H20 was implemented in 1979-1980[2] together with the new SEB switchyard[3]. In the SEB mode the debunched circulating beam is extracted to over 1 to 2 sec (spill length) for the typical repetition time of 2.4 to 4.0 sec by exciting a third integer resonance ($3 \cdot Q_h = 26$). The long and uniform spill allows HEP counter experiments to take data smoothly without overloading their detectors. For typically good SEB operation during the pre-Booster HEP program, the measured extraction inefficiency was 3-5 % for the extracted beam intensity of $1.0\text{-}1.6 \cdot 10^{13}$ ppp at $p = 25\text{-}29$ GeV/c[4,5]. In FY1995-6 HEP/SEB runs, we have had a 1-3 % inefficiency for the record high intensity of $5.0\text{-}6.0 \cdot 10^{13}$ ppp at $p = 25.5$ GeV/c according to the AGS Daily Report (see Fig. 1). This SEB performance looks great if it is the case. However, the unexpected large beam losses were observed in the Switchyard, indicating that the extracted beam emittance might have become too large. This note describes the results from a MAD tracking study of AGS SEB extraction up to the middle of straight section F13 (ssF13md) using the FY96 HEP/SEB run parameters and an estimated high intensity beam emittance, hoping they might shed some light on the problem.

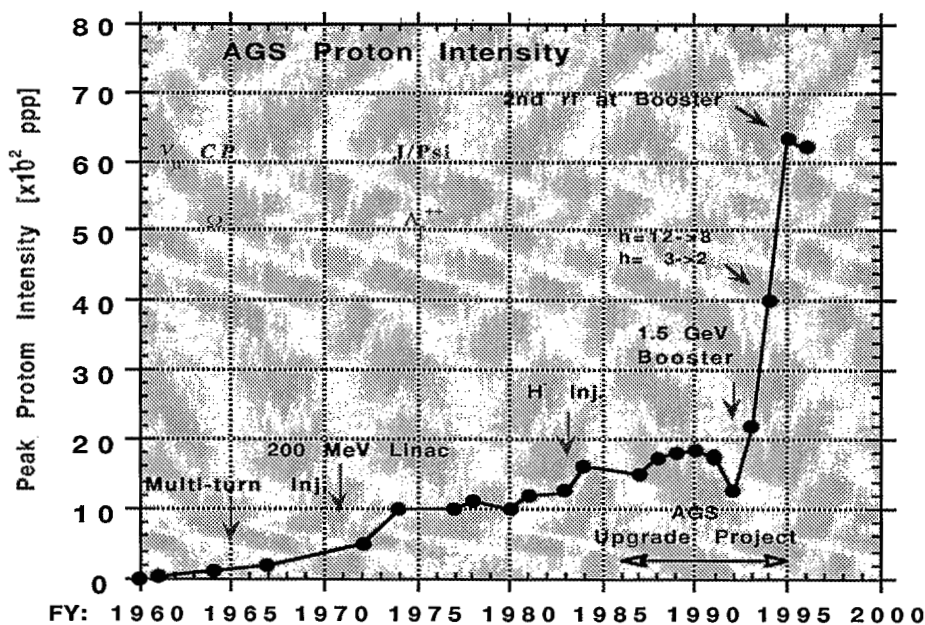


Fig. 1. AGS Proton Intensity Performance from 1960 to 1996.

2. SEB Extraction

2.1 SEB Extraction System

The present SEB system consists of an electrostatic wire septum at sssH20 [ESH20], a thin septum magnet at F05 [SMF05], and a thick ejector magnet at F10 [SMF10], together with two $3/2 \lambda$ local orbit deformations [HPBLW, FPBLW] and 4 drive sextupoles [DS] which excite the third-integer resonance [2,4]. The layout of the main SEB extraction elements in the ring are schematically shown in Fig. 2.

