

Stopband Correction of the AGS Booster Down Feed Effect Observed at the AGS Booster and Octupole Imperfection

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**Title: Stopband Correction of the AGS Booster
Down Feed Effect Observed at the AGS Booster and Octupole
Imperfection**

Study Period: April - July 1993**Participants:** C. Gardner and Y. Shoji**Reported by:** Y. Shoji**Machine:** AGS Booster**Beam:****Tools:**

Aim: We had already reported about our studies and data. Here we report a progressing analysis about the normal octupole imperfection.

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I Introduction

We have already reported about the beam loss by the 4th order resonances. At low intensity we observed the considerable beam loss by a 4th resonance [1-4]. It can be happen in the future that we will upgrade the stop-band correction system to correct octupole errors. A special extra windings of the correction package can afford octupole corrections [5]. But we did not know the field strength of octupole imperfections because we did not have any octupole stop-band correction strings. If we knew the field strength of octupole imperfections, it would help to design the power of octupole correction system.

The momentum dependence of the resonance correction was maybe a new phenomena observed at the AGS Booster. It was a down feed from a higher order resonance to a lower with a dispersion displacement. It produced an undesired residual beam loss [6-8]. But on the other hand a measurement of this effect would be a good tool to estimate the strength of higher order imperfections. In this report we will estimate the strength of octupole imperfections of the AGS Booster.

II Down Feed to the Sextupole Resonance

We can estimate the strength of a octupole imperfection from the radial dependence (we will refer it "slope") of sextupole correction. The slope of $3Q_x=14$ correction was measured at $B=1.7\text{kG}$ ($B_p=2.36\text{Tm}$), $dB/dt=0\text{ G/ms}$, $\xi_x=0$ and $\xi_y=0$ [9] as

$$\begin{aligned}\delta N(\cos 14X) / \delta dR_{set} &= 69 / \text{cm} \\ \delta N(\sin 14X) / \delta dR_{set} &= -63 / \text{cm} .\end{aligned}$$

Here $N(\cos 14X)$ and $N(\sin 14X)$ are control digit of the correction used in the computer control program. They are

$$\begin{aligned}N(\cos 14X) &= 125(e/cP) \sum B_x \sqrt{\beta_x} S \cos 14\theta_x \\ N(\sin 14X) &= 125(e/cP) \sum B_x \sqrt{\beta_x} S \sin 14\theta_x\end{aligned}$$

[10] here

$$2S = \int_s^{s+\Delta s} \frac{\partial^2 B_y}{\partial x^2} ds' \quad T/m$$

$$(cP/e) = 3.335641 \quad Tm/(\text{GeV}/c)$$

$$\theta_x = \int_0^s (1/\beta_x) ds' \quad / \quad \int_0^{2\pi R} (1/\beta_x) ds' .$$

The s' is 0 at the beginning of A-super period. The dR_{set} is a control digit of the radial steering by the RF used in the computer control program. It is [9]

$$dR_{set} (\text{cm}) = 319 (dP/P) . \quad (\text{II-3})$$

The down feed of the octupole imperfection to the sextupole imperfection by the dispersion is

$$N(\cos 14X) = 125 (e/cP) \sum B_x \sqrt{\beta_x} 3 O_h (dP/P) \cos 14\theta_x$$

here h is a dispersion function and

$$6 O = \int_s^{s+\Delta s} \frac{\partial^3 B_y}{\partial x^3} ds' \quad T/m^2 .$$

Then the slope is

$$\delta N(\cos 14X) / \delta dR_{set} = 125 (e/cP) \sum B_x \sqrt{\beta_x} 3 O_h \cos 14\theta_x / 319 .$$

The expected octupole imperfection is

$$\sum O \cos 14\theta_x = 319/3/125 (cP/e) [\delta N(\cos 14X) / \delta dR_{set}] / \langle B_x \sqrt{\beta_x} h \rangle .$$

We will use Twiss parameters calculated by A. Luccio and M. Blaskiewicz [11].

$$\langle B_x \sqrt{\beta_x} h \rangle \approx 47 \text{ m}^{5/2} .$$

The integrated harmonic strength of 14th octupole imperfection is

$$\begin{aligned} |\Sigma Oe^{14j\theta x}| &= 319/3/125 \cdot 3.3356 \sqrt{(69^2+63^2)} / 47 \\ &= 5.6 \text{ T/m}^2. \end{aligned}$$

