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Longitudinal stability of individual bunches in the AGS - Booster

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LONGITUDINAL STABILITY OF INDIVIDUAL BUNCHES IN THE AGS - BCOSTER

AD Booster Technical Note No. 98

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To estimate the longitudinal stability of individual bundes in the AGS Booster one calculates the following complex quantity

$$U'-iV' = -i \frac{2e \operatorname{I}_{P} \beta^{2}(Z/n)}{\pi |2| E (\Delta E/E)_{FWHM}} \cdot \frac{Q}{A}$$
 (1)

where

e, electron charge

Q , particle charge state

A , partide mass number

B, particle relocity to light velocity ratio = V/c

Z, the complex beam-environment coupling impedance

n, the harmonic number of the instability

E, the particle total energy /a.m.u.

(DE/E) FWHM, the full-width half-maximum velchie bunch energy spread Also

$$\eta = \gamma_{\tau}^{-2} - \gamma^{-2}$$

where

Eo, particle rest energy /a.m.u.

and of is a evaluated at the Booster transized tion energy-

Finally Ip is the bunch peak current

$$T_{p} = \frac{Ne \beta c Q}{\sqrt{2\pi'} \delta}$$

N, number of particles in a bunch o, rms bunch length

We are assuming bundes with bi-gaussian distribution.

$$I_{\infty} = \frac{Nec}{2\pi R} Q$$
, average current per sunch for $\beta \to 1$

$$B = \frac{\sqrt{2\pi'} \circ}{2\pi R}$$
, bunding factor

with
$$\sigma_z = \sigma/\beta e$$
 and $\delta = \Delta E/2.355$, sms energy spread

and
$$f_p = c/2\pi R$$
, revolution frequency for $\beta \rightarrow 1$

Hen we can also write (1) as

$$U'-iV' = -0.18i \frac{e \Gamma_{\infty} E_{o} \beta \gamma B(Z/n)}{f_{\infty}^{2}/2/S^{2}} \cdot \frac{Q}{A} \qquad (2)$$

In our notation Z = X + i Y with X > 0 a resistance and Y is positive for a capacitive reactance and negative for an inductive reactance.

In the AGS-Booster we can conceive three modes of operation

- 1. Acceleration of protons from 200 MeV to 1.5 GeV with 3 bundes Each bunch has 5×10^{12} protons and an rms bunde area of 0.10 eV-sec.
- 2. Acceleration of heavy-ions for RHIC in one single bunch. The beam parameters are those appearing in the RHIC COR and also reported in Table I Each heavy-ion learn is accelerated to $\beta y = 0.9$ that is $\beta = 0.669$ The individual bunch rms area is May $0.05 \, \mathrm{eV}$ -s/amu
- 3. Acceleration of heavy-ion for fixed target experiments out of the AGS. In this mode there are 3 bundes each with an intensity three times smaller than in the previous mode for RHIC, and each with an ims bunch area of 0.015 eV-sec/a.m.u. The acceleration is also up to B=0.669.

For all bese modes of operation, the final energy is always below the Booster transition energy ($t_7 = 4.5$).

Table I. Heavy-Ion Beams for RHiC

	Carbon	Sulfur	Copper	Todine	Gold
A Q N × 10 ⁹	12 6 22	32 14 6.7	63 21 4.7	127 29 3.2	197 33 2.2
Binjection	0.1262	0.1002	0.0782	0.0595	0.0463
Bextraction	<		0.669		
S eV-s/amn			0.05		
Z/n in Kohm from space charge Q injection Q extraction		1.86	2.40 0.70	3.16	4.06

The largest contribution to the coupling impedance for the Booster is the "space charge"

$$\frac{Z}{n} = i \frac{Z_0 g}{2\beta g^2} \qquad Z_0 = 377 \text{ ohm}$$

where $g = 1 + 2\log \frac{1}{a}$ is I at injection and 4.5 at extraction for all mode of operation. The value of this impedance for the heavy-ion cases is given in Table I. For the proper beam case

Z/n = i 226 ohm at injection = i 136 ohm at extraction

This impedance is so large that it is hard to imagine an inductive wall impedance of the same magnitude.

Provisional Condusion: If Here is no resistance, the reactance being positive (capacitive) and the accelerating cycle always below the Booster transition energy the individual burdes are always stable.