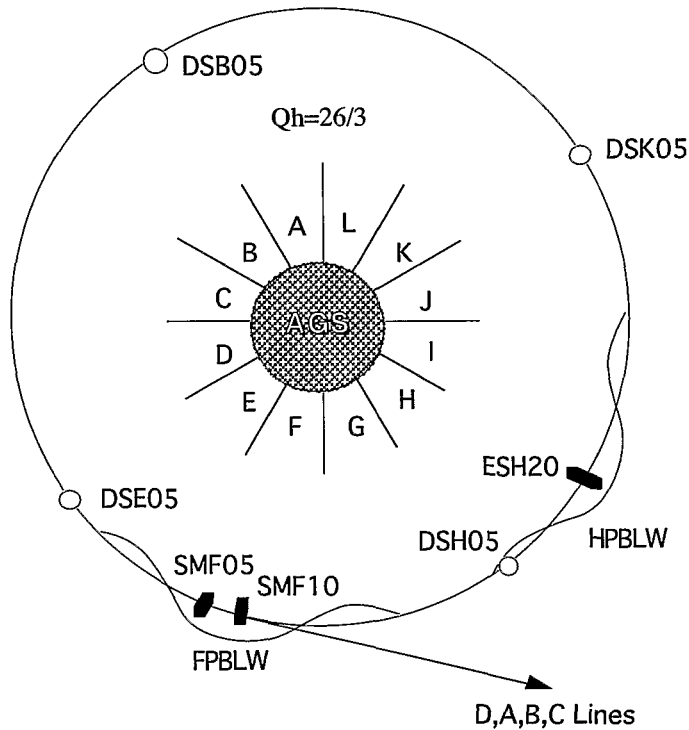


Fig. 2. The layout of the SEB extraction elements.

2.2 Third-integer Resonant Extraction Principle

AGS SEB extraction is performed by controlled excitation of nonlinear third integer betatron resonance of the main ring at $3 \cdot Q_h = 26$, which corresponds to the tune value Q_h of the unperturbed ring near the center of the vacuum chamber. The factor "3" arises from the drive sextupole fields which excite the resonance and the integer 26 from the 26-th Fourier harmonic component of the fields. For each momentum, there is a set of three linear separatrices which define a stable horizontal phase space area with a triangle shape. The stable triangular area can be reduced by adjusting the tune to be in the vicinity of $Q_h = 26/3$ either directly by varying the tune quadrupole strength or indirectly through slow beam steering or by increasing the drive sextupole strength. At the AGS, the tune quadrupole and drive sextupole strengths are kept constant during extraction, and the beam is pushed slowly from its stable position to the resonance by a slightly negative main magnet flattop slope so that different momentum particles progressively move on to the resonance. Extraction starts when the beam emittance fills the

stable area for the higher end of momentum distribution. As the stable area slowly shrinks to zero for these particles, other parts of the momentum distribution are subsequently extracted.

The extraction inefficiency depends mainly on the fraction of particles which strike the septum of the first extraction device (ESH20) and is given by $\text{ineff.} = f \cdot w/s$, where w is the effective septum thickness, “ s ” is the step size (spiral pitch, the growth of the resonant betatron amplitudes in the final few turns before extraction) and f is a factor (~ 1.5) to account for the horizontal beam density variation. The step size s of a particle depends on its momentum, initial amplitude and its proximity to the resonant tune value. The maximum value of the step size is limited by the available aperture at ESH20, downstream and the circulating beam emittance. For the analytical aspect of SEB extraction, see ref. [1].

2.3 SEB Extraction Process at AGS

Once the beam is accelerated to the 25.5 GeV/c flattop, the rf is turned off and the beam is debunched and the the momentum spread is increased. Extraction devices ESH20, SMF05 and SMF10 as well as high field quadrupoles are turned on before rf turn-off, and the 4 Drive Sextupoles [DS], Backleg Windings [FPBLW ,HPBLW] are energized shortly after the rf turn-off.

The equilibrium orbit of the beam at this time is adjusted to be located at approximately 8.0 mm inside of the resonance radius ($R_0 = 0$) and $Q_h = \sim 8.75$. Since the horizontal chromaticity is negative, the beam spirals outward to the radius near the center corresponding to the $Q_h = 8^{2/3}$ resonance by slowly decreasing the the flattop fields in the AGS main magnets. As the beam tune approaches the resonance, the stable region (separatrix) in the horizontal phase plane shrinks for each momentum. For each momentum, there is a set of three separatrices which forms a stable triangular area. The lower the momentum, the larger the stable triangular area is.

Extraction starts when the separatrix area for the highest momentum particles are equal to the phase plane area enclosed by the orbit of the particles of that momentum with the highest betatron oscillation amplitude. Slightly later, particles of a higher momentum with smaller amplitudes will be extracted simultaneously with the particles of a lower momentum with larger amplitudes. Particles outside the separatrix execute oscillations of rapidly increasing amplitude and move outward each of three asymptotes associated with the third integer resonance separatrix. For every *three* turns the change of the particle position (or step size) increases rapidly as it approaches the septum of ESH20 in the phase space. The particle eventually jumps across the wire septum and is kicked outward. After deflection by ESH20, the particle turns around about $2^{3/4}$ orbits and then arrives at SMF05 in the proper phase for further deflection by this septum. At the distance = 24 full betatron oscillations, the angular deflection dx' at ESH20 will, in the linear case, produce no net spatial separation at SMF05 since $dx(F5) = dx' \cdot \sqrt{[\beta(H20)\beta(F5)]} \cdot \sin(\Delta\mu) = \sim 0$, where $\Delta\mu =$ the phase advance between H20 and F5 = $\sim 24 \cdot 2\pi$.