III Down Feed to the Quadrupole Resonances

The parabolic dependence of the quadrupole corrections ($2Qx=9$ and $2Qy=9$) on dRset was very small. But it existed. We re-analyze the dRset dependence data [6,7] by fitting with function

$$N(**9*) = Co + Cr (dRset) + Crr (dRset)^2. \quad (VII-1)$$

The Crr were listed in Table I. Here $N(**9*)$ is

$$\begin{aligned} N(\cos 9X) &= (10^5/2\pi) (e/cP) \Sigma \beta_x Q \cos 9\theta x \\ N(\sin 9X) &= (10^5/2\pi) (e/cP) \Sigma \beta_x Q \sin 9\theta x \\ N(\cos 9Y) &= (10^5/2\pi) (e/cP) \Sigma \beta_y Q \cos 9\theta y \\ N(\sin 9Y) &= (10^5/2\pi) (e/cP) \Sigma \beta_y Q \sin 9\theta y \end{aligned}$$

[10] and

$$Q = \int_s^{s+\Delta s} \partial B_y / \partial x \, ds \quad \text{Tesla}$$

$$\theta y = \int_0^s (1/\beta_y) ds \quad / \quad \int_0^{2\pi R} (1/\beta_y) ds.$$

The fitted values were comparable to errors or smaller than errors. But obviously the parabolic term existed. The sign of parabolic terms were always the same. It was not the non-linearity of dRset because the ratio of Crr to Cr was not constant. It was reasonable that the observed Crr on $2Qx=9$ and $2Qy=9$ had roughly the same amplitudes and the horizontal correction phase advanced to the vertical correction phase by a few tens degrees [12].

The down feed of the octupole imperfection to the quadrupole imperfection $N(\cos 9X)$ by the dispersion is

$$N(\cos 9X) = (10^5/2\pi) (e/cP) \Sigma \beta_x 3 O [\underline{h}(dP/P)]^2 \cos 9\theta x.$$

The expected octupole imperfections can be calculated as followings.

$$\Sigma O \cos 9\theta x = 319^2/3/(10^5/2\pi) (cP/e) Crr <\beta_x \underline{h}^2>$$

$$<\beta_x \underline{h}^2> \approx 28 \text{ m}^3$$

$$<\beta_y \underline{h}^2> \approx 21 \text{ m}^3$$

$$\begin{aligned} |\Sigma Oe^{9j\theta x}| &= 319^2/3/(10^5/2\pi) \cdot 3.3356 \sqrt{(13^2+10^2)} / 28 \\ &= 4.2 \text{ T/m}^2 \end{aligned}$$

$$\begin{aligned} |\Sigma Oe^{9j\theta y}| &= 319^2/3/(10^5/2\pi) \cdot 3.3356 \sqrt{(10^2+15^2)} / 21 \\ &= 6.1 \text{ T/m}^2 \end{aligned}$$

The three estimated strength of the octupole imperfection were roughly the same. They were not necessarily the same because the weight functions were not the same. But they should be close because the error sources were the same. The conclusion is;

the integrated (or accumulated) harmonic normal octupole imperfection at B=1.7kG and dB/dt=0 was about 5 T/m² .

An estimation error depends on the character of the field error. When the error sources are many (= N) and random, the estimation error is a random error; $1/\sqrt{N}$. And when error source was only one, the error is the same as a variation of twiss parameters;

$$(\underline{h}_{\max} - \langle \underline{h} \rangle) / \langle \underline{h} \rangle \text{ and } (\beta_{\max} - \langle \beta \rangle) / \langle \beta \rangle .$$

Table I

Dependence on second order term of dRset of the half integer corrections; Crr (/cm²).

Date	N(cos9X)	N(sin9X)	N(cos9Y)	N(sin9Y)
April 21	-6 ± 7	2 ± 7		
April 23	-9 ± 12	14 ± 17		
April 24	-20 ± 25	0 ± 25		
April 27	-9 ± 25	25 ± 20		
April 28			-9 ± 11	19 ± 11
April 30	-40 ± 30	20 ± 20		
May 7	-20 ± 20	10 ± 25	-12 ± 15	10 ± 15
weighted mean	-13 ± 7	10 ± 7	-10 ± 9	15 ± 9
(Cr (/cm)	-110	-40	-76	-25)

IV Octupole Imperfection by the Random magnet Variation

The octupole harmonic imperfection produced by the random variation of the magnets can be estimated using the equation

$$|\Sigma O e^{jn\theta}| = \sqrt{\langle (\Sigma O e^{jn\theta})^2 \rangle} \\ = \sqrt{N} \delta O_{rms}$$

The estimated values from each imperfections are listed in Table II. We used the rms errors reported by E.Blessner and R.Thern [13]. The strength of quadrupoles and sextupoles were set to the typical values. The down feed effect from the decapole field of the dipole magnets were also calculated. Here the rms of the horizontal displacement was assumed to be 1mm.