Only the presence of a resistance in the coupling impedance can cause the lunder to be unstable. We can calculate the tollerances on X/n.

ou X/n.
Observe that the energy dependence of U' is given by the quantity

By B(Y/n) ~ & B

for the space charge ingedance and rince one is so well below the transition energy for all cases. Since B decreases with increasing energy, it is seen that U' has indeed only a very week dependence with energy. We will take B = 0.3 at injection and B=0.03 at top energy for all cases. The results are given in Table II. We show the values of U' with space charge at injection and extraction for each case. Based on the stability diagram shown in Fig. 1 we can then infere the maximum

allowed values for V' - Sina we are below the

transition energy sign (Ko) >0 -

The choice for V' depends critically on the shape of the energy distribution.

The sange of U' for the grotus beam driving the acceleration cycle is shown in Fig. 1 - With the exception of a truncated corne distribution (8) and a tefrist-order parabola distribution (9), the learn bunch is always stable provided V' = 0.4, the limit being set by a second-order parabola distribution (7) at top energy. This corresponds to the resistive impedance limit by X/n < 60 ohm.

Heavy-ion bundes are even mor stable than proben lundes. The varye of U' for all heavy-ion cares is also shown in Fig. 1. At very most U' = 0.11 for Carbon at injection in fixed-target mode. All the distribution considered in Fig. 1 are stable provided V' < 0.5. The tollerance on the resistive inyedamce is very digh: tens of Kohm? It is possible to double the number of heavy ions per lunch and reduce considerally the initial bunds area.

Table II.	Bunch	Stability_	Reguirement	in the Boe	sker
	U'		\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \	X/4	
	Injec.	Extr.			
Proton	0.67	0.90	0.4	60 ohn	1
Fixed Turget:					
Carbon	0.11	0.066	0.5	5.3 K	ohns
Sulfur	0.066	0.041	0.5	8.5	
Copper	0.052	0.033	0.5	10.6	
Iodine	0.034	0.021	0.5	16.7	
Gold	0.019	0.012	0.5	29.2	
RHIC:					
Carbon	0.029	0.018	0.5	19 K	shus
Sulfur	0,018	0.011	0.5	32	
Copper	0.014	0.009	6. 5	39	
Idine	0.009	0.006	0.5	58	
Gold	0.005	0.003	0.5	117	

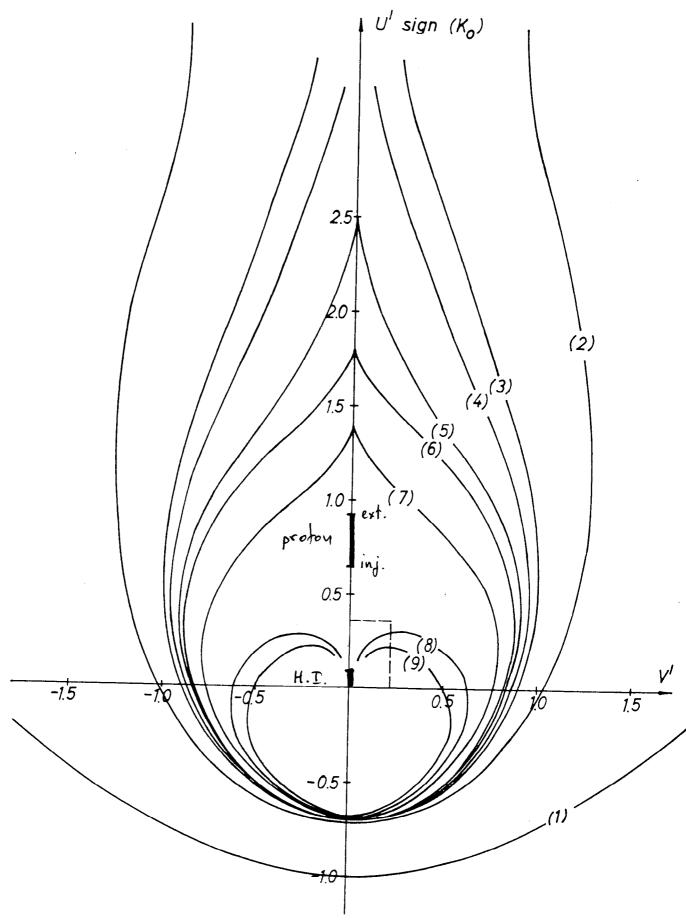


Fig. 1. Stabilitz Dyagram



Diotributions in Fig. 1

	Lorentaian
2	Ganssian
3.	5th-order Parabola
4.	4th-order Parabola
	3 rd - order Paraboly
	Squared Cosine
7	2nd-order Parabola
8	Truncated Cosine
9	1 st- order Parabola