However, the non-linear effects of the sextupole moments of the drive sextupoles and of the main magnets are significant enough to produce the required displacement at SMF05. The beam kicked by SMF05 is finally extracted to the external beam line by the ejector magnet SMF10. A comprehensive description of the AGS SEB extraction system is given in ref.[4].

2.4 SEB Extraction Setup Parameters

During the FY95/96 HEP/SEB runs, the typical machine setup parameters on the SEB flattop were

•Main Magnet	4484 A	! 1.0052 T or $p = 25.725 \text{ GeV}/c$
•HBLW	~267 A	
•FBLW	~297 A	
•DS	~300 A	
•ESH20	~50 kV	! $x_{\text{wire}} \approx 47.0 \text{ mm}$ (upstream)
•SMF05	~2.2 kA	
•SMH10	~4.4 kA	

While there are no recent reliable measurements of the transverse emittance of the circulating high intensity proton beam in the AGS, the consensus is that the 95 % normalized emittance should be $\epsilon_{h,v}^*(95\%) = (75 - 95) \pi \text{ mm-mrad}$, significantly higher than the pre-booster values $\epsilon_{h,v}^*(95\%) = 6 \cdot \sigma^2/B \cdot (p/m) = (25 - 40) \pi \text{ mm-mrad}$. During the FY96 HEP/SEB run, the chromaticity is corrected to be $\{\xi_h, \xi_v\} = \{-1.8, -0.22\}$ with $\{I_{Sh}, I_{Sv}\} = \{150, 70\} \text{ A}$ and $\{Q_h, Q_v\} \approx \{8.67, 8.78\}$ with $\{I_{Qh}, I_{Qv}\} \approx \{80, -160\} \text{ A}$ on the flattop at $R_0 = 0$ to cope with the observed vertical beam instability. Table 1 summarizes some basic parameters of ESH20, SMF05 and SMF10.

Table 1. Parameters of the SEB extraction devices

	ESH20	SMF05	SMF10	
gap(v) x width(h)	20 x 10	31.75 x 17.5	60.0 x 25.9	mm x mm
length	2.30	0.667	2.06	m
I_{max}	(80 kV)	2.2	5.0	kA
B_{max}	(80 kV/cm)	0.15	0.97	T
Septum width	0.051(W/Re)	0.76(Cu)	15.9(Cu)	mm
Kick(nominal)	0.43	1.1	18.5	mrad

3. MAD Simulation

The simulation of the AGS SEB extraction was made by using the BNL MAD code[6] including the following elements:

- Main Combined Function Magnets containing multipoles up to sextupoles,
- High Field Tune Quadrupoles,
- Chromaticity Sextupoles,
- SEB Extraction devices (HPBLW, FPBLW, DSs, ESH20, SMF05, SMF10).

MAD was run to obtain the working point parameters and extraction closed orbit deformations using the actual FY96 HEP/SEB run parameters. Results are shown in MAD FTWISS output format in Table2 and in Fig. 3.

Table 2. FTWISS output for the SEB working point ($p = 25.5$ GeV/c and $dp/p = 0.0068$ at $I_{DS} = 300$ A, $\{I_{Qh}, I_{Qv}\} = \{80, -160\}$ A, $\{I_{Sh}, I_{Sv}\} = \{150, 70\}$ A, $I_{HPBLW} = 270$ A and $I_{FPBLW} = 300$ A).