The calculated strength of octupole imperfection was only 0.36 T/m², which was much smaller than the observed strength. There should be the other source of the imperfection. The eddy current correction windings, which had strong higher order multipoles, did not work during the measurement because dB/dt was 0. The systematically distributed sextupoles could have produced only 6n-th harmonic octupole components and no 9th neither 14th harmonic components. The quadrupole magnets have rather large skew octupole component. But it is skewed and is systematic. A remanent field is a possible source of the imperfection. But we have no proof. The measurement of the slopes at the various B and dB/dt will give us more information about it.

Table II
Octupole imperfections produced by magnet imperfections.

magnet	N	(δO) rms/Br	reference	$ \sum O e^{jn\theta} $
variation of magnets				
B	36	$2\pi/36 \times 0.14$	13	0.35
QF	24	0.270×0.01	11,13	0.03
QD	24	0.276×0.01	11,13	0.03
SexF	24	0.157×0.02	13	0.04
SexD	24	0.237×0.02	13	0.05
down feed from decapole				
B	36	$2\pi/36 \times 9.8 \times 1.0E-3$	13	0.10

V Down Feed from the Sextupole Imperfections

We will check if the slope of quadrupole imperfections could predict the strength of sextupole imperfections. The strength of corrections were fitted with functions

$$N(\text{xxxx}) = C_o + C_b B + C_{bt} (dB/dt) \quad . \quad (II-1)$$

Here N(xxxx) is an computer control unit of correction selected for a convenience. The xxxx in the bracket is replaced by a name of resonance and a harmonic numbers. C_o , C_b and C_{bt} are fitting parameters. The unit of B, dB/dt, P are kG, G/ms=kG/s and GeV/c respectively.

The results of correction parameters were listed in Table III, which had reported in our study reports [6,9,14,15,16]. The strength of integrated correction field when N(xxxx)=1 is listed in Table IV. The strength of harmonic sextupole and skew sextupole imperfections were calculated from parameters listed in Table III and Table IV. The results were listed in Table V.

The strength of imperfections calculated from the observation of different resonances were roughly the same. The Cbt of normal sextupole field did not agreed with each other. That suggested that the error field located mainly at the low dispersion point.

$$\underline{h} \approx (2.59+2.74)/(7.57+7.0) <\underline{h}> \\ \approx 0.62 \text{ m}$$

And β_x/β_y might be close to 1 because the strength of $3Q_x=14$ and $Q_x+2Q_y=14$, and the slope of $2Q_x=9$ and $2Q_y=9$ were close. The location which satisfies the above conditions is dipole magnet at 1 and 8 of each super period. The Co of skew sextupole field also had disagreement. The remanent error field might locate at high dispersion and high β_y point. The random error could have produced that kind of disagreement. But the probability was very low.

The strength of normal and skew sextupole imperfections were roughly the same for Co and Cb. But the Cbt of normal sextupoles was larger than that of skew sextupoles. Which suggested the imperfection of eddy current correction system.

Table III

Stop band correction parameters.

resonance string		Co	Cb (/kG)	Cbt (/(G/s))
normal sextupole field				
3Qx=14	N(cos14X)	48±70	-31±34	3.49±0.43
	N(sin14X)	-129±34	40±16	6.00±0.20
Qx+2Qy=14	N(cos14XY)	5±29	14±11	4.74±0.20
	N(sin14XY)	-103±24	17± 9	2.64±0.19
2Qx=9	δN(cos9X)/δdRset	75±40	-3±12	1.06±0.29
	δN(sin9X)/δdRset	52±40	13±12	0.45±0.29
2Qy=9	δN(cos9Y)/δdRset	49±25	21± 9	0.94±0.18
	δN(sin9Y)/δdRset -	22± 9	1± 3	-0.44±0.06
skew sextupole field				
2Qx+Qy=14	SV4*20	720±120	-152±42	6.82±0.70
	SH4*20	604± 81	30±30	-0.30±0.64
Qx+Qy=9	δN(cos9XY)/δdRset	19.9±1.0	-0.4±0.6	0.024±0.03
	δN(sin9XY)/δdRset	9.8±1.0	1.6±0.6	0.044±0.03