ELEMENT POS. NO.	SEQ ELE NAME.	I I	HORIZONTAL						VERTICAL				
			DIST [m]	β_h [m]	α_h [1]	μ_h [2 π]	x [mm]	x' [.001]	D_x [m]	D_x' [1]	I I	β_v [m]	α_v [1]
BEGIN	AGS	1	0.000	21.750	-1.790	0.000	37.255	2.465	1.695	0.118	11.575	1.037	0.000
138	DSK05	1	150.529	27.328	0.101	1.592	0.400	0.059	2.245	-0.014	9.404	-0.068	1.658
313	DSB05	1	352.298	18.550	0.044	3.770	2.155	0.068	1.811	-0.014	10.816	-0.012	3.846
488	DSE05	1	554.067	24.802	-0.304	5.967	3.096	0.008	1.659	0.008	9.367	0.064	6.035
551	SMF05us	1	620.666	23.970	0.161	6.648	39.367	1.040	1.922	-0.016	10.206	0.016	6.767
552	SMF05ds	1	621.326	23.776	0.133	6.652	40.053	1.040	1.911	-0.016	10.229	-0.049	6.777
569	SMF10us	1	636.795	15.940	1.258	6.824	37.497	-4.030	1.656	-0.115	13.301	-1.149	6.948
570	SMF10md	1	637.810	13.555	1.094	6.835	33.410	-4.030	1.544	-0.115	15.812	-1.326	6.959
571	SMF10ds	1	638.825	11.502	0.929	6.848	29.322	-4.030	1.432	-0.115	18.683	-1.503	6.969
580	ssF13md	1	648.043	23.514	-0.083	6.957	15.375	-1.919	1.965	0.007	10.058	-0.020	7.064
690	DSH05	1	755.836	24.085	-0.107	8.112	-18.121	-0.087	2.330	0.019	9.960	-0.003	8.236
742	ESH20us	1	804.826	14.677	-1.356	8.646	31.714	2.465	1.435	0.118	17.148	1.441	8.759
743	ESH20ds	1	807.076	21.750	-1.790	8.667	37.255	2.465	1.695	0.118	11.575	1.038	8.785
END	AGS	1	807.076	21.750	-1.790	8.667	37.255	2.465	1.695	0.118	11.575	1.038	8.785
C = 2 π R =			807.075641 m	Q _h =			8.666550	Q _v =			8.784913		
Δ C =			8.323980 mm	Q _h ' =			-15.424595	Q _v ' =			-1.953253		
α =			0.141348E-01	β_h^{\max} =			29.304740 m	β_v^{\max} =			24.930812 m		
γ_{tr} =			8.411152	D_x^{\max} =			2.664284 m	D_y^{\max} =			0.000000		
				x^{\max} =			43.607278 m						

$$\dagger Q' = \xi \cdot Q = dQ/(dp/p)$$

snapshot

96/10/02 09.58

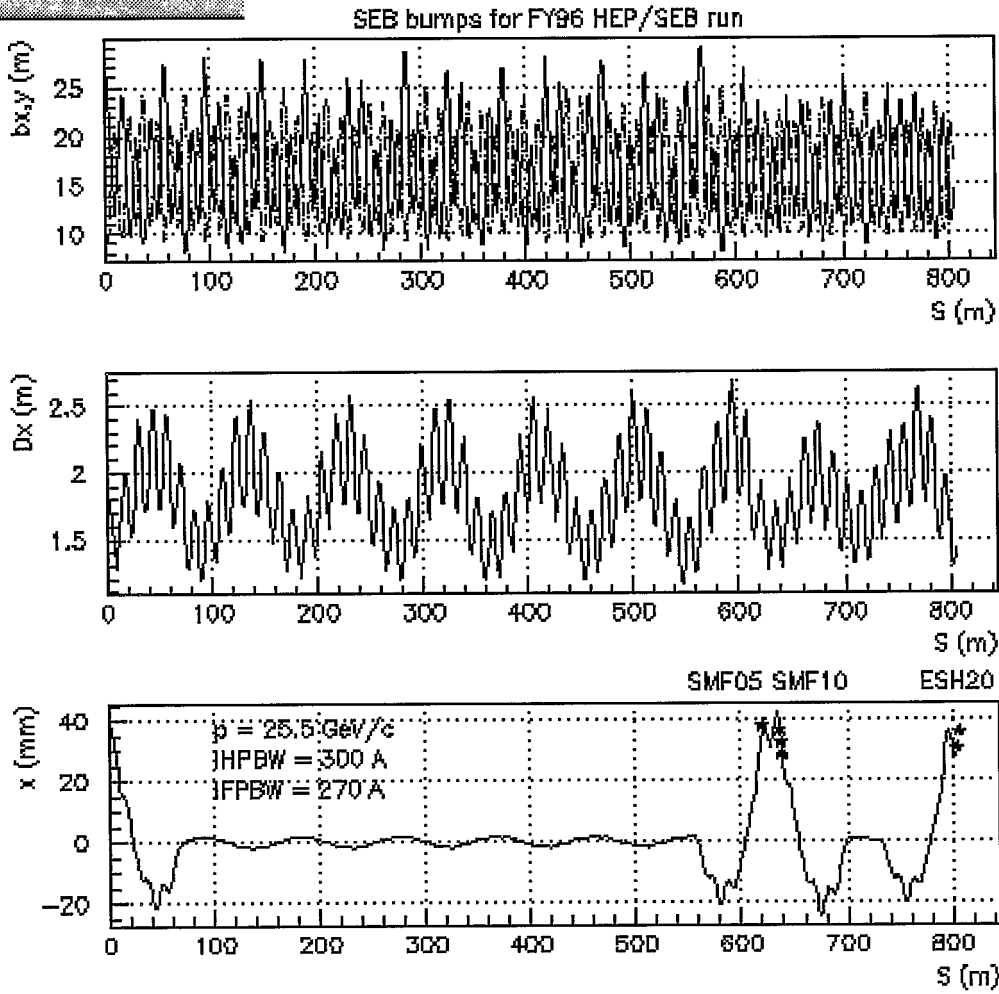


Fig. 3. $\beta_{h,y}$, D_x and x (SEB bumps) as a function of S starting from ESH20ds at the SEB working point.

3.1 Tracking of the Resonance

Particle tracking in the horizontal plane has been performed using MAD with an interactive graphical display utility[6]. During the extraction process, the stability of particles depends on both their momentum dp/p and betatron amplitude (x, x') . The dp/p values corresponding to the zero and nominal emittance of the circulating beam have been obtained by trying as close as possible to the resonance for the zero emittance case, and by trying out various sets of initial values for the nominal emittance case. Tracking was performed for each case, varying the initial values in small steps until two adjacent trajectories corresponding to the largest stable and the smallest unstable emittance were obtained as illustrated in the Fig. 4.0.

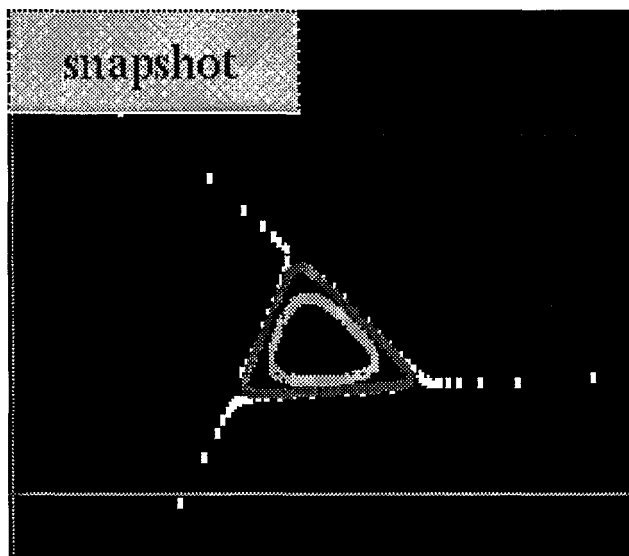


Fig. 4.0. Phase space (x, x') output at SMF05 from tracking particles from various initial conditions for the fixed dp/p .

The nominal emittance of the circulating beam is estimated by calculating an area of the largest stable triangle. The tracking results show that the dp/p value for the zero emittance and for the nominal emittance ($\epsilon_h = \sim 3.3 \pi$ mm-mrad which is corresponding to $\epsilon^*_h = \sim 90 \pi$) are $dp/p = 0.000650$ and $dp/p = 0.0000785$, respectively. Tracking the unstable particles is used to calculate the step size at ESH20. Tracking was stopped once the particle had jumped the first wire septum. The positions of the successive turns are interpolated along the extraction separatrix to find the value of x' at the septum position and the next (x, x') values after 3 turns. Fig. 4.1 shows the results at the beginning of ESH20 and also at the end with a septum kick of 0.35 mrad. The step size is calculated to be ~ 7.5 mm for the zero emittance case and ~ 3.7 mm for the nominal one at $x_{sep} = 47.0 \text{ mm}^\dagger$, and the instantaneous external beam emittance (an

[†] Practically all the beam loss occurs at the upstream end of the septum. The setp size is (4.6 - 9.5) mm at the end of the septum.

area of the trapezoid) is $\sim 0.60 \pi$ mm-mrad. Tracking continues with starting values (x, x') at four points which define the external beam emittance. These four particles are traced through the extraction channel and receive 0.35 mrad kick at ESH20, 1.45 mrad kick at SMF05 and 19.0 mrad kick at SMF10 up to the middle of ssF13, where the extracted beam should be ~ 380 mm away from the central orbit, free from the fringe field of the ring magnets.

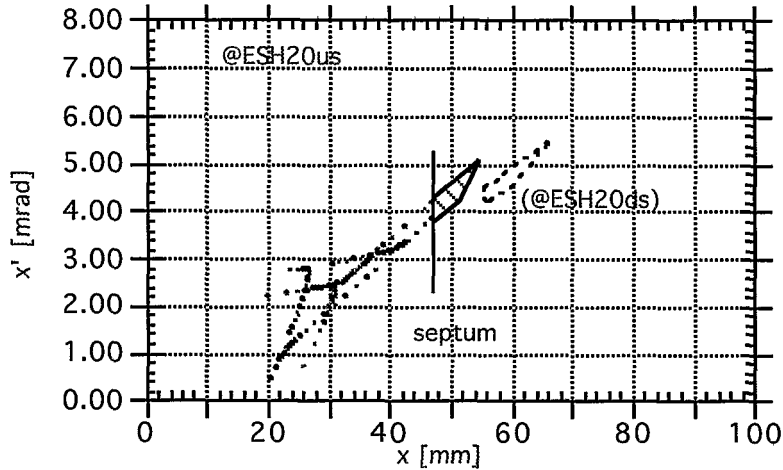


Fig. 4.1. Phase space (x, x') at the beginning of ESH20.