Table IV

Strength of correction parameters.

resonance	string	strength(mT/m)	definition
normal sextupole			
3Qx=14	N(14X)	1.09	$\Sigma \delta S \ Bx/Bx \ e^{j14\theta} / \langle Bx/Bx \rangle$
Qx+2Qy=14	N(14XY)	1.29	$\Sigma \delta S \ By/Bx \ e^{j14\theta} / \langle By/Bx \rangle$
2Qx=9	$\delta N(9X)/\delta Rset$	2.25	$\Sigma \delta S \ Bx_h \ e^{j9\theta} / \langle Bx_h \rangle$
2Qy=9	$\delta N(9Y)/\delta Rset$	2.64	$\Sigma \delta S \ By_h \ e^{j9\theta} / \langle By_h \rangle$
skew sextupole			
2Qx+Qy=14	SV4*20	.160	$\Sigma \delta S' Bx/By \ e^{j14\theta} / \langle Bx/By \rangle$
	SH4*20	.163	$\Sigma \delta S' Bx/By \ e^{j14\theta} / \langle Bx/By \rangle$
Qx+Qy=9	$\delta N(9XY)/\delta Rset$	21.3	$\Sigma \delta S' /Bx/By_h \ e^{j9\theta} / \langle /Bx/By_h \rangle$

Table V

Strength of sextupole imperfections

imperfection sextupole	resonance	Co T/m	Cb T/m/kG	Cbt T/m/(G/s)
normal	3Qx=14	150	55	7.57
	Qx+2Qy=14	130	28	7.00
	2Qx=9 (slope)	205	30	2.59
	2Qy=9 (slope)	141	55	2.74
skew	2Qx+Qy=14	100	25	1.09
	Qx+Qy=9 (slope)	472	35	1.07

REFERENCES

- [1] Y.Shoji and C.Gardner,
'Tune Space Survey at Low Intensity (1)', AGS SR-290, 1993.
- [2] Y.Shoji, C.Gardner and C.Whalen,
'Observed Loss by 4th Resonances', AGS SR-291, 1993.
- [3] Y.Shoji and C.Gardner,
'Tune Space Survey at Low Intensity (2)', AGS SR-296, 1993.
- [4] Y.Shoji and C.Gardner,
'Intensity Dependence of resonances', AGS SR-297, 1993.
- [5] G.Danby,
informal meeting on Oct.4, 1993.
- [6] Y.Shoji and C.Gardner,
' $2Q_x=9$ Correction Data Before May 7', AGS SR-287, 1993.
- [7] Y.Shoji and C.Gardner,
' $2Q_y=9$ Correction Data Before May 7', AGS SR-288, 1993.
- [8] Y.Shoji and C.Gardner,
'Integer Coupling ($Q_x+Q_y=9$) Correction Data', AGS SR-289, 1993.
- [9] Y.Shoji and C.Gardner,
'14th Normal Sextupole Correction', AGS SR-286, 1993.
- [10] C.Gardner,
'Booster Stopband Corrections', Booster TN-217, Jan.6, 1993.
- [11] A.Luccio and M.Blaskiewicz,
'AGS Booster Parameters (MAD Output)', Booster TN-196, July 23, 1991.
- [12] Y.Shoji and C.Gardner,
'Summary of the stop band correction parameters', to be reported as AGS/AD Tech Note.
- [13] E.Blessner and R.Thern
'Analysis of Magnetic Field Measurement Results for the AGS Booster Magnets', IEEE PAC 1991, p.45.
- [14] Y.Shoji and C.Gardner,
'9th Skew Sextupole Correction Test', AGS SR-294, 1993.
- [15] Y.Shoji and C.Gardner,
'Skew Sextupole Correction for $3Q_y=14$ and $2Q_x+Q_y=14$ ', AGS SR-295, 1993.
- [16] Y.Shoji and C.J.Gardner,
'Quadrupole and Sextupole Correction Parameters for $2Q_y=9$ ', AGS SR-298, 1993.