3.2 Tracking of the Extracted Beam

From the first septum (ESH20) to the second septum (SMF05), tracking is done for $2^{3/4}$ turns for particles that jumped the wire septum and were kicked outward by 0.34 mrad by ESH20. The clearance between the circulating beam and the extracted one at SMF05us is ~ 1.14 mm, marginal but acceptable as the SMF05 septum thickness is 0.9 mm. The step size is (9.5 - 31) mm so that the beam could fill up the available horizontal aperture of 32 mm at SMF05. The external beam emittance ϵ_H^{ext} is calculated to be $\sim 1.8 \pi$ mm-mrad, increased from $\sim 0.60 \pi$ mm-mrad at ESH20. The results are shown in Fig. 4.2 with the beam kicked further outward 1.46 mrad by SMF05. The particles kicked by SMF05 are then traced through the third septum SMF10. The clearance at SMF10us is ~ 23.1 mm, quite adequate for the 16 mm thick septum as seen Fig. 4.3. The external emittance is calculated to be $\sim 2.0 \pi$ mm-mrad at this point, which corresponds to $\epsilon_H^{\text{ext}} = \sim 54 \pi$ mm-mrad, 60% of the internal emittance of 90π mm-mrad..

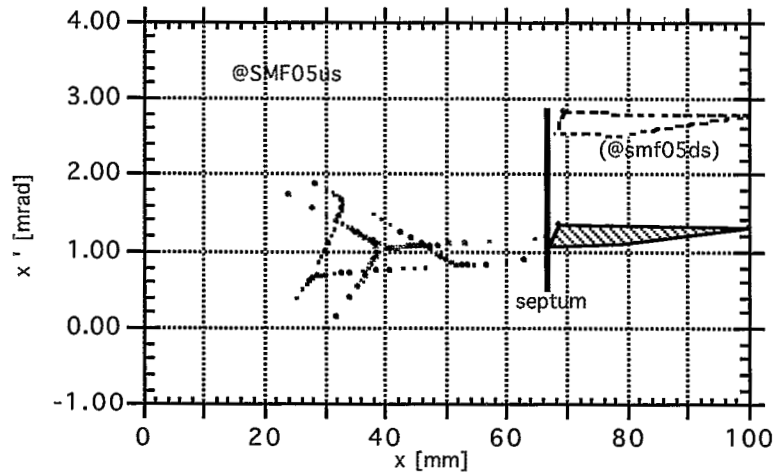


Fig. 4.2. Phase space at the beginning of SMF05.

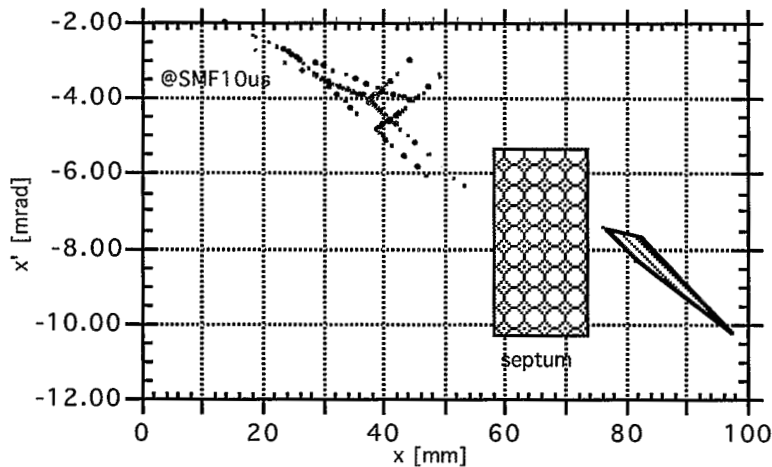


Fig. 4.3. Phase space at the beginning of SMF10.

The beam extracted by 19 mrad kick at SMF10 will cross the stray field of the two horizontally defocusing AGS main magnets (F11 and F12) and pass through the fringe field of the F13 main magnet at $x > 250$ mm, where its gradient reverse sign and becomes defocusing. Since the standard AGS model uses the field parameterization which is good up to ± 100 mm, we built a simple model of parameterizations of F11, F12 and F13 magnet based on the values calculated by the POISSON code and the knowledge of where the beam must pass through at each magnet. Fig. 4.4 shows the phase space plot at ssF13md. The external emittance has not changed and is $\sim 2.0 \pi$ mm-mrad. Fig. 4.5 and Table 3 show the Twiss parameters and beam positions as a function of S from SMF05ds to ssF13md.

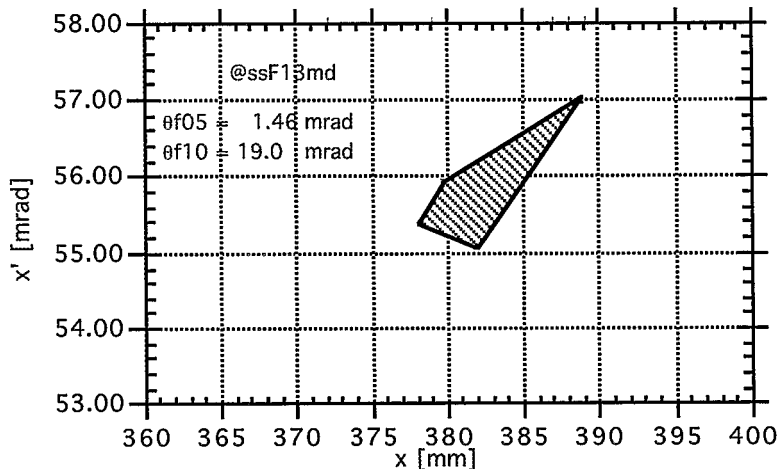


Fig. 4.4. Phase space at ssF13md.

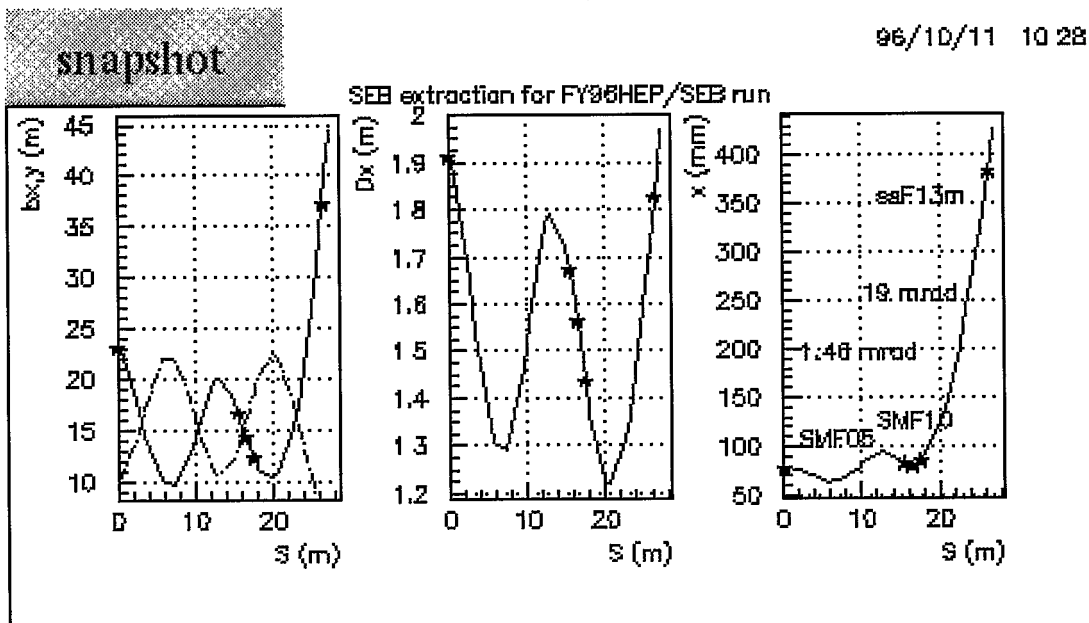


Fig. 4.5. $\beta_{h,y}$, D_x and x as a function of S starting from SMF05ds for the extracted beam.

Table 3. FTWISS output for the extracted beam with $\theta_{f05} = 1.46$ mrad and $\theta_{f10} = 19.0$ mrad.

LATTICE PARAMETERS FOR BEAM LINE: "SEB F05ds-to-F13md ", DELTA(P)/P = 0.000036 SYMM = F

ELEMENT SEQ I		HORIZONTAL							VERTICAL				
POS. ELE	I	DIST	β_h	α_h	μ_h	x	x'	D_x	D_x'	I	β_v	α_v	μ_v
NO. NAME .	I	[m]	[m]	[1]	[2 π]	[mm]	[.001]	[m]	[1]	I	[m]	[1]	[2 π]
00	SMF05ds	1	0.000	23.000	0.124	0.000	77.000	2.660	1.911	-0.017	10.320	-0.048	0.000
17	SMF10us	1	15.469	16.674	1.255	0.171	83.126	-8.215	1.670	-0.109	13.302	-1.128	0.170
18	SMF10md	1	16.484	4.287	1.098	0.182	79.610	1.281	1.563	-0.119	15.766	-1.301	0.181
19	SMF10ds	1	17.499	12.218	0.941	0.194	85.730	10.778	1.436	-0.128	18.582	-1.474	0.191
22	MMF11	1	20.015	10.419	-0.373	0.232	124.107	22.500	1.226	0.018	22.497	0.418	0.210
25	MGF12	1	22.631	15.141	-1.672	0.266	200.588	40.631	1.311	0.128	17.475	1.656	0.230
27	MGF13	1	25.628	30.736	-4.015	0.289	339.742	55.955	1.688	0.238	8.047	1.450	0.269
28	ssF13md	1	26.390	37.175	4.439	0.293	382.360	55.955	1.827	0.238	6.061	1.156	0.287

4. Results

• Lets first compare our tracking results at ssF13md for the FY96 HEP/SEB with the values quoted as the nominal values used for design in the AGS SEB switchyard design report[3] and 1979 Weisberg's measurements[2]:

	Design Report	Weisberg	FY96 MAD	
$\epsilon^{*ext}_h(95\%)^\dagger$	31.4 π	37.7 π	54.0 π	mm-mrad
β_h	36.68	58.42	37.175	m
α_h	-4.565	-6.6	-4.439	
$\epsilon^{*ext}_v(95\%)^\dagger$	52.3 π	45.5 π	(90.0 π)	mm-mrad
β_v	3.76	3.30	6.106	m
α_v	1.032	0.87	1.156	
D_x			1.827	m

\dagger NB: The values in ref.[2] and [3] are unnormalized 99% values at $p = 29$ GeV/c and are converted to $\epsilon^*(95\%)$. The MAD values assume that $\epsilon^{*int}_{h,v}(95\%) = \epsilon^{*ext}_{h,v}(95\%) = 90 \pi$ mm-mrad (no vertical blow-up).

The FY96 MAD simulation results are consistent with both the switchyard design report values and Weisberg's values except the emittance, which is much higher, merely reflects the assumed larger internal emittance of $\epsilon^{*int}_{h,v}(95\%) = 90 \pi$ mm-mrad for the high beam intensity of $5\text{-}6 \cdot 10^{13}$ ppp. It should be noted that the MAD tracking predicts that the external horizontal emittance should be $\sim 60\%$ of the internal emittance once the momentum sweep during spill is corrected. These values should be compared with the recent measurements once they are available[7].

• Due to the chromaticity correction made in FY96 to cope with the vertical beam instability,

- the instantaneous momentum spread $\Delta p/p$ should have increased from 0.023 % to 0.057 % for a total $\Delta p/p \approx 0.4\%$ after debunching as it is proportional to $1/\xi_h$,
- the difference between the step size for the zero and for the nominal emittance increased from 3.0 mm to 3.8 mm at ESH20us and 14.6 to 22.1 mm at SMF05us. The horizontal beam size could be already too large at SMF05, causing extra beam loss.

- The 1.14 mm clearance at SMF05 is marginal. It could be improved if we could increase the ESH20 kick though it is better to use the ESH20 kick directly rather than using the non-linear effect to create the gap at SMF05[8].
- The overall general characteristics of the SEB extraction obtained in this study are consistent with the previous simulation studies [2,8].

5. Conclusions

The AGS SEB extraction had been studied by the MAD tracking using the actual FY96 HEP/SEB machine setup parameters and an estimated emittance $\epsilon^*_h(95\%) = 90 \pi$ mm-mrad for the high intensity beam of $5\text{-}6 \cdot 10^{13}$ ppp to predict the extracted beam parameters at ssF13md.

- The predicted beam parameters (α, β) at ssF13md are consistent with the switchyard design values and also with Weisberg's.
- However, the external emittance is *substantially larger* than the previous values simple because we use the estimated value of $\epsilon^*_{int_h}(95\%) = 90 \pi$ mm-mrad rather than $(30 - 40) \pi$ mm-mrad.
- It is crucial for us to know exactly what the internal beam emittance is. We should make every effort to measure both the horizontal and vertical beam emittance on the flattop and its intensity dependency in a *controlled fashion*.
- Both horizontal and vertical tracking from SMF05 to ssF23md for the extracted particles can be repeated using a full field map rather than the simple parameterization of the F13 main magnet and taking particle distribution into account [9].

Acknowledgments

The author would like to thank Jim Niederer for his continuous efforts on improving the BNL MAD code and developing the user-friendly interactive graphical display utility. Without his help this tracking study would be more tedious and time consuming work.

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